

## Introduction of OMEG Institute

- 천체와 우주의 진화에서 보는 원소 및 우주물질의 기원 -

**Myung-Ki Cheoun**

(**Soongsil University**, **OMEG Institute**, Seoul, Korea)

# Since 2021

## 우주물질연구소

OMEG institute

### 원소 및 물질의 기원 연구

#### Fermi

희귀동위원소에 대한  
이론 모형 연구

핵자 사이의 핵력  
(BCS 이론 등)

핵구조 및 열핵반응

핵구조 및 열핵반응  
이론 연구로 확장

고밀도 핵물질

### 항성 및 은하의 진화 연구

#### Hubble

중성미자의 역할 연구

입자 및 강입자  
이론 연구 기반

초기우주 시뮬레이션

자기유체역학

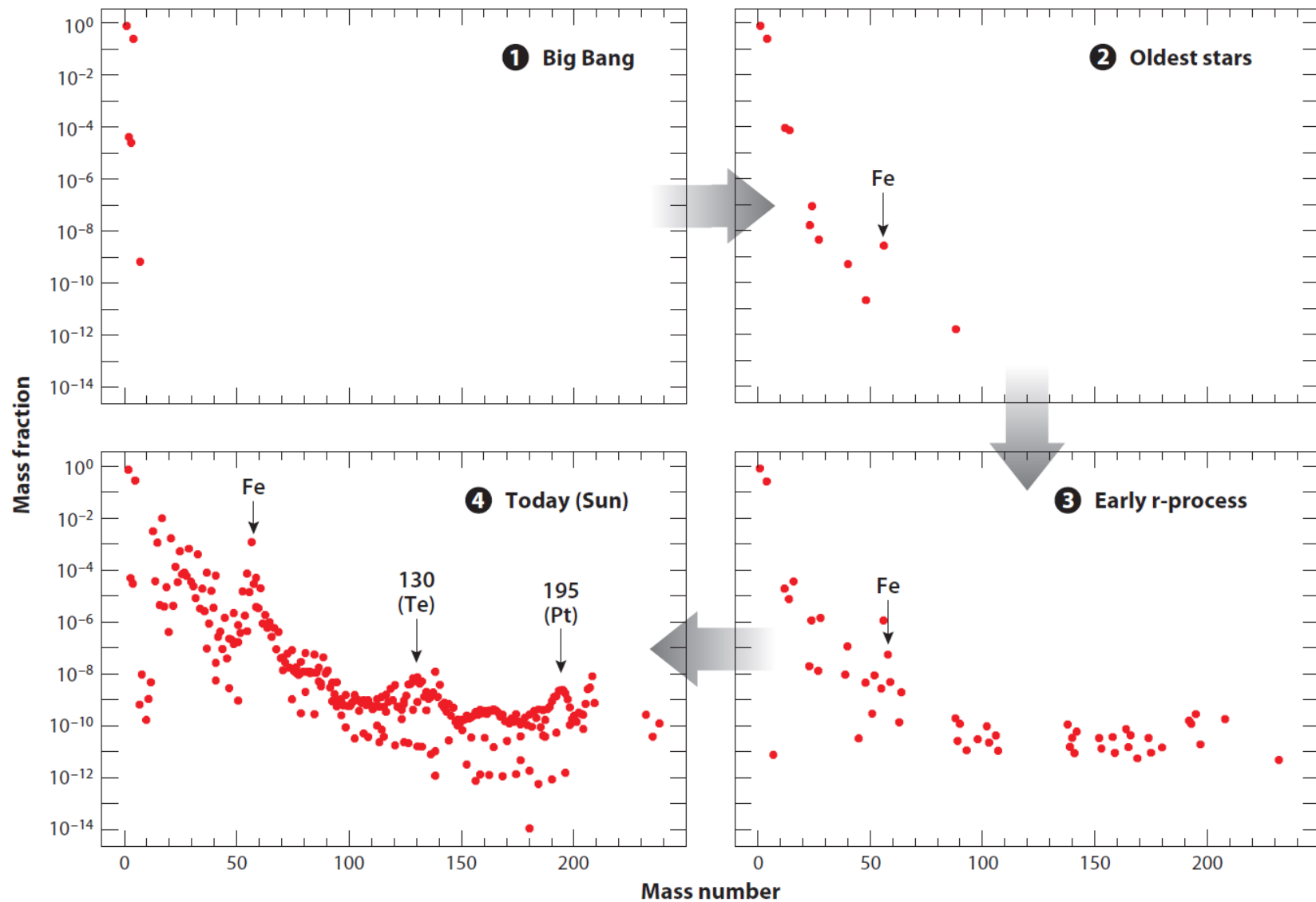
마그네타 하전 입자의  
양자화 가능성

우주를 구성하는 원소와 물질의 기원에 대한 종합적 이해

## Periodic Table



# Preliminary Evolution of Element Abundances in the Universe Evolution



M. Wiescher et al., Annu. Rev. Astro. Astrophys. (2012)



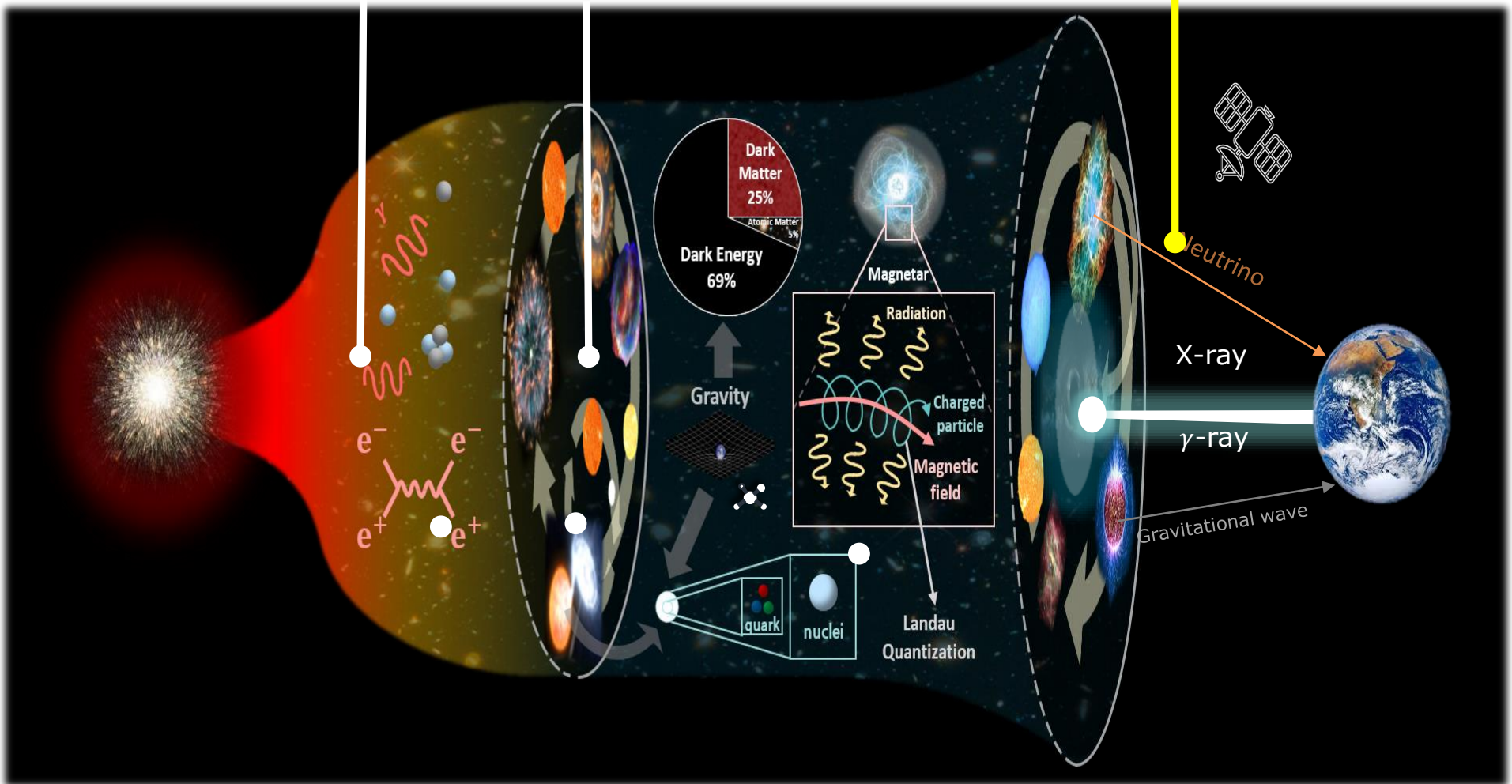
# Preliminary Cosmological Sources of elements in the Universe Evolution

FAUST (Femtoscale Astrophysics for the Universe Study)

Early Universe & Big Bang

Stellar and Galaxy Evolution

Multi-messenger Astrophysics

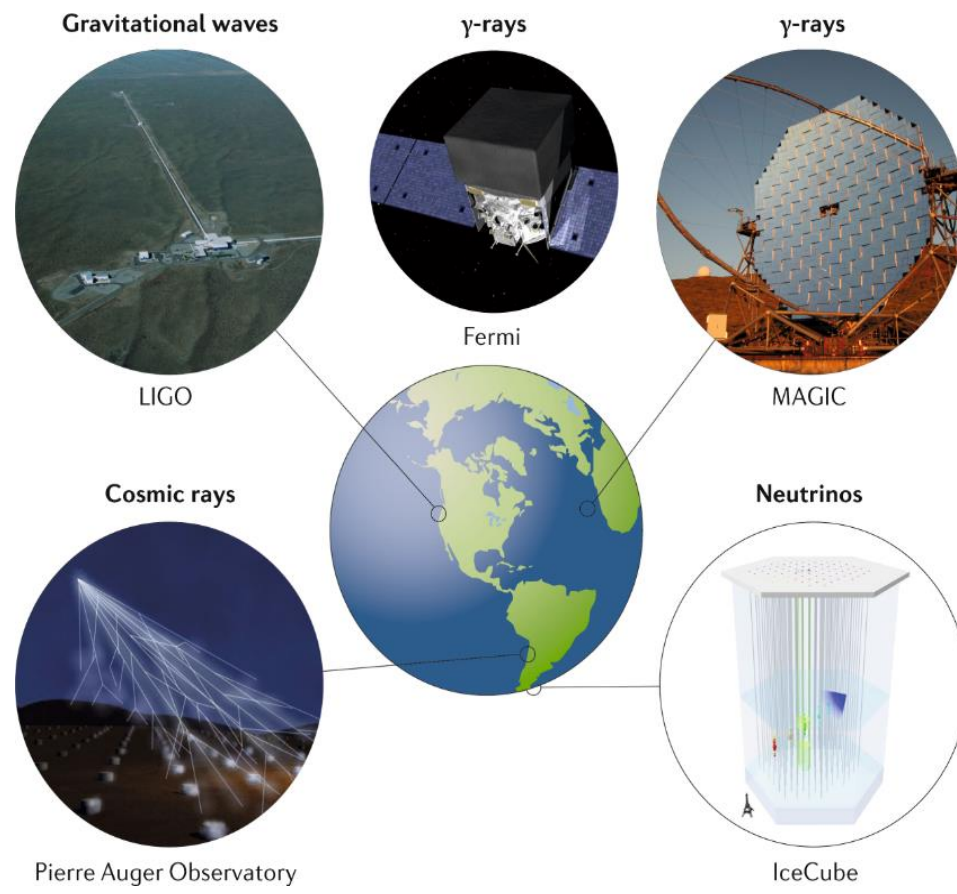
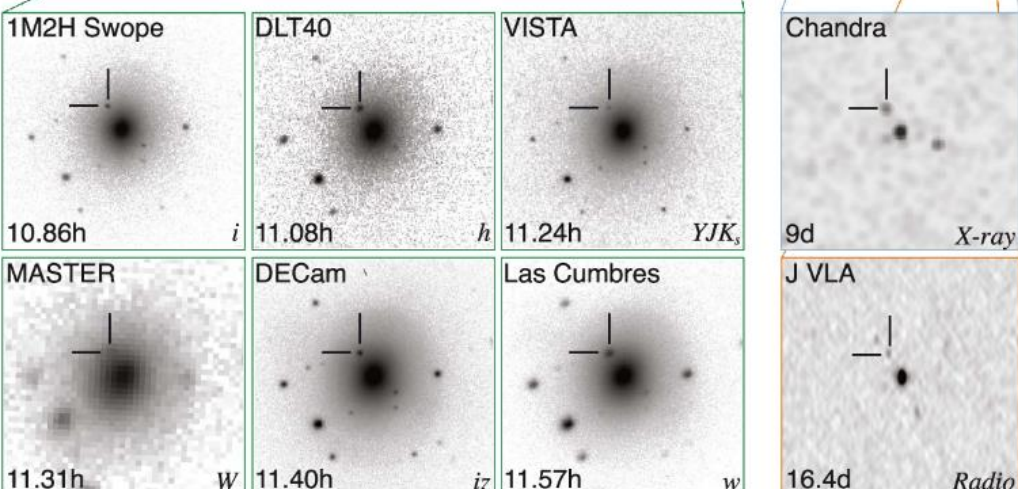
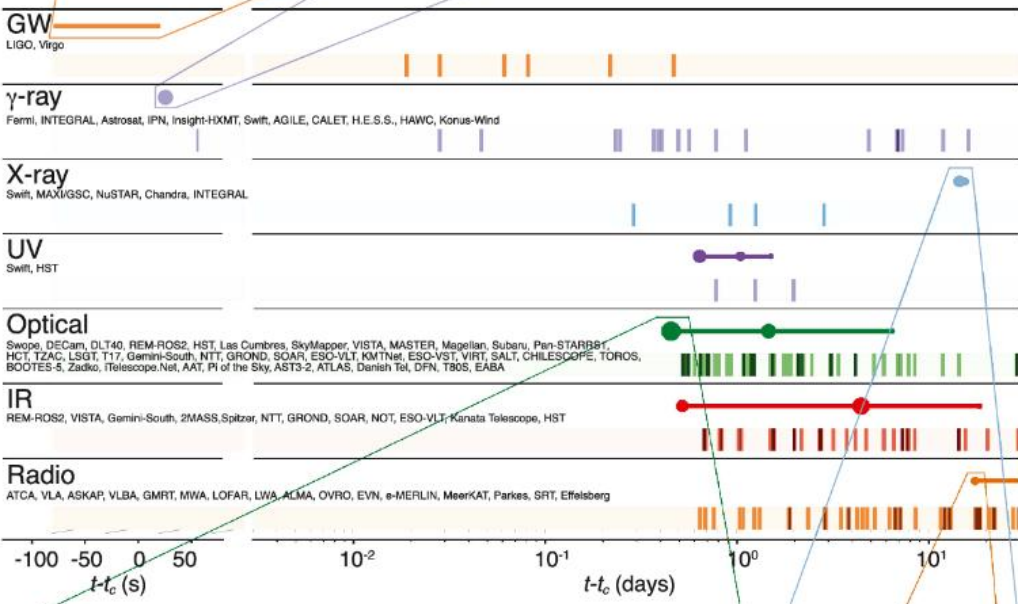
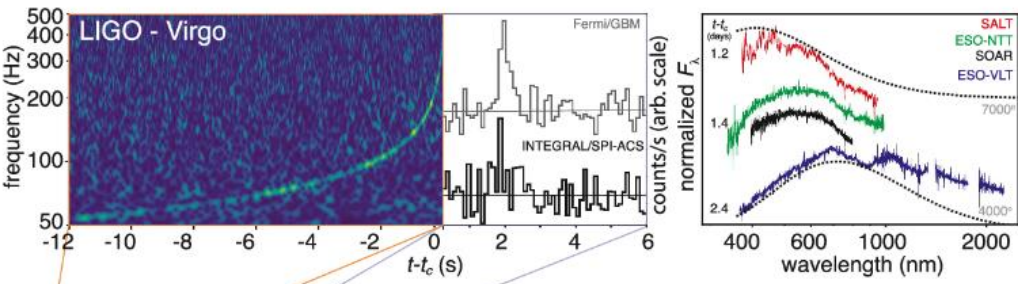


Primordial Nucleosynthesis

Nucleosynthesis in the Stellar evolution

Galaxy Chemical Evolution & Exotica Matter

# Preliminary Multi-messenger Astrophysics



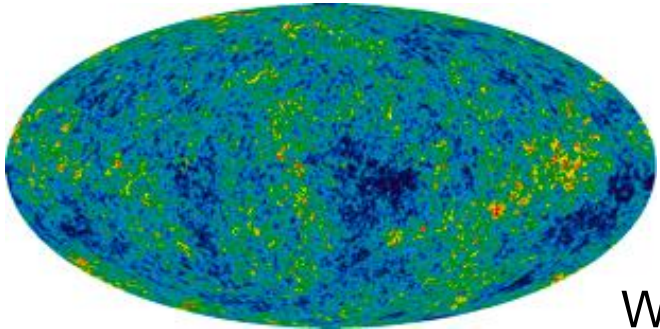
...e, Nov. 16-18,

**Big Bang Nucleosynthesis**

**Primordial abundances in BBN**



- Standard big bang nucleosynthesis (SBBN) parameter: baryon-to-photon ratio  $\eta$
- **Observation** of CMB  $\rightarrow$  constraint on  $\eta$



WMAP Science Team

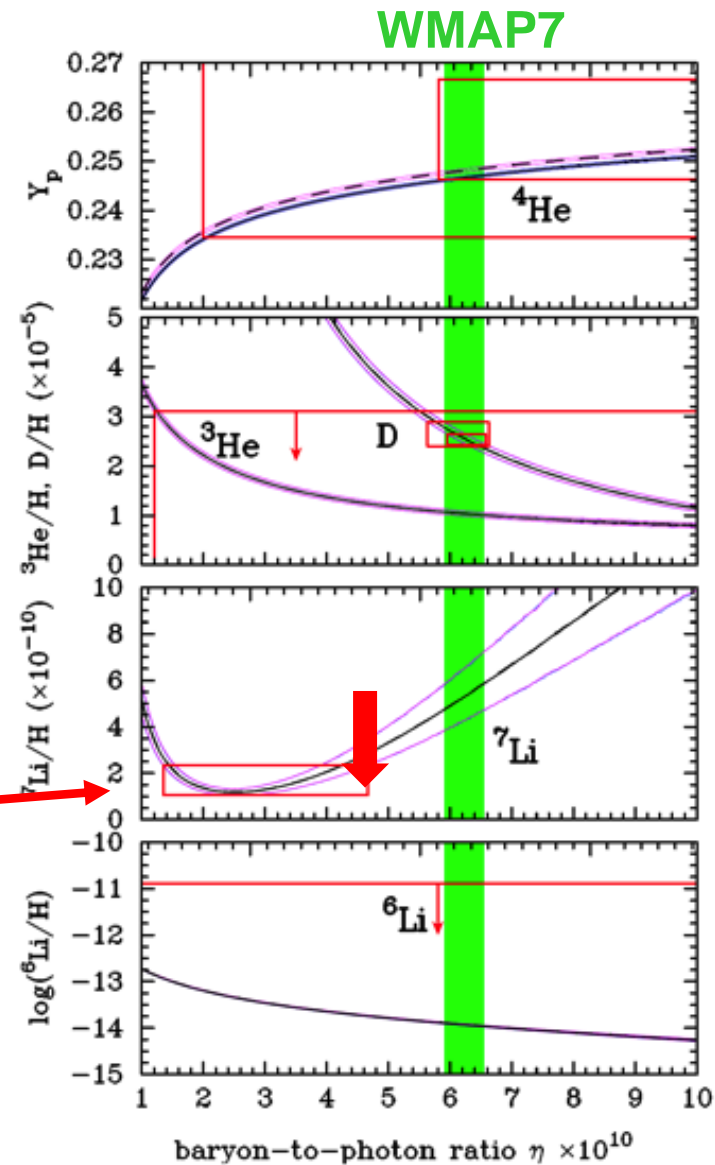
- **Observation** of metal-poor stars (MPSs)
  - ✓  ${}^7\text{Li}$  abundance is smaller than theory by a factor of  $\sim 3$

**Signature of new physics?**

## Goals

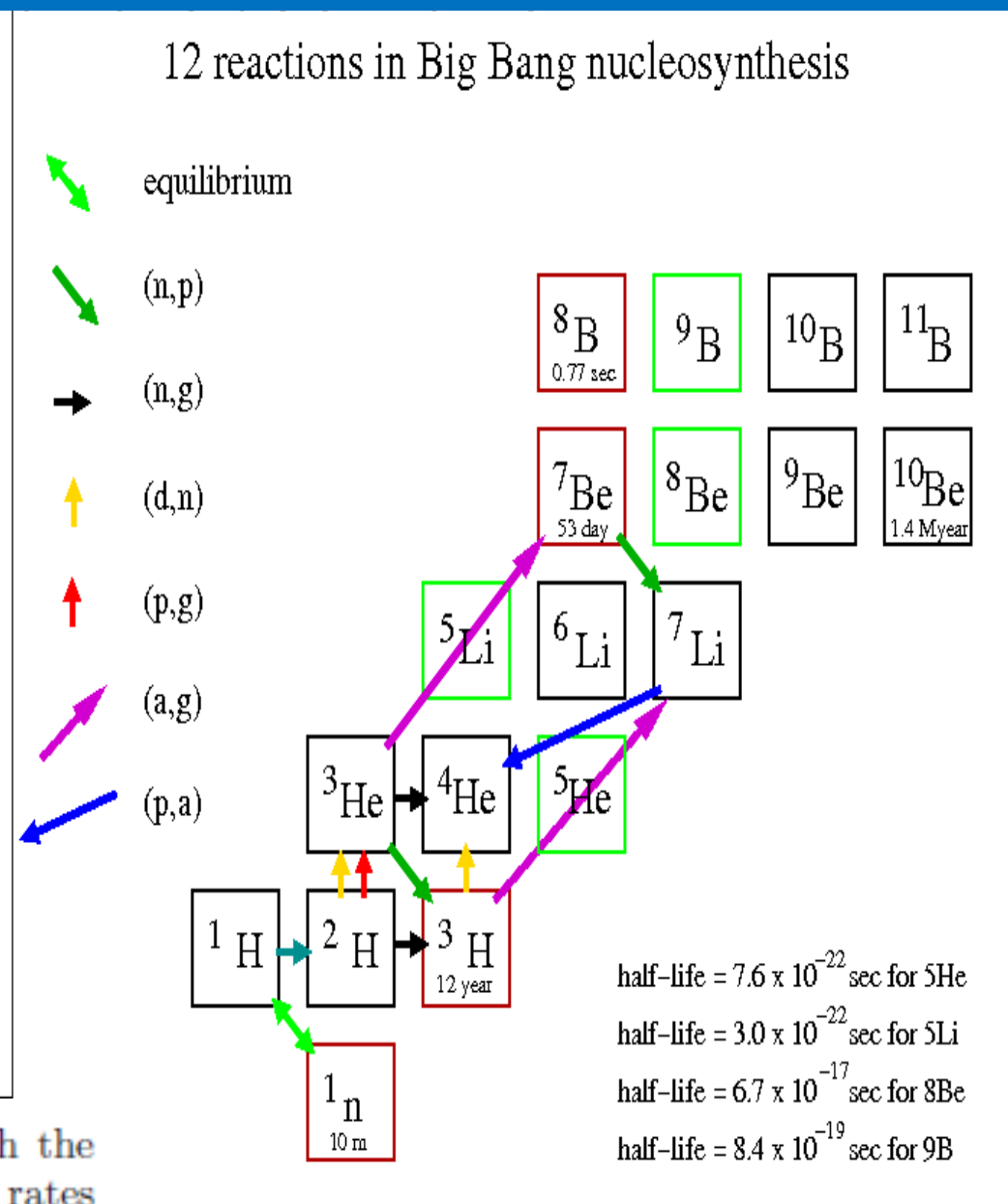
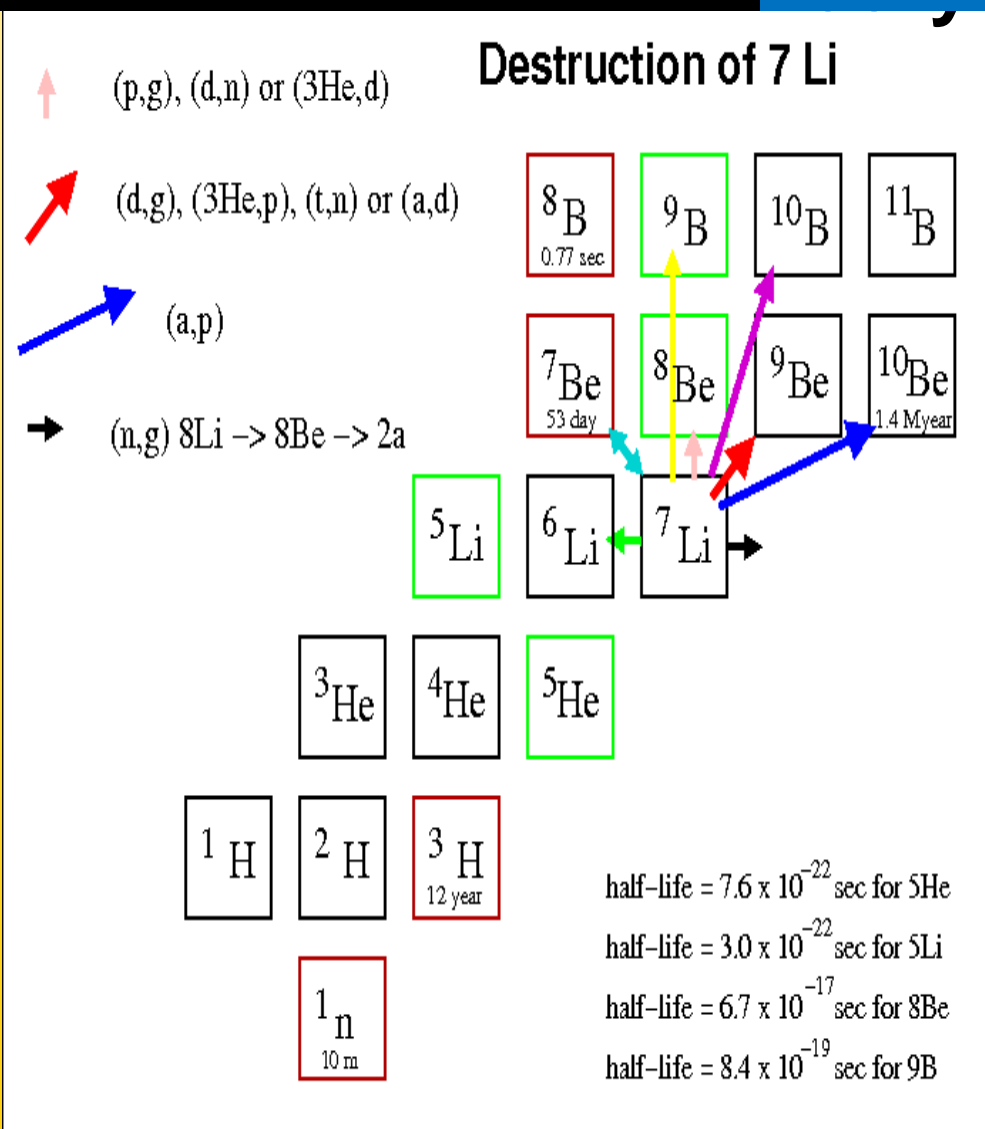
- Find the solution of the  ${}^7\text{Li}$  problem, & identify the processes in the early universe
- Derive constraints on **particle models** and **modified GR ??**

KSHEP Fall Conference, Nov. 16-18, 2022



Kawasaki & MK, PRD 86, 063003 (2012)

# Big Bang Nucleosynthesis Main Nuclear Reactions for BBN



We use the BBN code by Kawano [86, 87] with the Sarkar's correction [88] to  ${}^4\text{He}$  abundance. Reaction rates relating to light nuclei of mass number  $A \leq 10$  are updated with the JINA REACLIB Database V1.0 [89]. We adopt the neutron lifetime of  $878.5 \pm 0.7_{\text{stat}} \pm 0.3_{\text{sys}}$  s [90].

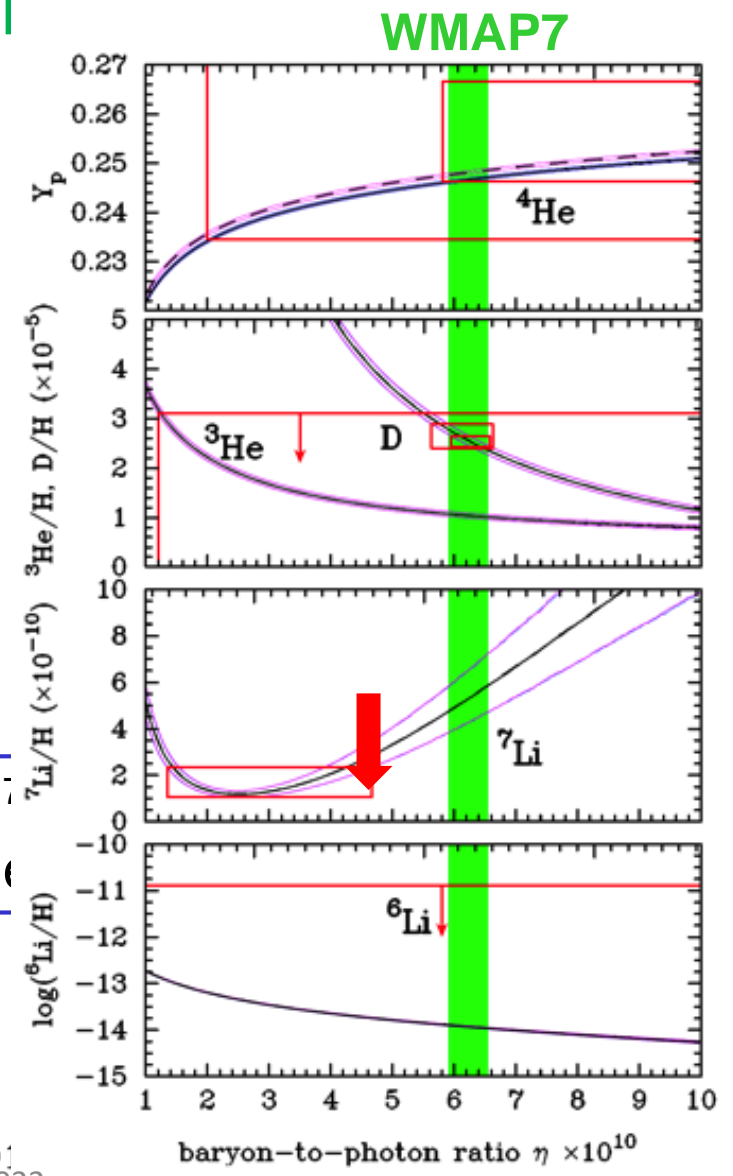
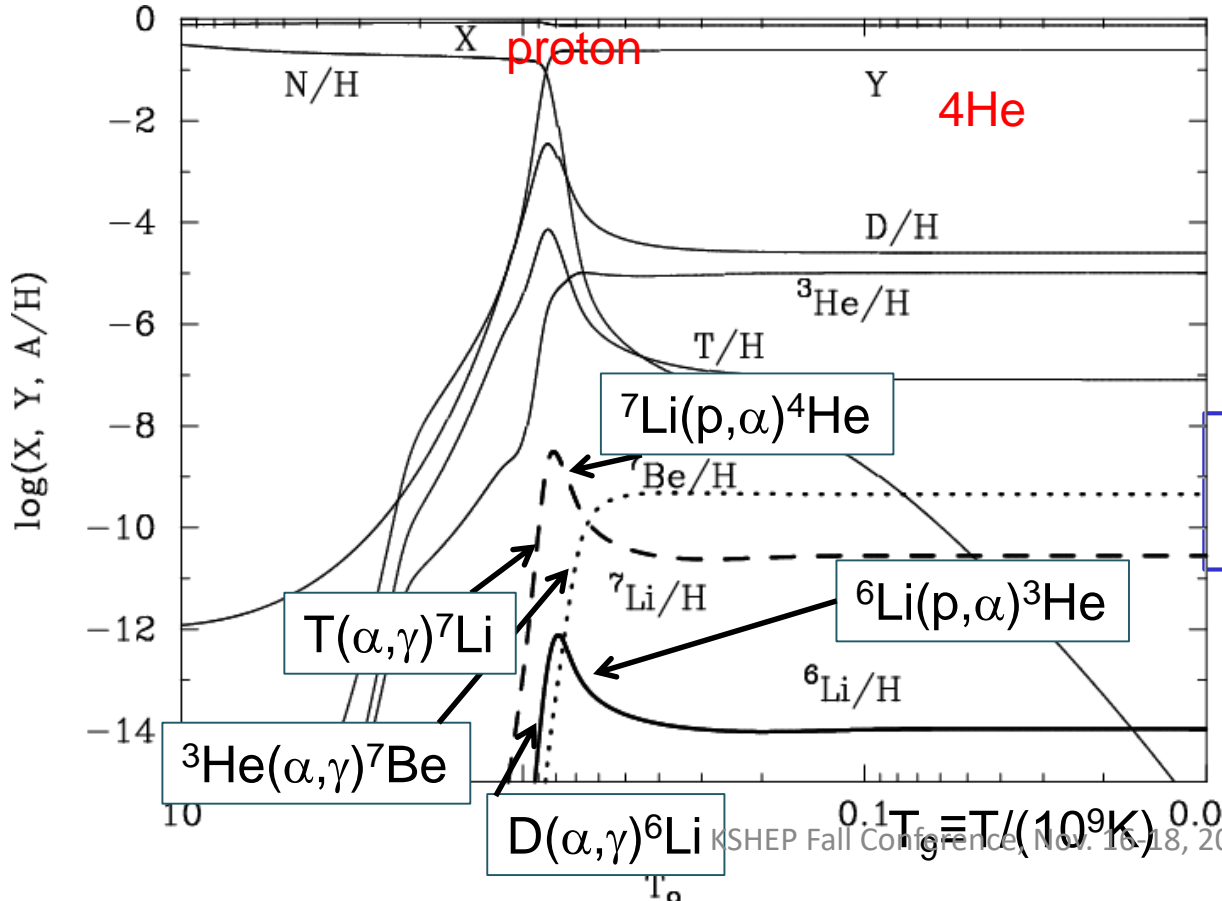
KSHEP Fall Conference, Nov. 16-18, 2022

[89] R. H. Cyburt, A. M. Amthor, R. Ferguson, Z. Meisel, K. Smith, S. Warren, A. Heger, R. D. Hoffman, T. Rauscher, A. Sakharuk, et al., *Astrophys. J. Suppl.* **189**, 240 (2010).

- $n \leftrightarrow p$  equilibrium  $(n/p)_{EQ} = \exp(-Q/T)$   $Q \equiv m_n - m_p = 1.293 \text{ MeV}$
- $t \sim 1 \text{ sec}, T = T_F \sim 1 \text{ MeV}$  (weak interaction freeze-out)

- ✓  $\nu \bar{\nu} \leftrightarrow e^+ e^- \leftrightarrow \gamma \gamma$
- ✓  $n \leftrightarrow p$
- ✓  $e^\pm \rightarrow \gamma \gamma$  ( $T \sim m_e/3$ )

$$(n/p)_{\text{freeze-out}} = \exp(-Q/T_F) \sim 1/6$$



$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} f(R) + S_m(g_{\mu\nu}, \phi_m), \quad (1)$$

where  $\kappa^2 = 8\pi G$  is defined, with  $G$  Newton's constant,  $g_{\mu\nu}$  the metric tensor,  $g$  the determinant of the metric tensor

and  $S_m$  the action for matter, account radiations, equation for radiation with respect

$f'R_{\mu\nu}$

where  $f' =$  momentum

Here,  $\mathcal{L}_m$  is related to  $S_m$

Even though Einstein's general relativity (GR) has successfully passed the observational tests in the solar system scale, lots of efforts to generalize GR for cosmology have been continued. One of them is to introduce additional higher-order derivative terms due to both theoretical and phenomenological reasons. For instance, sixth order  $R\Box R$  [1] as well as fourth order  $R^{ab}R_{ab}$  terms [2–7, 17] are considered here with the action

$$S = \int d^4x \sqrt{-g} \left[ \frac{1}{16\pi} (f(R) + BR^{ab}R_{ab} + CR\Box R) + L_m \right], \quad (1)$$

where  $f(R)$  is a polynomial function of the Ricci scalar  $R$

$$f(R) = \sum_{n=1}^N A_{(n)} R^n = R + A_{(2)} R^2 + A_{(3)} R^3 + \dots, \quad A_{(1)} \equiv 1, \quad A_{(2)} \equiv A, \quad (2)$$

$A_{(n)}, B, C$  are constants,  $R_{ab}$  is the Ricci tensor. The d'Alembertian of  $R$  is  $\Box R \equiv g^{ab}R_{;a;b} = R^{;c}_c$  where the semicolon denotes covariant derivative, and the matter part Lagrangian is defined as  $\delta(\sqrt{-g}L_m) \equiv \frac{1}{2}\sqrt{-g}T_{(m)}^{ab}\delta g_{ab}$ . In this paper we follow the Hawking-Ellis [5]

We assume the spatially flat Friedmann-Lemaître-Robertson-Walker metric, as supposed in the standard cosmological model,

$$ds^2 = dt^2 - a(t)^2(dx^2 + dy^2 + dz^2). \quad (5)$$

For matter, on the other hand, we assume a perfect fluid described by a time-dependent energy density  $\rho(t)$  and pressure  $p(t)$ ,

$$T^\mu_\nu = \text{diag}(\rho, -p, -p, -p). \quad (6)$$

The 0-0 component of Eq. (2) then becomes

$$-3\frac{\ddot{a}}{a}f' - \frac{1}{2}f + 3\frac{\dot{a}}{a}f''\dot{R} = \kappa^2\rho. \quad (7)$$

giving

$$-f''\dot{R} = \kappa^2\rho. \quad (8)$$

energy conservation

$$(10)$$

$M_p^{2-2/n}$ , with  $M_p$  = power-law index  $n$  reduces to Einstein's been analyzed in the scale factor exist s model.

$$(11)$$

A formulation of this model is given using the same assumptions as those adopted in Refs. [1,2], as follows. It is assumed that the matter part is predominantly contributed by the radiation with  $p = \rho/3$ . In this case, Eq. (9) leads to a relation of  $\rho \propto a^{-4}$ . Here, we additionally constrain the model space by assuming the power-law solution of the scale factor, i.e.,

$$a(t) \propto t^\alpha. \quad (12)$$

It is found that  $\alpha = n/2$  must be satisfied in order to hold Eqs. (13) and (14) for any time  $t$ . Then we assume  $\alpha = n/2$



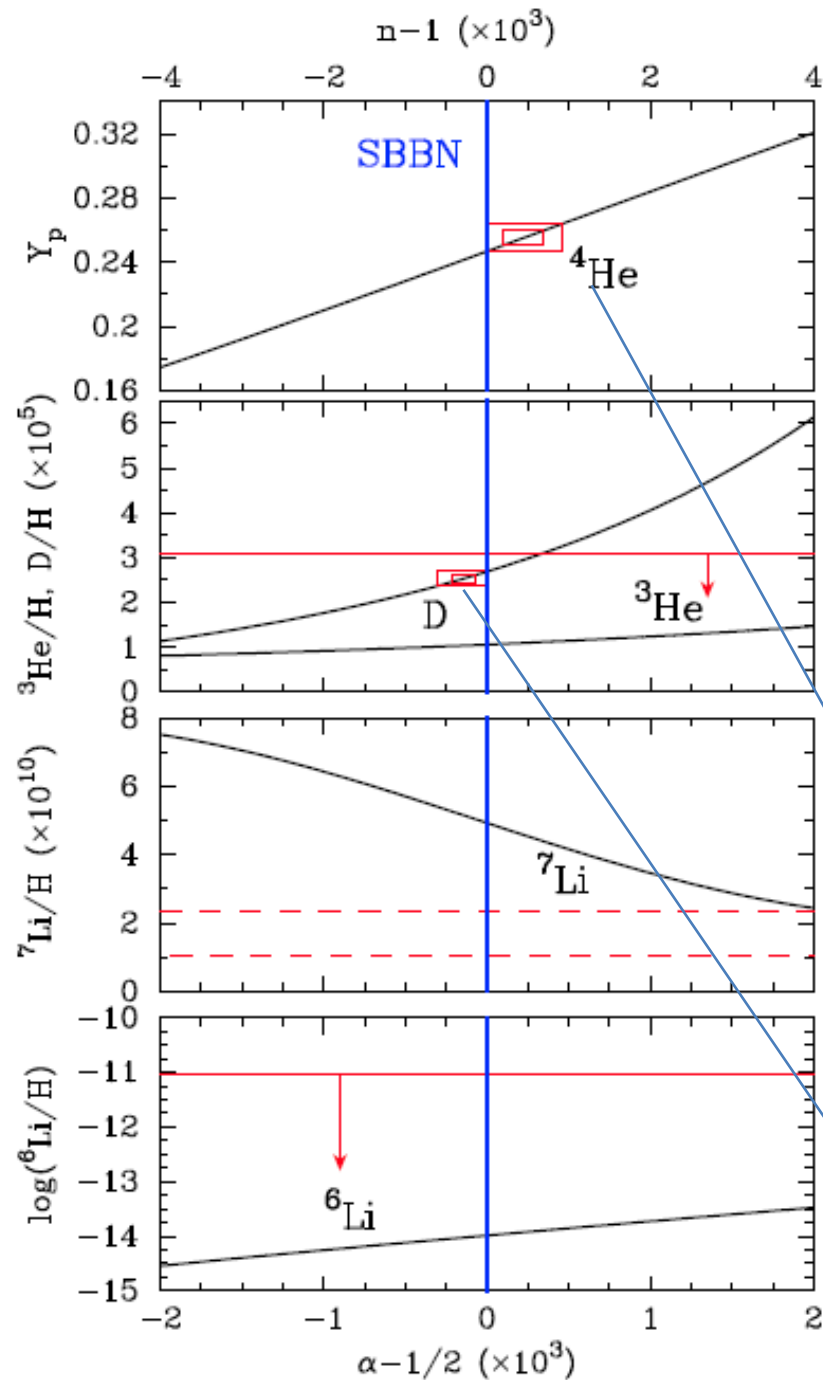


FIG. 1 (color online).  ${}^4\text{He}$  mass fraction  $Y_p$  and number abundance ratios of D,  ${}^3\text{He}$ ,  ${}^7\text{Li}$ , and  ${}^6\text{Li}$  relative to H as a function of the power-law index of the scale factor  $\alpha - 1/2$  or the  $f(R)$  function  $n - 1$ . Solid curves show calculated results for the  $f(R) \propto R^n$  model. The solid smaller and larger boxes of  ${}^4\text{He}$  and D abundances correspond to the  $2\sigma$  and  $4\sigma$  limits, respectively, from adopted observational constraints. The dashed box corresponds to the  $2\sigma$  limits of  ${}^7\text{Li}$  abundance. The horizontal lines with downward arrows of  ${}^3\text{He}$  and  ${}^6\text{Li}$  abundances show observational upper limits. The result on the vertical line for  $\alpha = 1/2$  or  $n = 1$  is for the SBBN model. Note that the  $f(R)$  function for  $\alpha < 1/2$  was corrected in this paper so that the cosmological model connects to the present  $\Lambda\text{CDM}$  model.

result of a changed expansion rate as follows. First, when the expansion rate is larger, the freeze-out of weak reactions occurs earlier. The neutron abundance remaining after the freeze-out is then higher. Second, the time interval between the freeze-out and the  ${}^4\text{He}$  synthesis is shorter because of faster cosmic expansion. Neutron abundances are larger because of the above two reasons. Almost all neutrons are processed to form  ${}^4\text{He}$  nuclei at the  ${}^4\text{He}$  synthesis epoch. The  ${}^4\text{He}$  abundance is therefore larger for larger values of  $\alpha$ .

$$-1 \times 10^{-5} \lesssim (\alpha - 1/2) \lesssim 5 \times 10^{-4}$$

$$-2 \times 10^{-5} \lesssim (n - 1) \lesssim 10^{-3}. \quad (28)$$

When the  $4\sigma$  limit of D is used, however, more stringent constraints are derived,

$$-3 \times 10^{-4} \lesssim (\alpha - 1/2) \lesssim 2 \times 10^{-6}$$

$$-6 \times 10^{-4} \lesssim (n - 1) \lesssim 4 \times 10^{-6}. \quad (29)$$



**Motivation**

**Stellar evolution & nucleosynthesis**

## Life cycle of stars

5. Nuclear Abundances

Seeds for the next generation stars

4. Sites for Nucleosynthesis

3. Nucleosynthesis

Shock wave by sudden gravitational collapse and neutrino driven wind

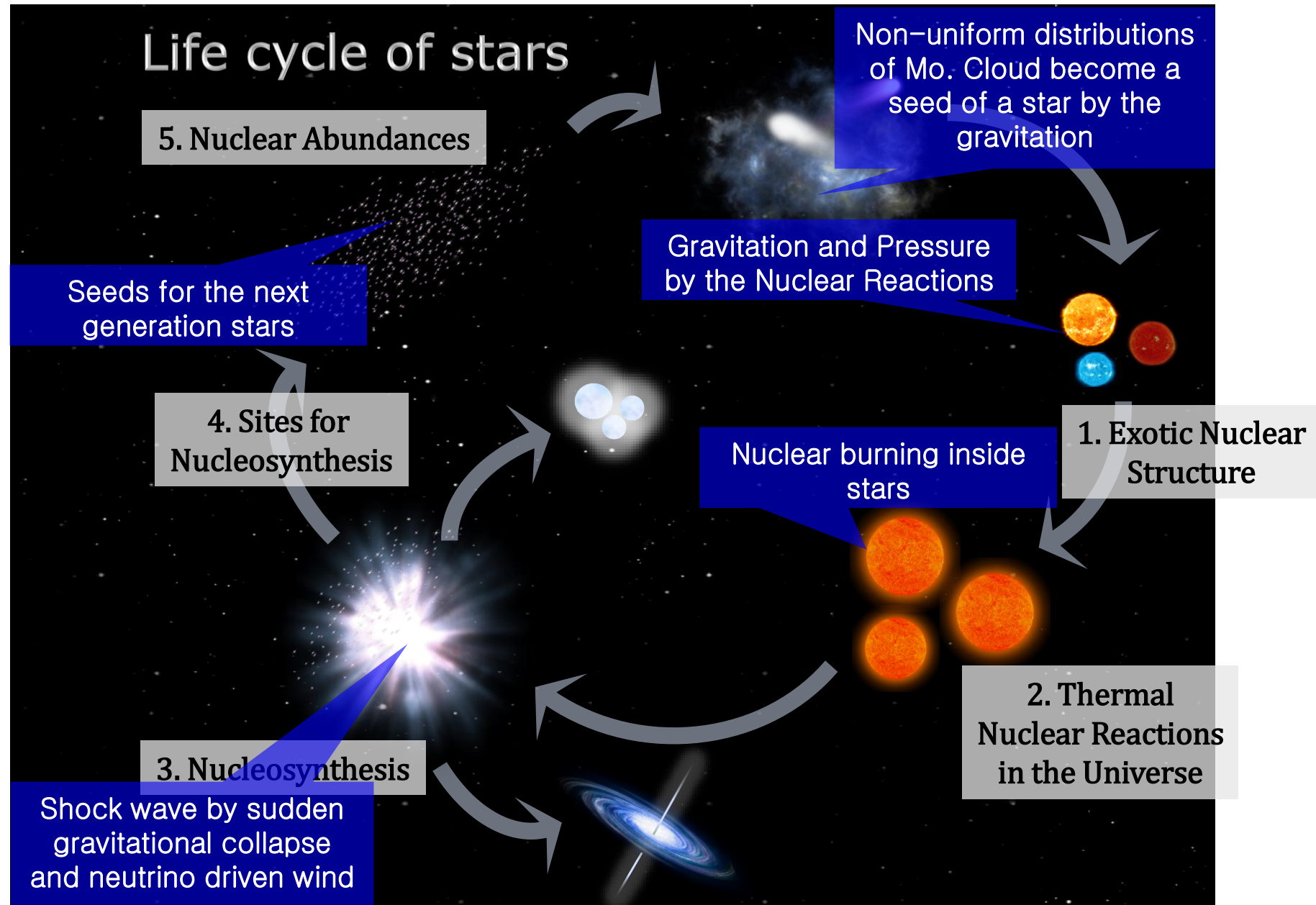
Nuclear burning inside stars

2. Thermal Nuclear Reactions in the Universe

1. Exotic Nuclear Structure

Non-uniform distributions of Mo. Cloud become a seed of a star by the gravitation

Gravitation and Pressure by the Nuclear Reactions



# Motivation Exotic nuclear structures and reactions for nucleosynthesis

## Life cycle of stars

Seeds for the next generation stars

Non-uniform distributions of Mo. Cloud become a seed of a star by the gravitation

Gravitation and Pressure by the Nuclear Reactions

### 1. Exotic Nuclear Structure

### 2. Thermal Nuclear Reactions

### 3. Nucleosynthesis

Deformed QRPA, RMF, SHF DFT ...

Nuclear Transitions

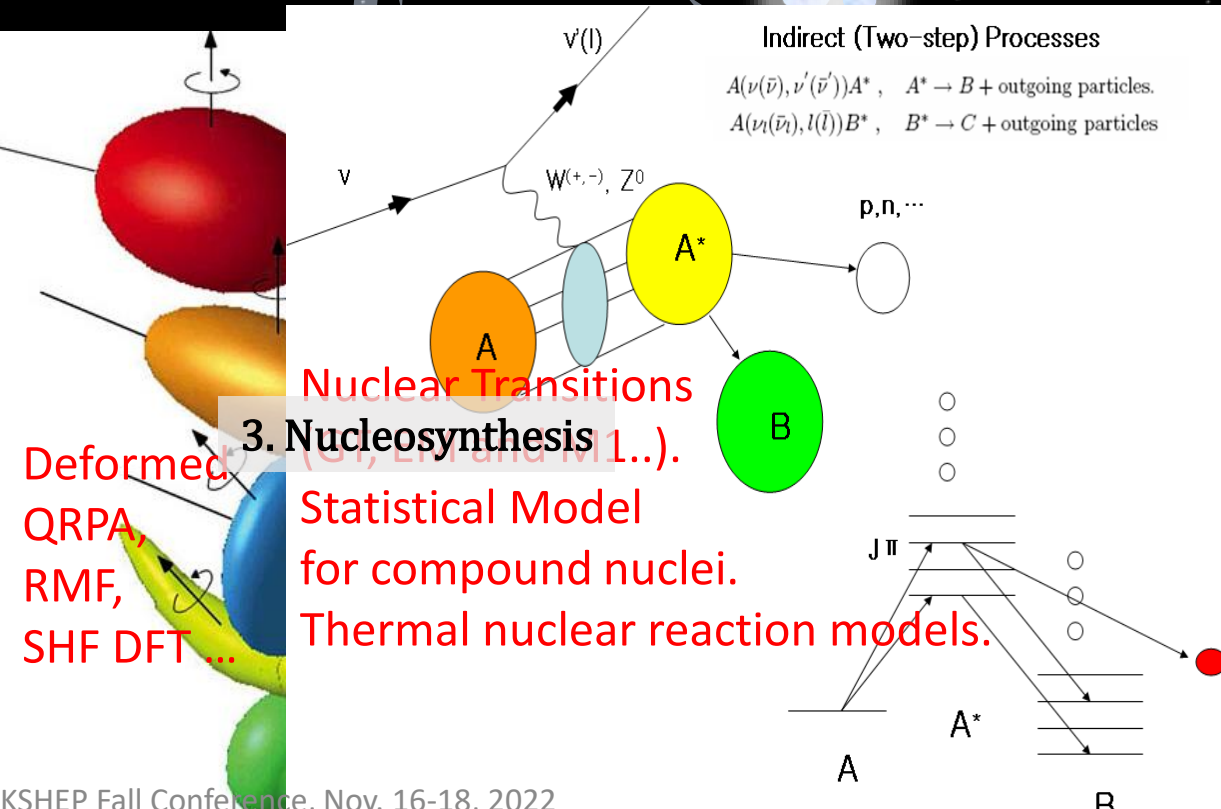
Statistical Model for compound nuclei.

Thermal nuclear reaction models.

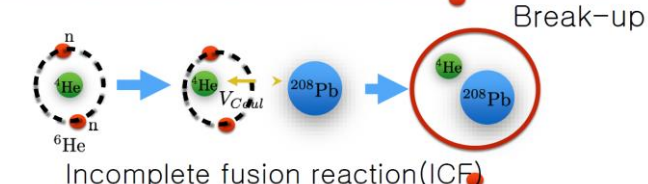
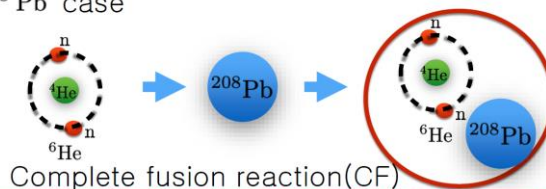
#### Indirect (Two-step) Processes

$A(\nu(\bar{\nu}), \nu'(\bar{\nu}'))A^* \rightarrow B + \text{outgoing particles.}$

$A(\nu_l(\bar{\nu}_l), l(\bar{l}))B^* \rightarrow C + \text{outgoing particles}$



${}^6\text{He} + {}^{208}\text{Pb}$  case

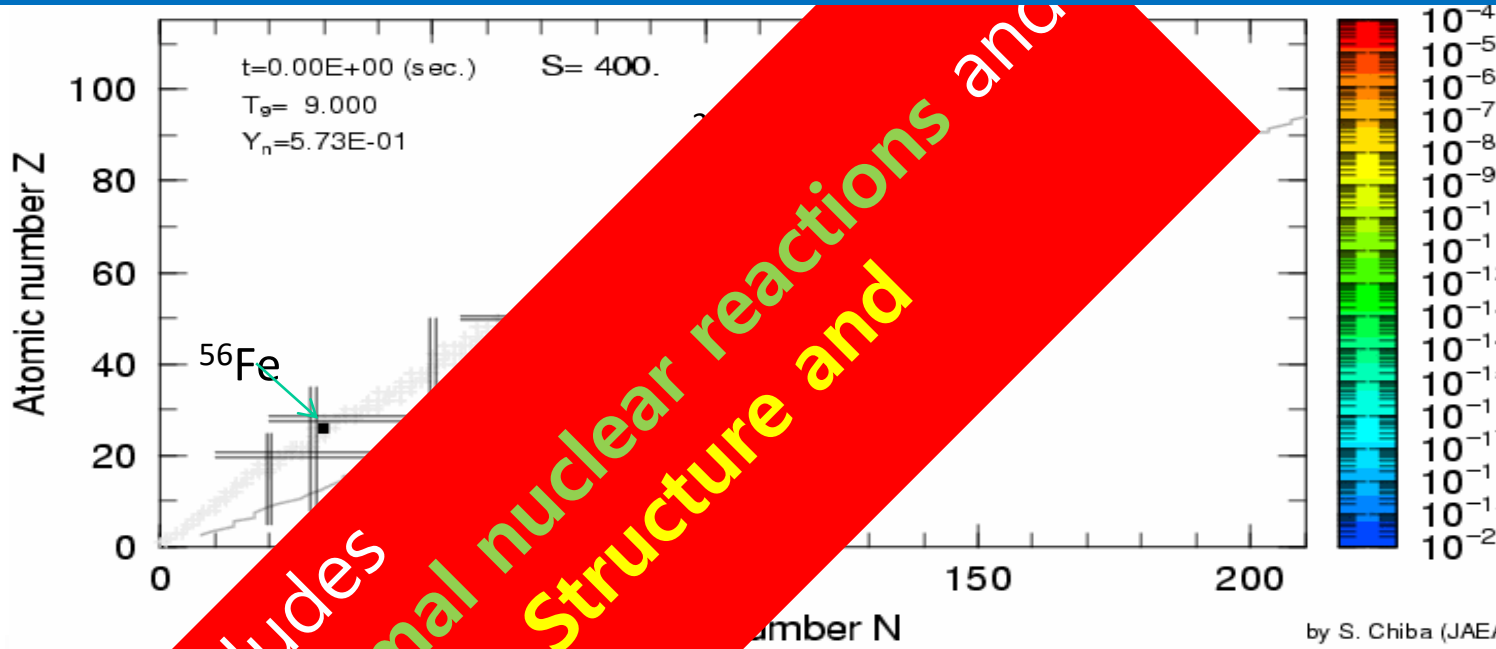
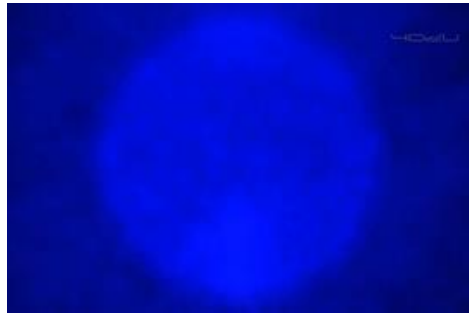


Break-up

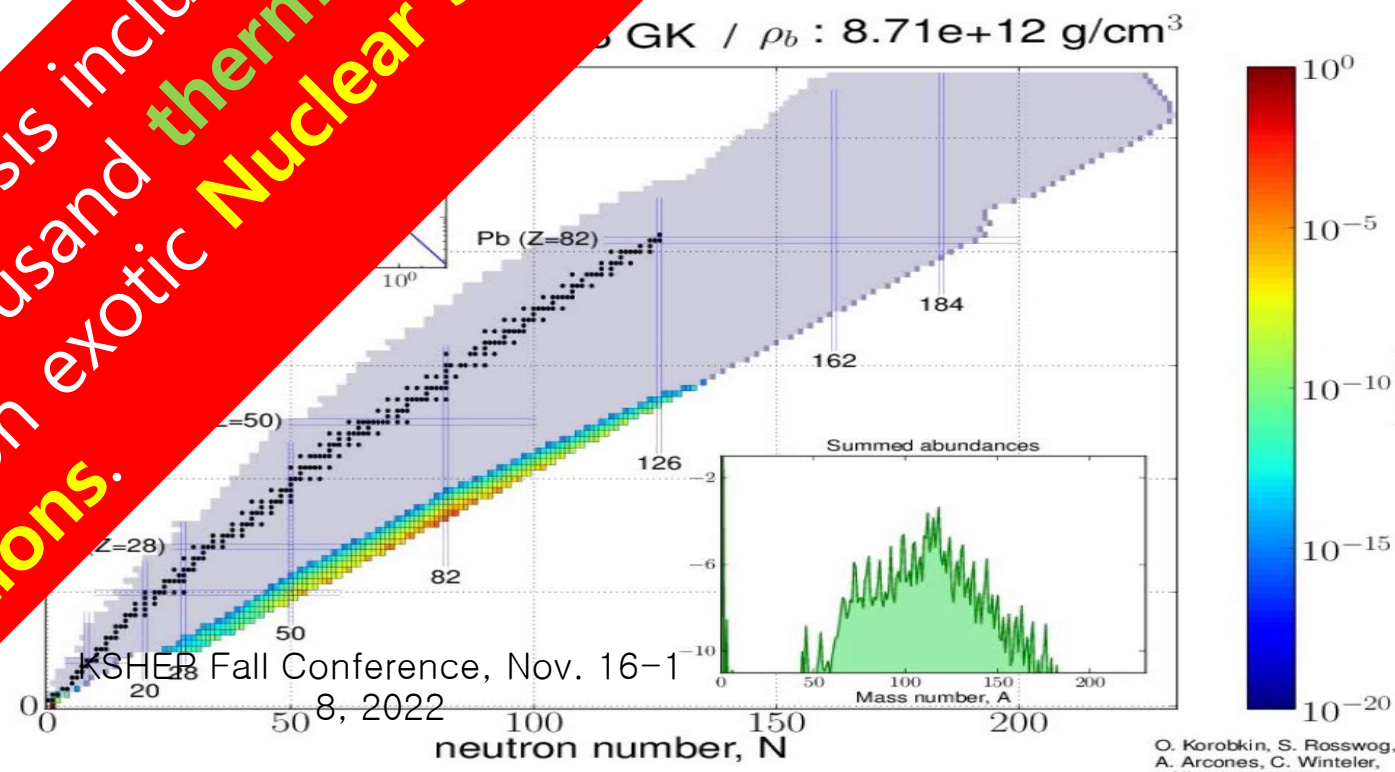
# Motivation

# Rapid process in SN and NS merge for nucleosynthesis

## SN



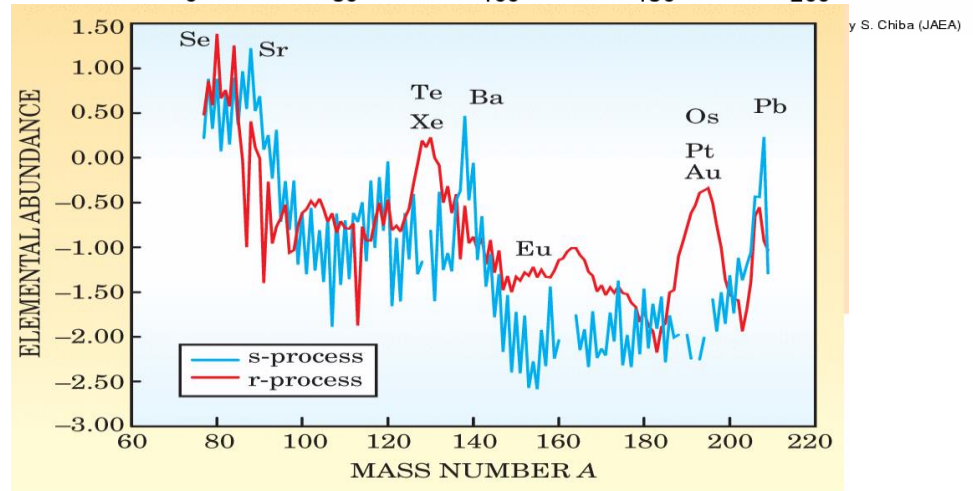
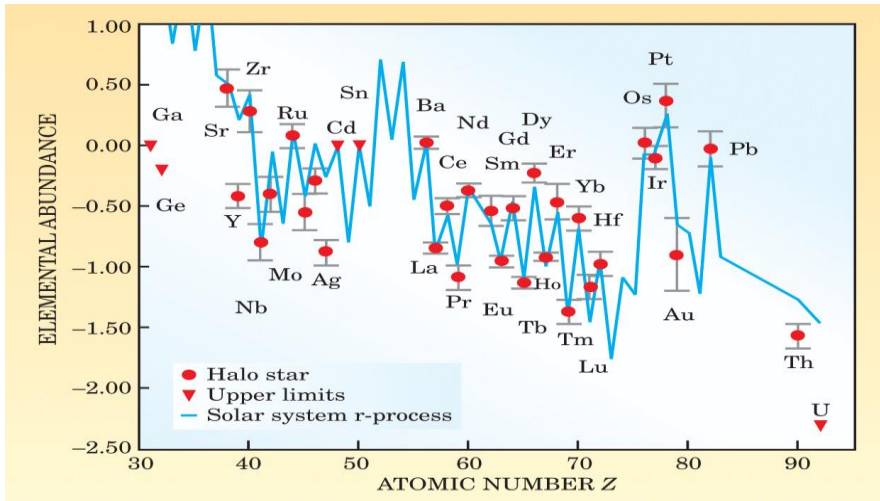
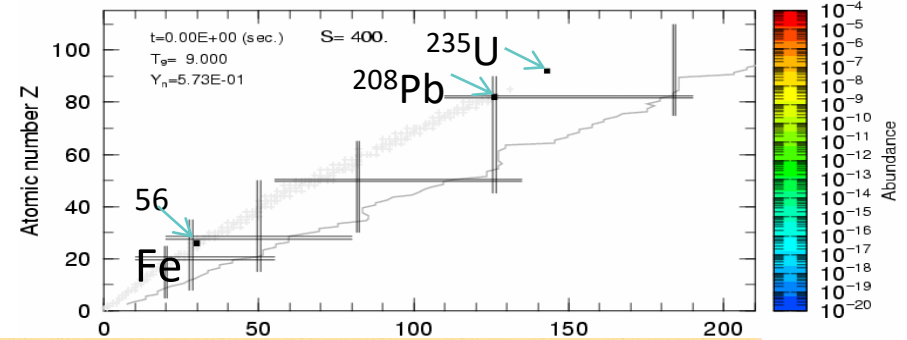
## NS Merger



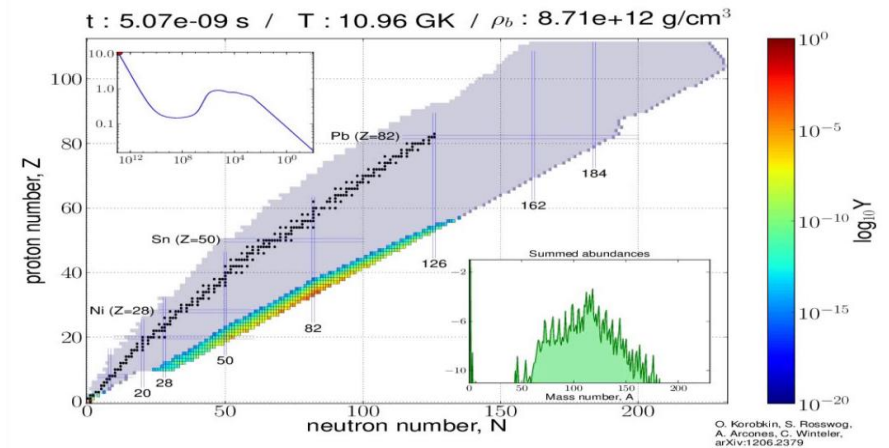
Nucleosynthesis includes  
tens of thousand thermal nuclear reactions and  
hinges on exotic Nuclear Structure and  
Reactions.

KSHER Fall Conference, Nov. 16-18, 2022

## r-process in SN

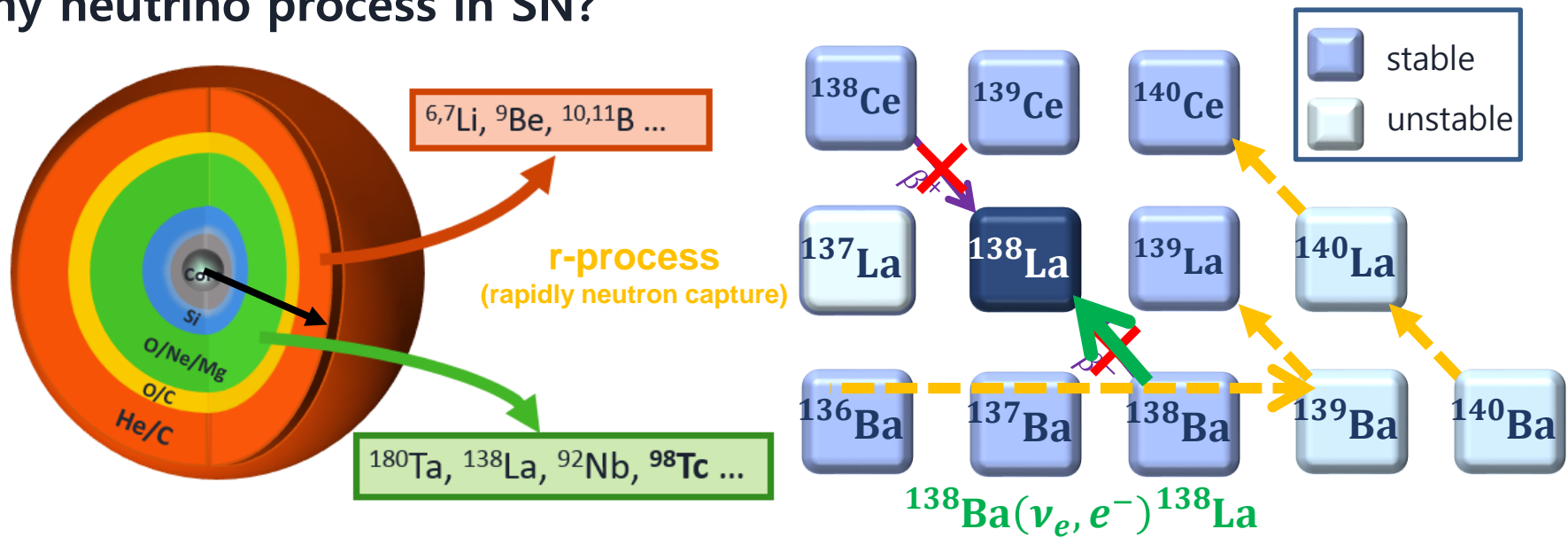


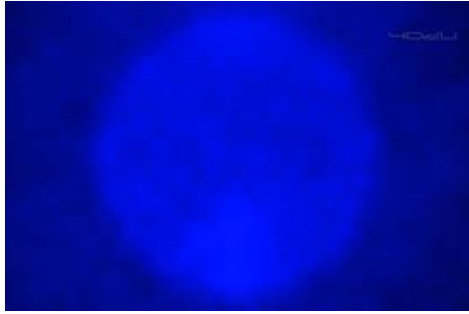
## Nucleosynthesis in Neutron Star Merge





## Why neutrino process in SN?





r-process in SN

PHYSICAL REVIEW LETTERS 121, 102701 (2018)

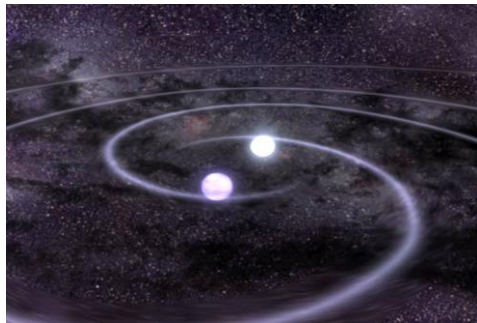
H. Ko, M.K. Cheoun et al.

Short-Lived Radioisotope  $^{98}\text{Tc}$  Synthesized by the Supernova Neutrino Process

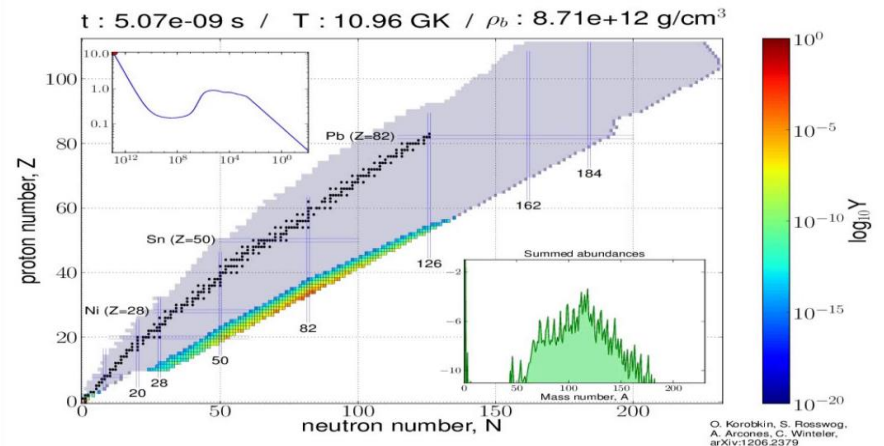
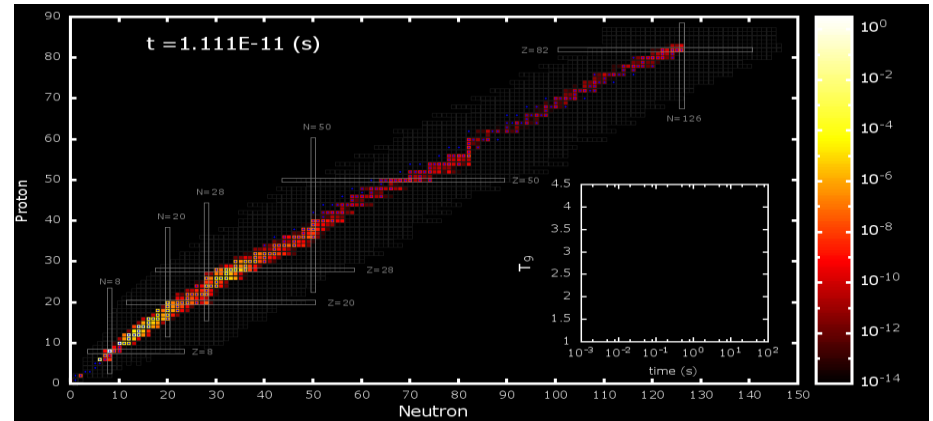
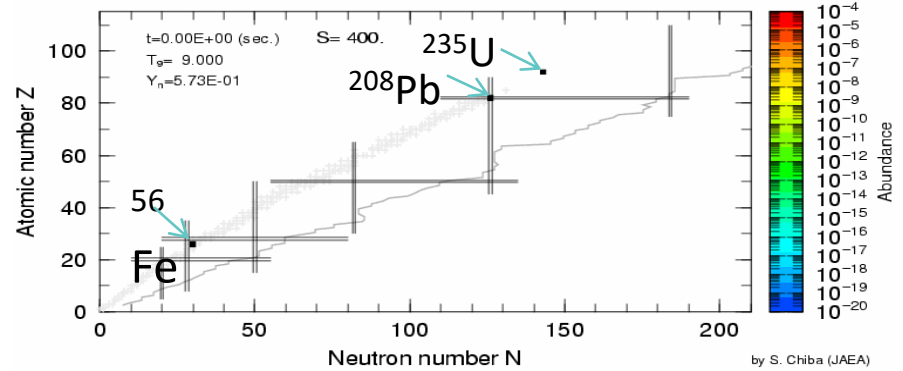
A Supernova Secret May Be Hidden Inside Meteorites

By Bill Andrews | September 4, 2018 3:50 pm

SCIENCE FOR THE CURIOUS  
**Discover**

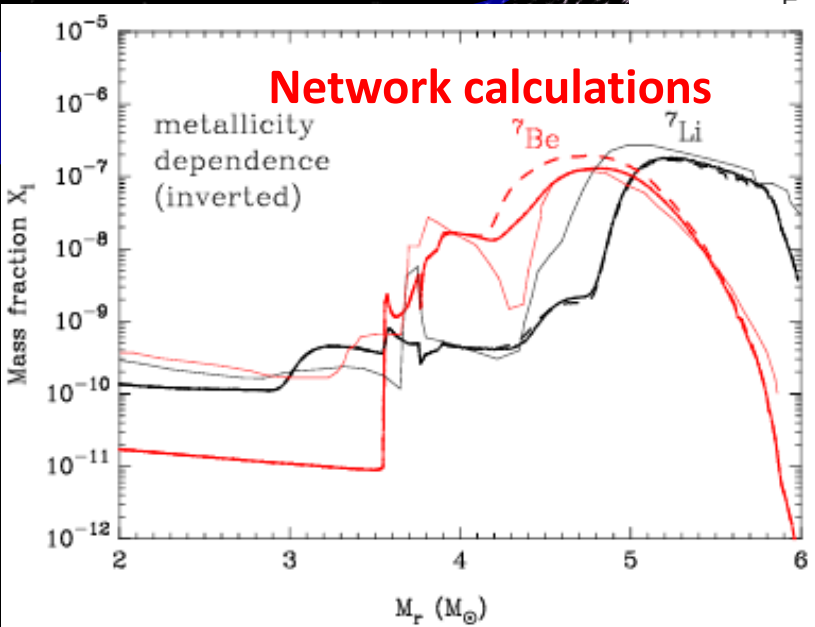


Nucleosynthesis  
in  
Neutron Star Merge



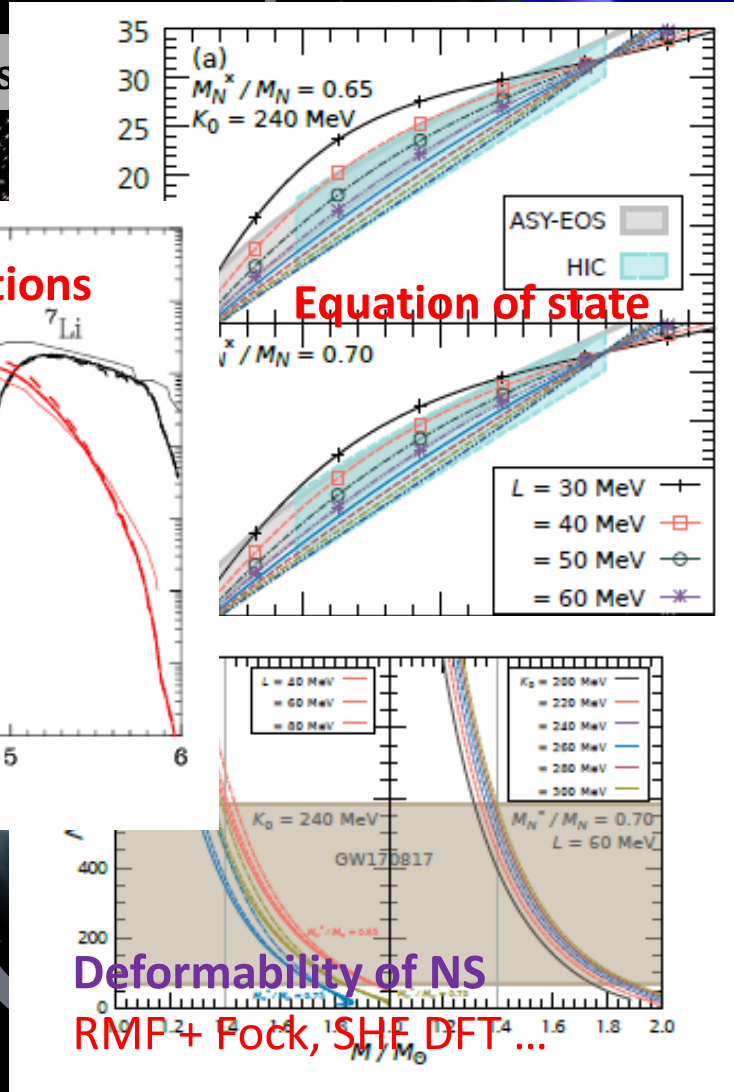
## Life cycle of stars

### 5. Nuclear Abundances



### 3. Nucleosynthesis

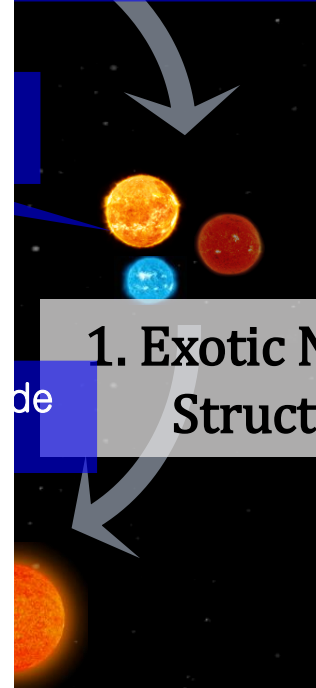
Shock wave by sudden gravitational collapse and neutrino driven wind



### Deformability of NS

RMP + Fock, SHE, DFT ...

Non-uniform distributions of Mo. Cloud become a star by the gravitation



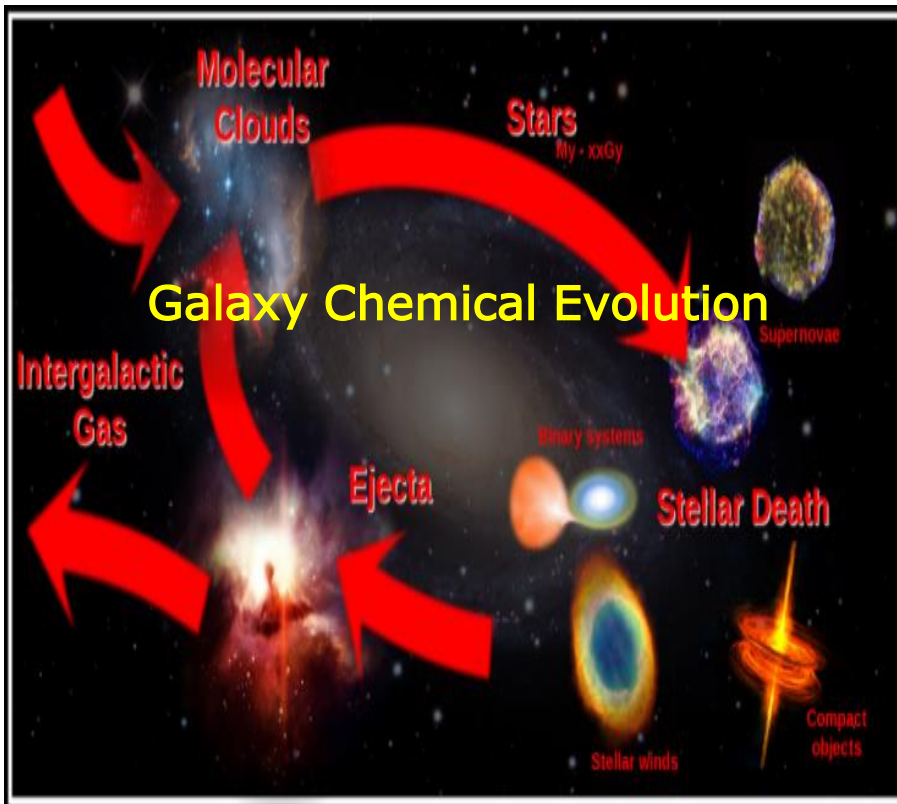
### 1. Exotic Nuclear Structure

### 2. Thermal Nuclear Reactions in the Universe



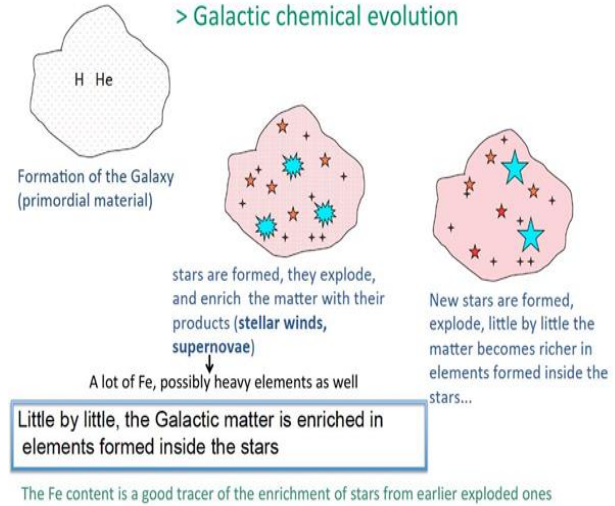
**Motivation**

**Galaxy chemical evolution for cosmological matter**



## Star Formation rate

Enrichment in elements over time ?  
> Galactic chemical evolution



The equation of chemical evolution for the element  $i$  in each galactic shell takes the following form:

$$\begin{aligned} \frac{dG_i(t)}{dt} = & -\psi(t)X_i(t) + \int_{M_L}^{M_{BM}} \psi(t - \tau_m)Q_{mi}(t - \tau_m)\phi(m)dm \\ & + A \int_{M_{BM}}^{M_{BM}} \phi(m) \left[ \int_{\mu_{min}}^{0.5} f(\mu)Q_{mi}(t - \tau_{m_2})\psi(t - \tau_{m_2})d\mu \right] dm \\ & + (1 - A) \int_{M_{BM}}^{M_{BM}} \psi(t - \tau_m)Q_{mi}(t - \tau_m)\phi(m)dm \\ & + \int_{M_{BM}}^{M_U} \psi(t - \tau_m)Q_{mi}(t - \tau_m)\phi(m)dm \\ & + \left( \frac{dG_i(t)}{dt} \right)_{infall}, \end{aligned} \quad (1)$$

where  $G_i(t) = M_{gas}(t) X_i(t)$  is the fractional mass of the element  $i$  at the time  $t$  in the ISM. The quantity  $X_i(t)$  is defined as the abundance by mass of the element  $i$ . By definition  $\sum_i X_i = 1$ .

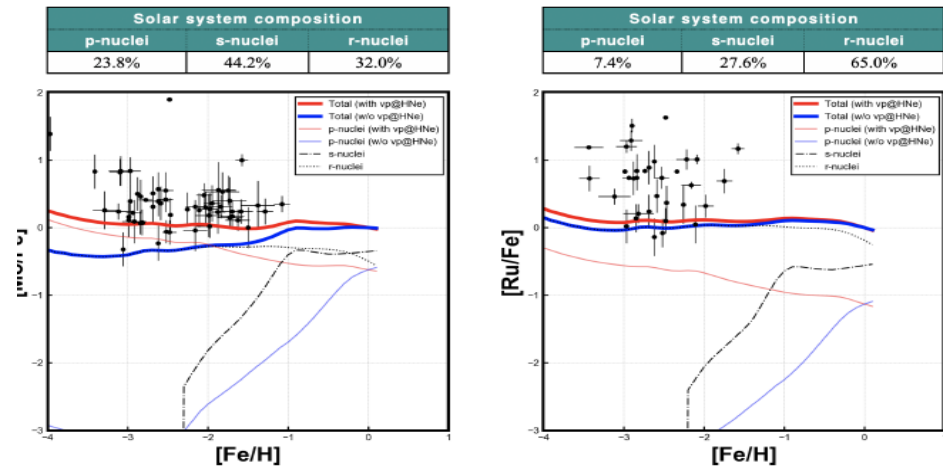


Figure 3. elemental (a) [Mo/Fe], (b) [Ru/Fe] vs. [Fe/H] //TODO

## IR spectra from ISM & Exotic Molecules

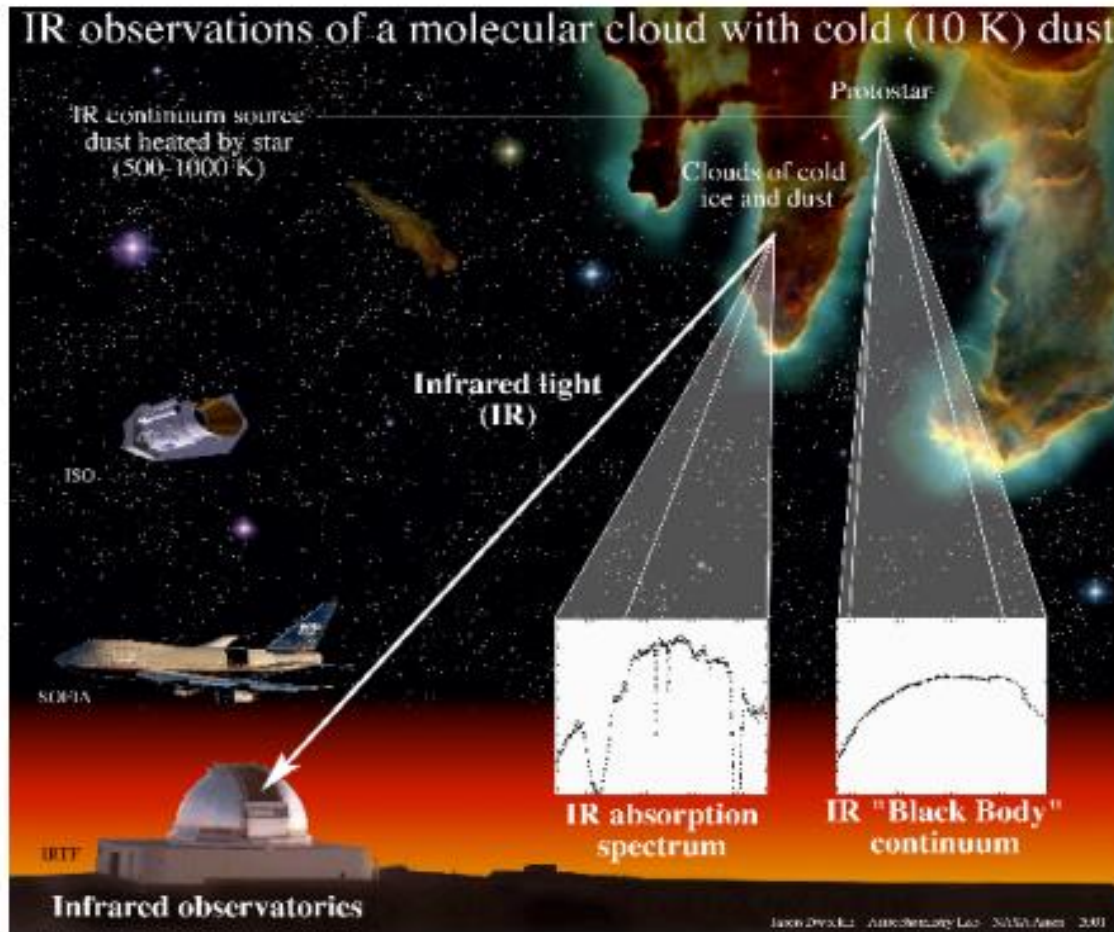
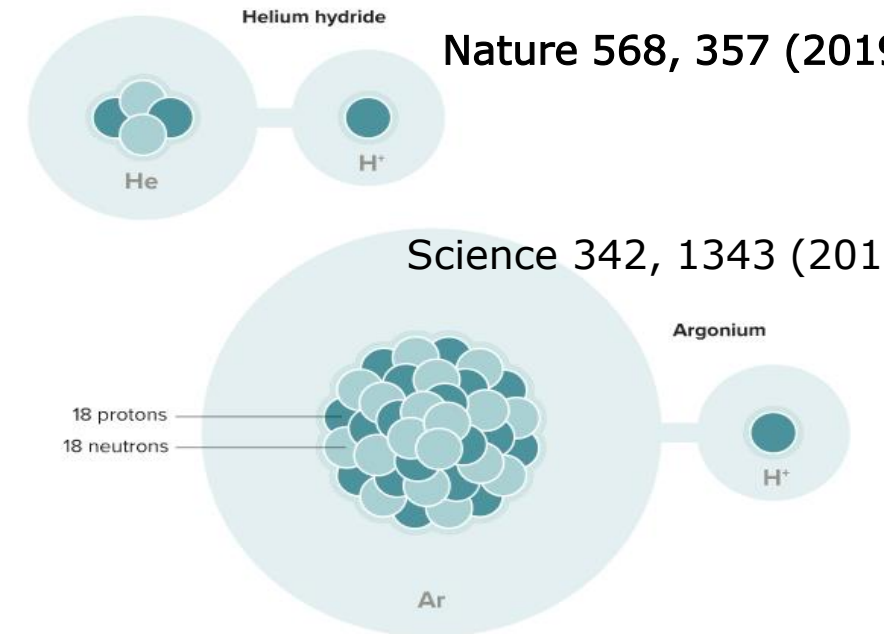


Plate 3. An illustration of how one measures the IR absorption spectra of interstellar clouds. A background star (protostar) serves as the IR source, the cloud is the sample, and the telescope gathers the light and sends it to a monochromator or spectrometer. The advent of airborne IR telescopes in the 70's and orbiting telescopes in the 80's made it possible to avoid IR absorptions by atmospheric H<sub>2</sub>O, CO<sub>2</sub>, and so on, opening a new window into the Cosmos. Figure courtesy of Dr. Jason Dworkin.

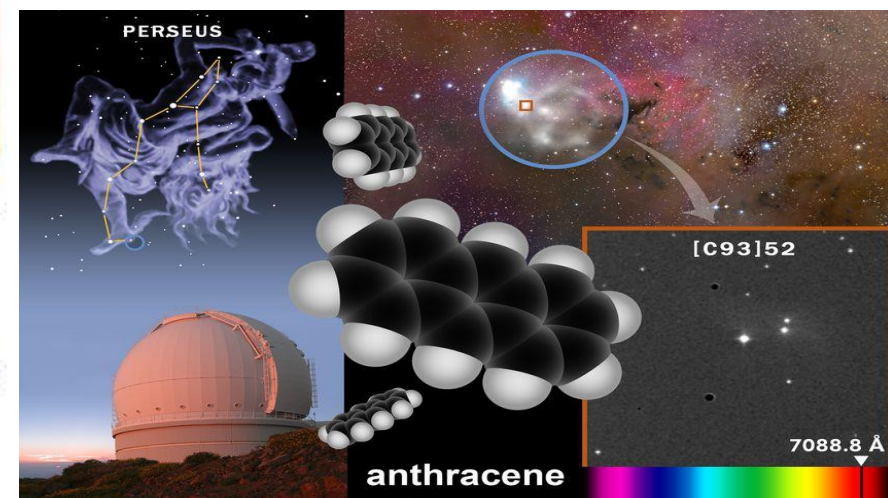
## Space molecules



SOURCE: REPORTING BY K. CROSWELL

KNOWABLE MAGAZINE

Helium hydride and argonium are the two noble gas molecules astronomers have found in space.





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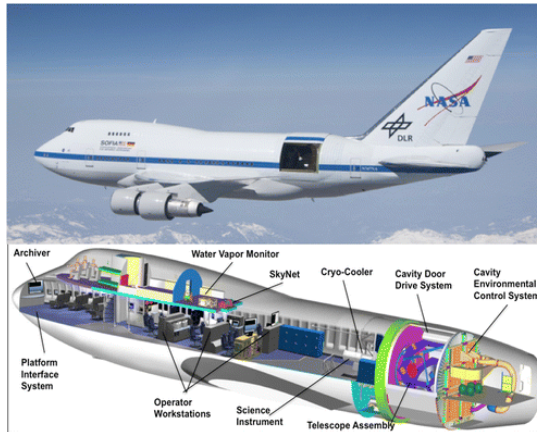
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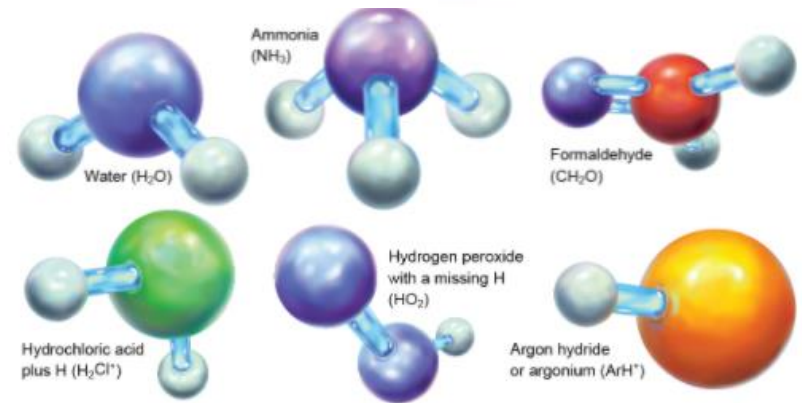
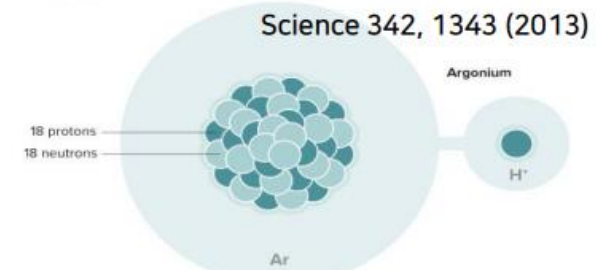
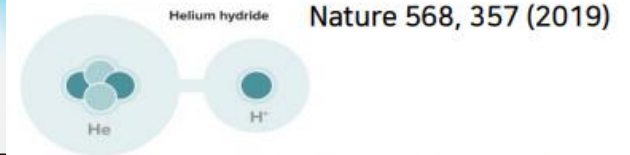
Credit: Mondolithic Studios

## The First Molecule in the Universe

Scientists have identified mystery molecules in space and the compounds started chemistry in the cosmos

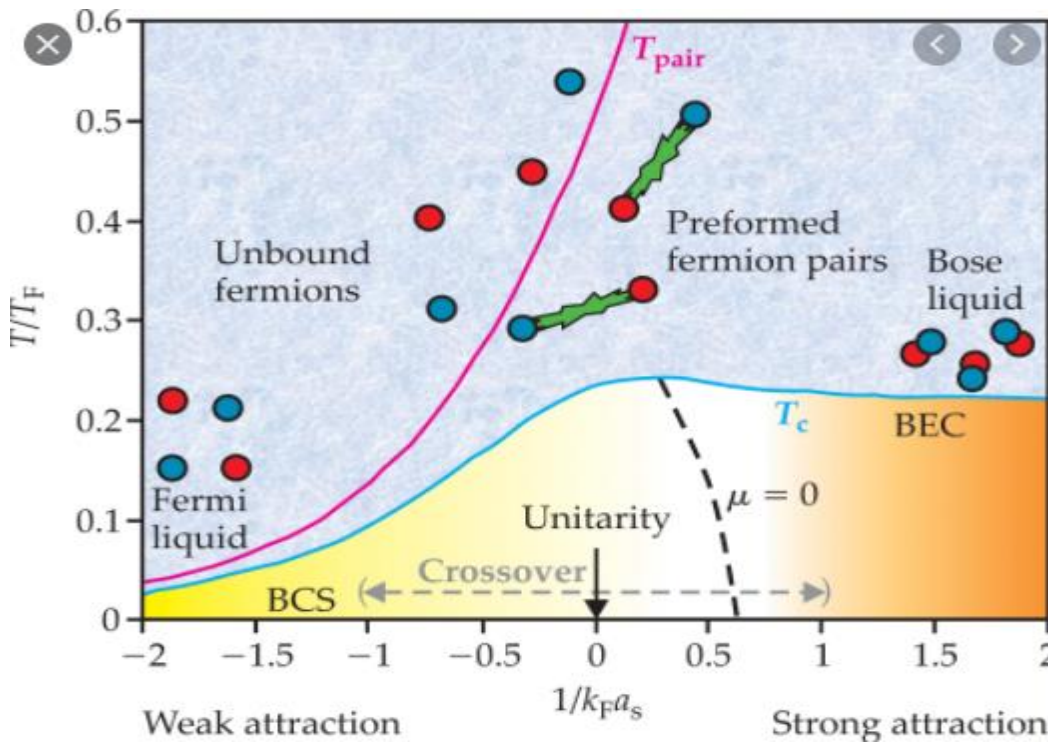
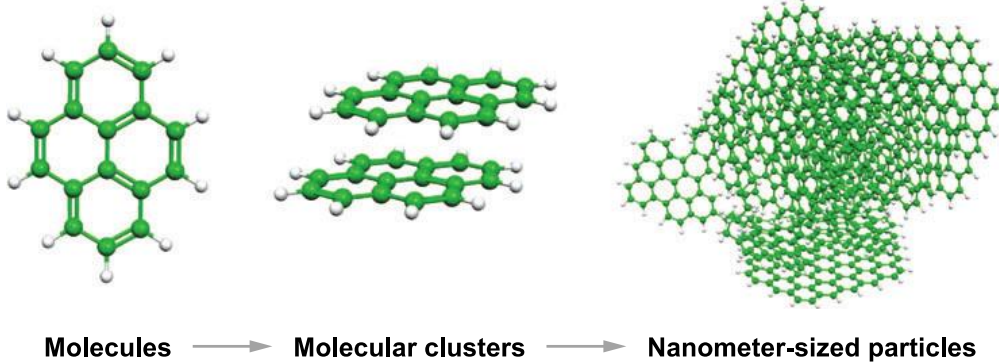


### Space molecules

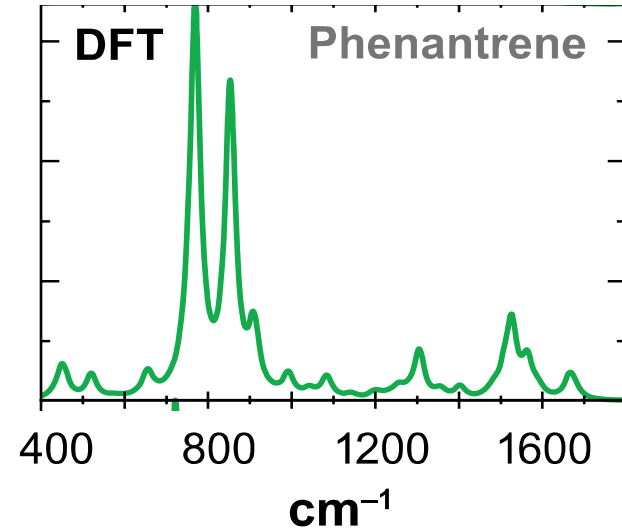


Credit: Elena Hartley

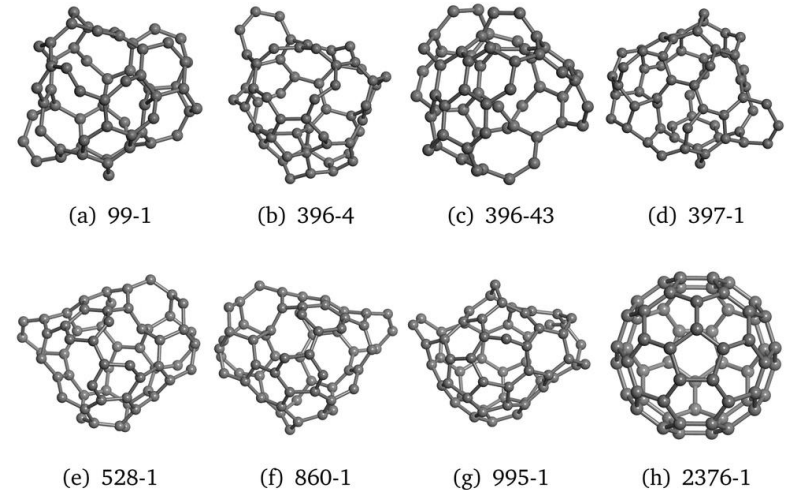
- 성간 분자 물질에 의한 적외선 스펙트럼 연구
- 크기에 따른 hydrocarbon 기반의 성간 우주물질



- 밀도 범함수 이론을 이용한 방출 스펙트럼 계산



- 탄소 기반의 국소적 안정 상태의 원자구조

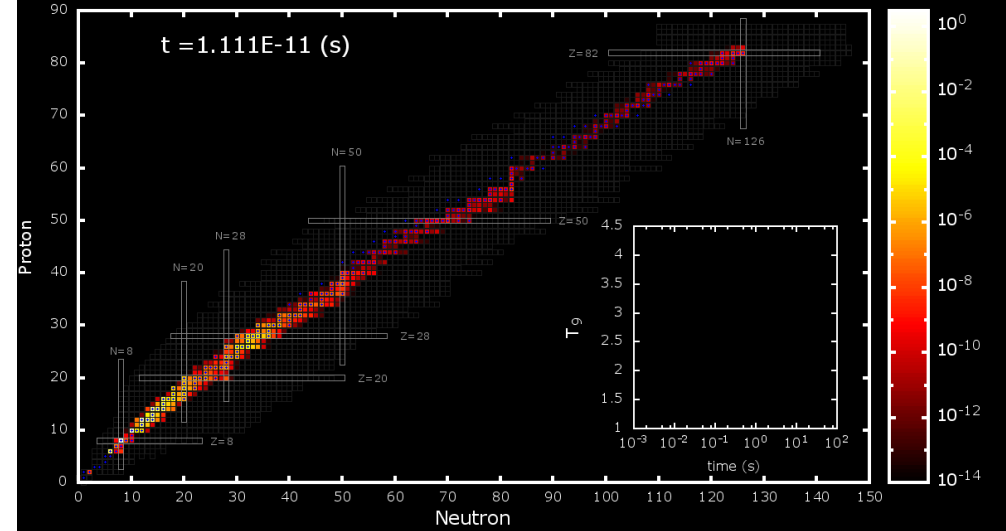


# Neutrino process Neut. Ham. for neutrino density propagation

JINA REALIB

Modified (n,g) Reactions

QRPA & Branching Ratios



# Numerical results for elements abundances

1987 SN model

Pre-supernova Model

Hydrodynamics Model: HCK18, KCK19

Modified Neutrino Flux by Self-interaction : w/ and w.o/

Neutrino Luminosity : EQ and NEQ

Mass Hierarchy : NH and IH

. 16-18,

## Network calculation for nucleosynthesis

$$\frac{dN_j}{dt} = N_i \lambda_{i,j} - N_j \lambda_{j,h} + \dots \rightarrow \frac{dY_j}{dt} = Y_i \lambda_{i,j} - Y_j \lambda_{j,h} + \dots$$

$$Y_j = \frac{N_j}{\rho N_A}$$

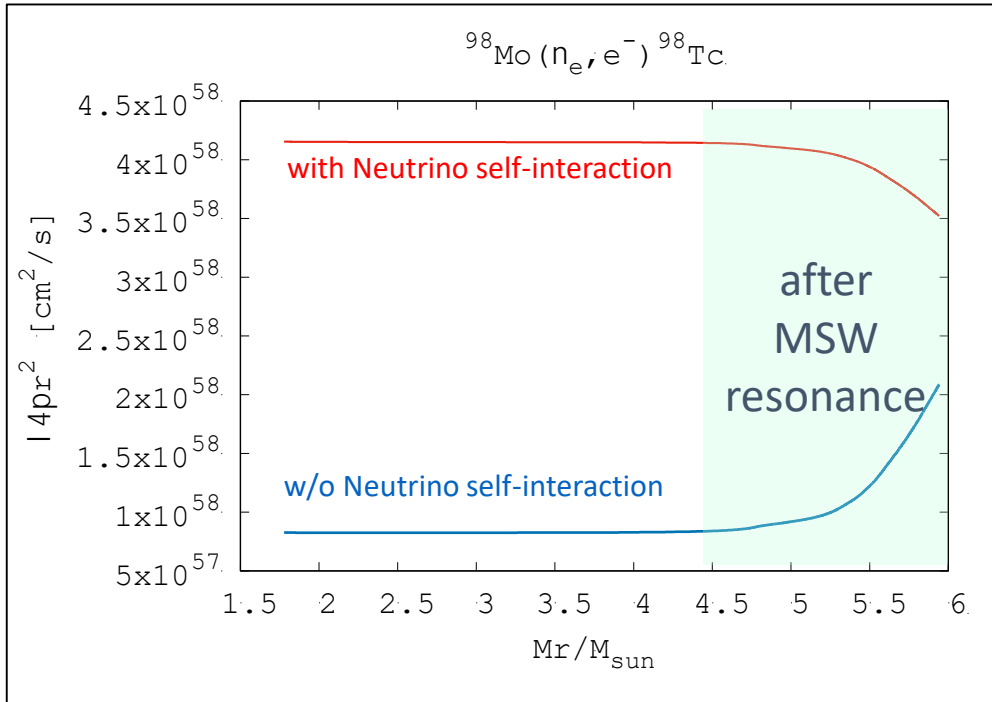
JINA REACLIB & Los Alamos (n,g) Data !

Part for neutrino reaction rates

Kyushu-Tokyo Progenitor Model !

$$\lambda_{\nu\alpha}(r) = \sigma \phi = \int_0^\infty \sum_{\alpha=e,\mu,\tau} \frac{d\phi_{\nu\alpha}}{d\epsilon_\nu} Br(\epsilon) \sigma_{\nu\alpha}(\epsilon_\nu) d\epsilon_\nu$$

Example: (EQ,IH)



## Cross section data using QRPA

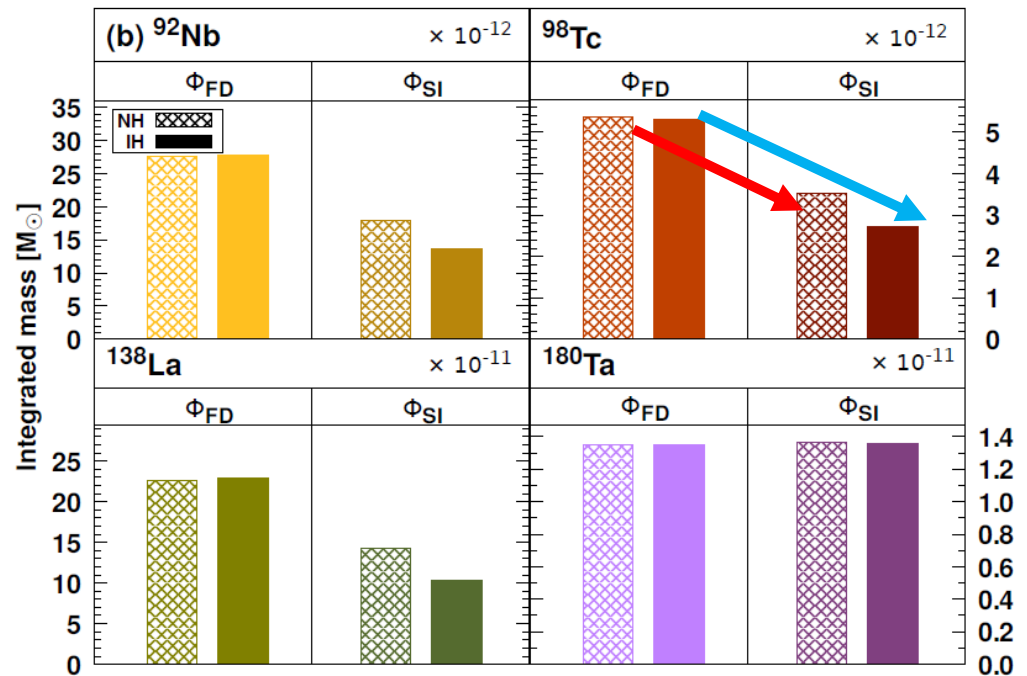
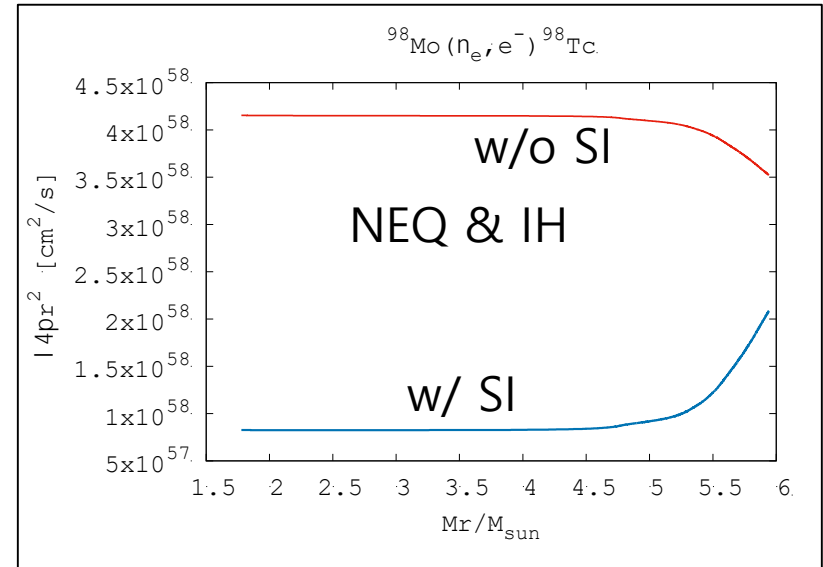
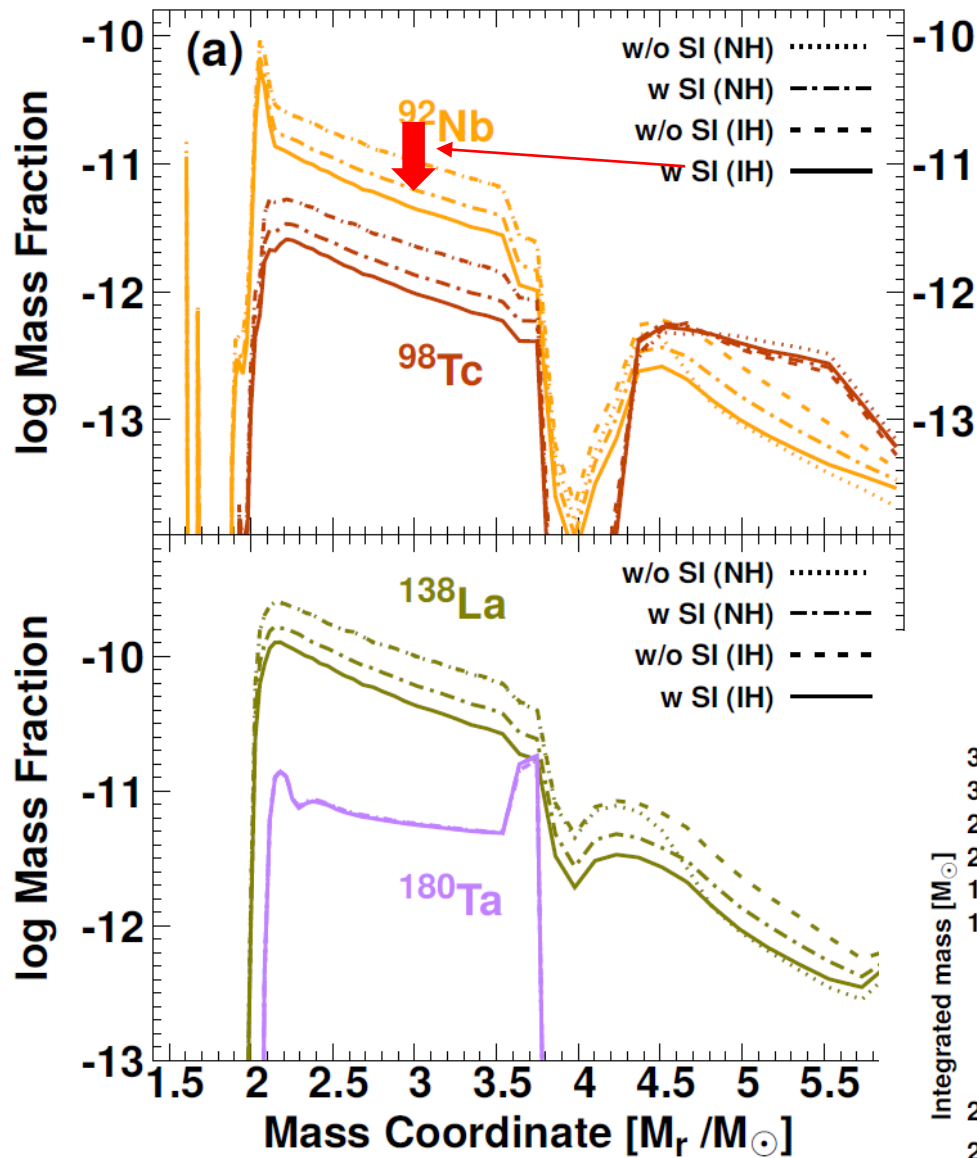
TABLE I. Averaged cross sections in units of  $10^{-42} \text{ cm}^2$  for  ${}^{98}\text{Mo}$  via CC and  ${}^{99}\text{Ru}$  via NC, and  ${}^{92}\text{Zr}$  via CC and  ${}^{93}\text{Nb}$  via NC with particle emission. Neutrino temperatures are taken from [4] and  $\langle E_k \rangle$  is calculated from  $\langle E_k \rangle / T \sim 3.1514 + 0.1250\alpha$  with  $\alpha = 0$  [31,42].

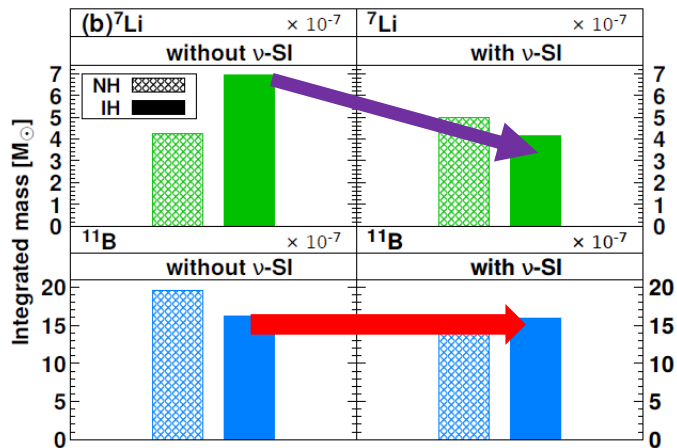
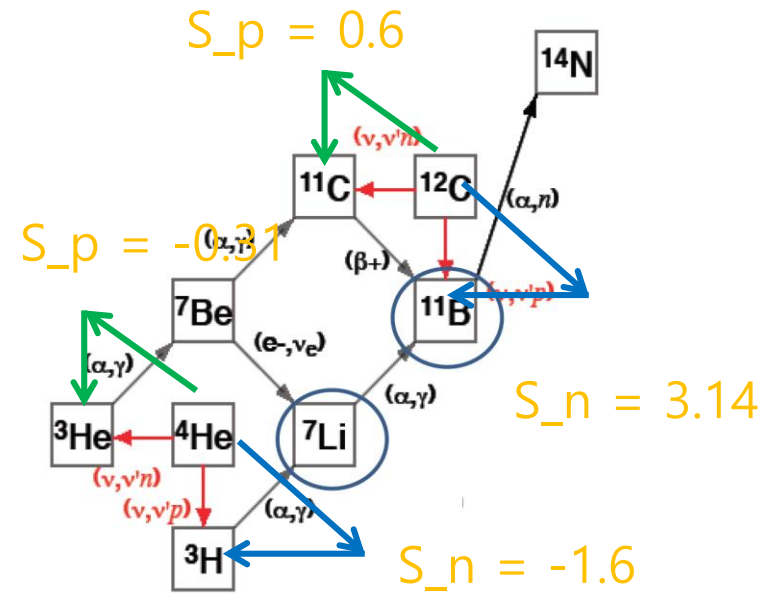
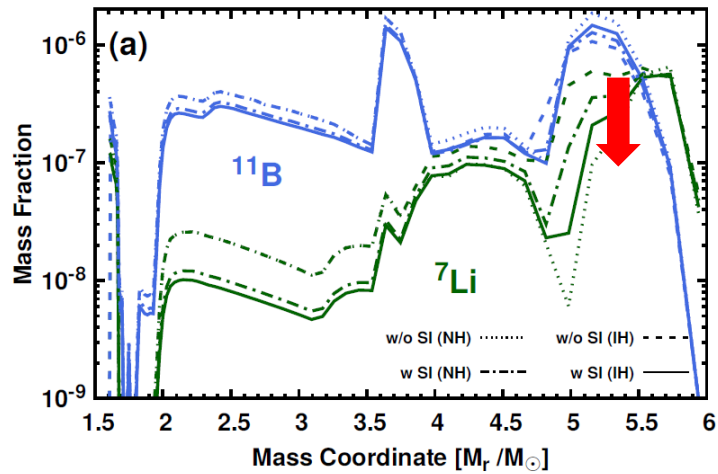
Reactions	$\langle E_k \rangle$ [MeV]	$T$ [MeV]	$\langle \sigma \rangle$
${}^{98}\text{Mo}(\nu_e, e^-){}^{98}\text{Tc}$	10.08	3.2	7.77
${}^{98}\text{Mo}(\nu_e, e^- p){}^{97}\text{Mo}$	10.08	3.2	1.90
${}^{98}\text{Mo}(\nu_e, e^- n){}^{97}\text{Tc}$	10.08	3.2	0.09
${}^{99}\text{Ru}(\bar{\nu}_\mu, \bar{\nu}'_\mu){}^{99}\text{Ru}$	18.90	6.0	78.5
${}^{99}\text{Ru}(\bar{\nu}_\mu, \bar{\nu}'_n){}^{98}\text{Ru}$	18.90	6.0	14.6
${}^{99}\text{Ru}(\bar{\nu}_\mu, \bar{\nu}'_p){}^{98}\text{Tc}$	18.90	6.0	1.70
${}^{99}\text{Ru}(\bar{\nu}_e, \bar{\nu}'_e){}^{99}\text{Ru}$	15.75	5.0	52.1
${}^{99}\text{Ru}(\bar{\nu}_e, \bar{\nu}'_n){}^{98}\text{Ru}$	15.75	5.0	10.5
${}^{99}\text{Ru}(\bar{\nu}_e, \bar{\nu}'_p){}^{98}\text{Tc}$	15.75	5.0	0.92
${}^{92}\text{Zr}(\nu_e, e^-){}^{92}\text{Nb}$	10.08	3.2	8.92

In MSW region, energetic e-neutrino is increased by the P(x-e) neutrino resonance ( w/o SI ).

But it is a bit decreased with the decrease of X-neutrino by the SI.







For  ${}^7\text{Li}$ , the main reactions are both  $e^-$ - and anti- $e^-$ - neutrino reactions via CC which are larger than NC. And  $e^-$ -CC reactions through  ${}^3\text{He}$  and  ${}^7\text{Be}$  from  ${}^4\text{He}$  are larger than anti- $e^-$  due to MSW.  $\Rightarrow$  Sensitive on the nu-SI.

But for  ${}^{11}\text{B}$  both electro- and antielectron-neutrinos via CC and NC work.  $\Rightarrow$  Insensitive to the nu-SI.

Hydrodynamics : HKC18 and KCK19

Luminosity : EQ and NEQ

Neutrino Self Interaction : FD and SI

Mass Hierarchy : NH and IH

**Table 4.** Integrated masses of the nuclei after 50 s in the mass range,  $M_r = 1.6-6 (M_\odot)$ . We used two hydrodynamics models (HKC18 and KCK19), two luminosity models (EQ and NEQ) and two cases without the  $\nu$ -SI (FD) and with the  $\nu$ -SI (SI) for the NH and IH case, by which the results for twelve different cases are tabulated. The last two results are quoted from our previous results. See texts for the details.

	Mass Hierarchy	${}^7\text{Li}$	${}^7\text{Be}$	${}^{11}\text{B}$	${}^{11}\text{C}$	${}^{92}\text{Nb}$	${}^{98}\text{Tc}$	${}^{138}\text{La}$	${}^{180}\text{Ta}$	Yield ratio	PF ratio
		$(10^{-7} M_\odot)$				$(10^{-12} M_\odot)$	$(10^{-11} M_\odot)$		$N({}^7\text{Li})/N({}^{11}\text{B})$	${}^{138}\text{La}/{}^{11}\text{B}$	
FD EQ	NH	1.256	4.953	5.576	2.048	4.903	1.048	3.395	0.845	1.280	0.1288
(HKC18)	IH	1.496	1.461	7.141	1.218	4.760	1.112	3.267	0.843	0.556	0.1130
FD EQ	NH	0.861	2.428	2.480	2.139	4.551	1.180	3.760	1.016	1.119	0.2354
(KCK19)	IH	1.017	0.936	3.099	0.883	4.226	1.218	3.436	1.012	0.771	0.2495
FD EQ Shock	NH	0.861	1.904	2.546	1.701	4.973	1.271	4.164	1.017	1.023	0.2835
(KCK19)	IH	0.949	1.027	2.922	0.937	4.271	1.215	3.485	1.012	0.805	0.2611
SI EQ <sup>a</sup>	NH	0.861	2.428	2.480	2.139	4.551	1.180	3.760	1.016	1.119	0.2354
(KCK19)	IH	0.920	2.057	2.852	3.874	15.07	3.259	13.58	1.052	0.695	0.5838
SI NEQ	NH	1.132	1.601	4.276	4.920	16.44	3.559	15.19	1.295	0.467	0.4776
(KCK19)	IH	1.261	1.206	4.623	4.283	12.29	2.854	11.31	1.281	0.435	0.3672
FD NEQ	NH	1.483	0.841	5.407	5.258	25.44	5.367	23.14	1.323	0.342	0.6274
(KCK19)	IH	0.959	2.303	3.946	6.566	26.15	5.302	23.94	1.331	0.488	0.6585
SI NEQ Ko et al. (2020)	NH	1.643	3.347	9.332	6.138	17.92	3.511	14.29	1.363	0.507	0.2671
(HKC18)	IH	1.792	2.372	10.33	5.524	13.59	2.720	10.41	1.358	0.413	0.1899
FD NEQ Ko et al. (2020)	NH	2.400	1.860	12.46	7.080	27.56	5.361	22.62	1.349	0.343	0.335
(HKC18)	IH	1.640	5.270	8.382	7.804	27.83	5.318	22.94	1.353	0.671	0.410

<sup>a</sup>Same as FD EQ (KCK19) NH result

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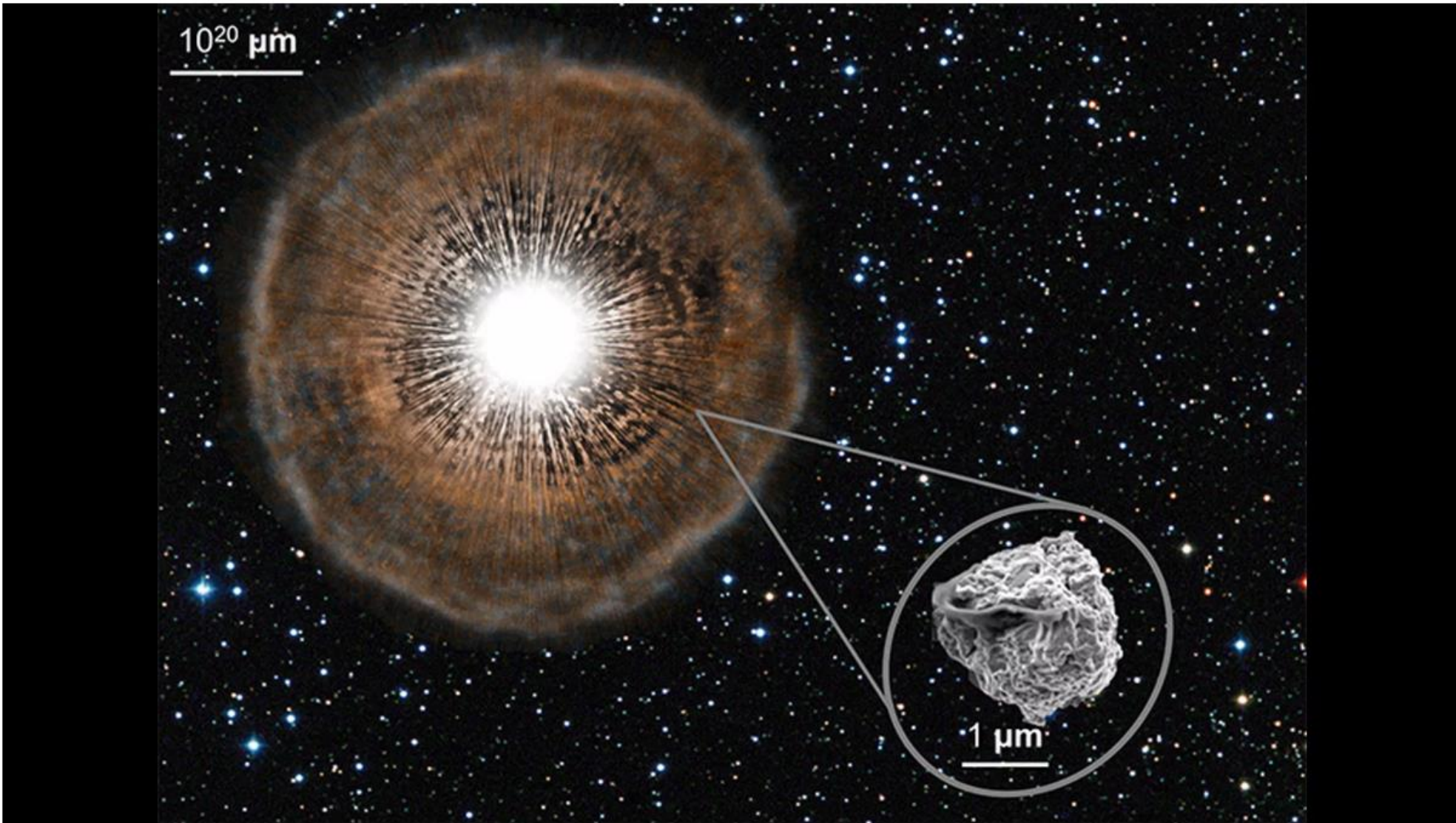


## Neutrino Process in Core-collapse Supernovae with Neutrino Self-interaction and MSW Effects

Heamin Ko<sup>1</sup>, Myung-Ki Cheoun<sup>1</sup>, Eunja Ha<sup>1</sup>, Motohiko Kusakabe<sup>2</sup>, Takehito Hayakawa<sup>3</sup>, Hirokazu Sasaki<sup>4</sup>, Toshitaka Kajino<sup>2,4</sup>, Masa-aki Hashimoto<sup>5</sup>, Masaomi Ono<sup>6</sup>, Mark D. Utsunomiya<sup>7</sup>, Satoshi Chiba<sup>7</sup>, Ko Nakamura<sup>8</sup>



# Meteorite Analysis : Analysis of SiC in stardust grain (Pre-solar grain) from Murchinson meteorite



An electron microscope image of a micron-sized silicon carbide stardust grain extracted from the Murchinson meteorite. (Image credit: NASA, Nan Liu and Andrew Davis)

# Mass Fraction ratio of ${}^7\text{Li}/{}^{11}\text{B}$ and PF ratio of ${}^{138}\text{La}/{}^{11}\text{B}$

- The yield ratio of  $[{}^7\text{Li}/{}^{11}\text{B}]$

$${}^7\text{Li}/{}^{11}\text{B} = -0.31 \pm 0.42 \quad < \mathbf{0.53} \text{ (2 sigma)}$$

Spectra	FD	+SI	FD	+SI
Mass Hierarchy	IH	IH	NH	NH
Yield Ratio	0.671(0.488)	0.413(0.435)	0.343(0.342)	0.507(0.467)

- The production factor ratio of  $[{}^{138}\text{La}/{}^{11}\text{B}]$

$$> \mathbf{0.41} \text{ (lower limit)}$$

$$\text{PF}[A] = X_A / X_{A\odot} \text{ with } X_A \text{ the mass fraction of } A$$

Spectra	FD	+SI	FD	+SI
Mass Hierarchy	IH	IH	NH	NH
PF ratio	0.410(0.6585)	0.1899(0.3672)	0.335(0.6274)	0.2671(0.4776)

# NH is favored !!!

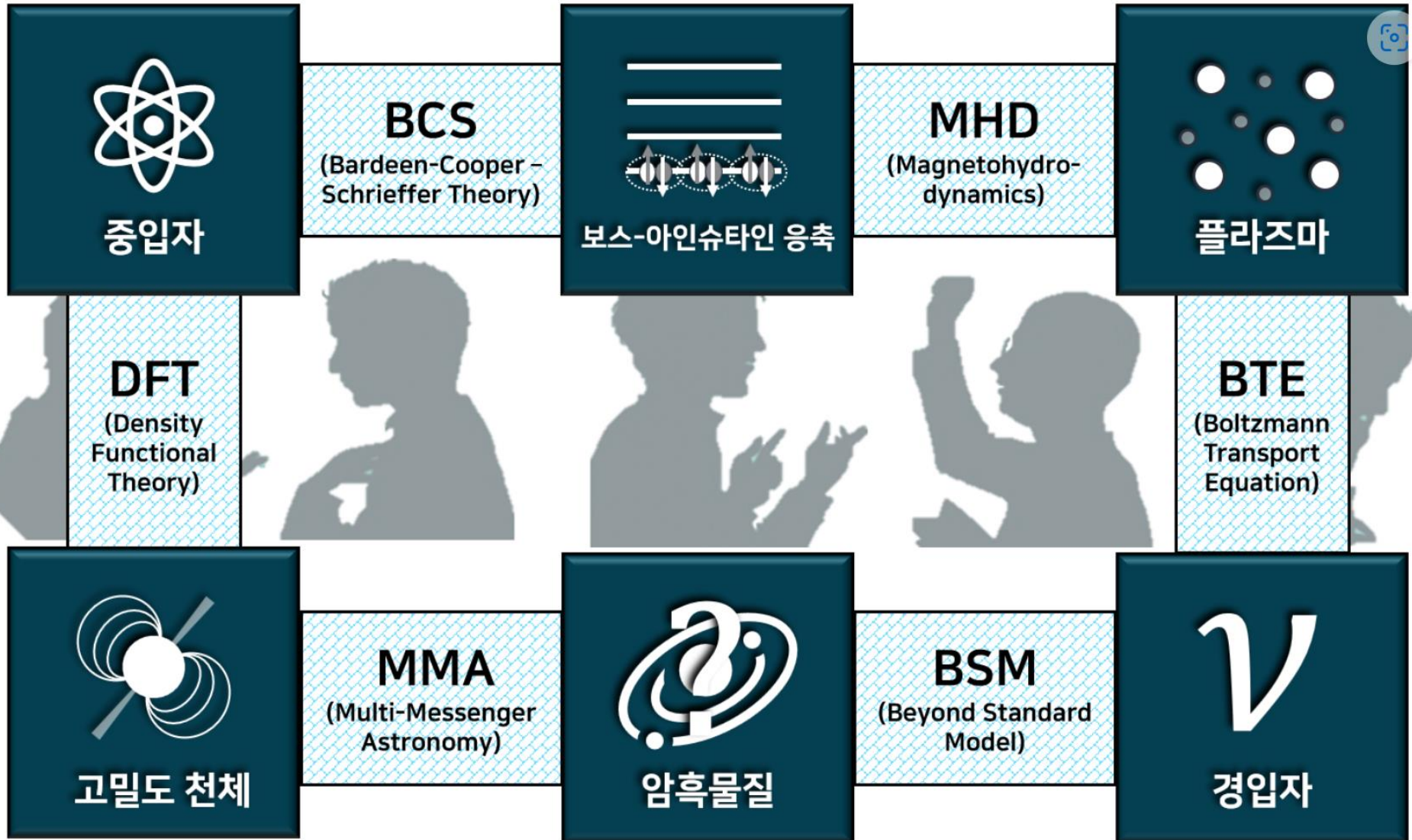
However, is this the last  
story ?

but

the least ??? Other effects ?

Sterile Neutrino,  
Magnetic field,  
Polarized Electron Density  
Neutrino Fast Oscillation..





# Member of OMEG (Origin of Matter and Evolution of Galaxies)



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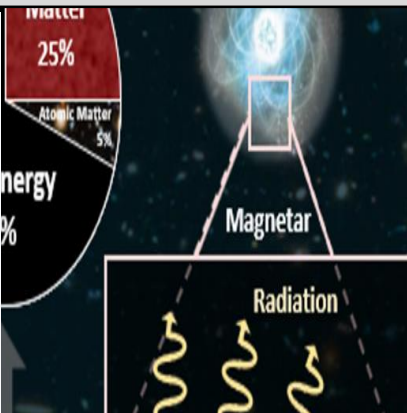
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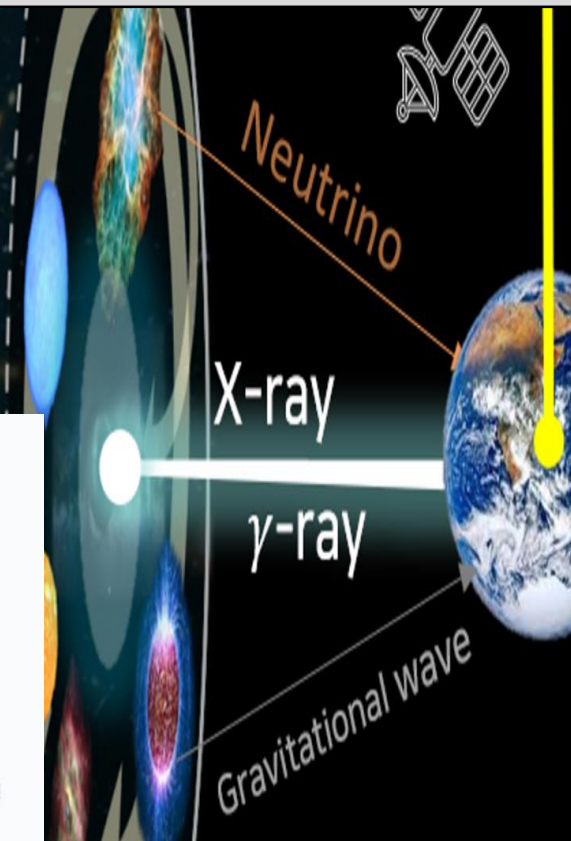
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T. Maruyama (Nihon), H. Sagawa (RIKEN), K. Hagino (Kyoto), E. Hiyama (Tohoku), T. Kawano (BNL)...




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 T. Maruyama (Nihon), H. Suga (RIKEN), K. Hagino (Kyoto), E. Hiyama (Tohoku), T. Kawano (BNL)...

대단히 감사합니다 !!



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