

How to detect ULDM

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Outline

1. ULDM (Fuzzy DM)
2. Gravitational detection of FDM
3. non-gravitational detection and SIULDM

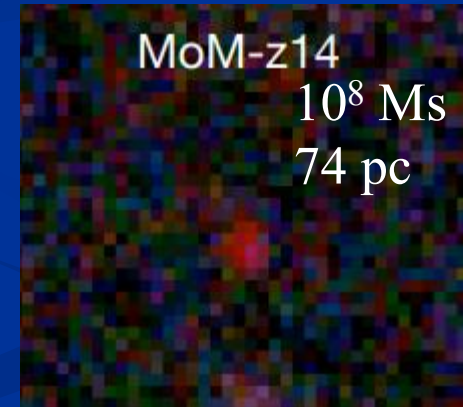
Challenges for Λ CDM

- Λ CDM was very successful but is slowly becoming non-standard?

cf) 2105.05208

1. **Small scale crisis** (on galactic scale and below) :
predicts too many small structures not observed
2. **Hubble tension $\sim 5\sigma$** :
 H_0 mismatch between CMB and SN
3. **S_8 tension $\sim 2-3\sigma$** : $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$, $z < 2$, Mpc
matter density fluctuation amplitude mismatch
between CMB and WL & Cluster
4. **DE is not Λ ? $\sim 4\sigma$** (ex, AP test, DESI, SN?)
5. **Too early SMBHs ($z \sim 11$) and galaxies ($z \sim 14, 32?$)**
6. Galaxy cluster collision (collision speed & DM-star offset)
7. Li problem, Cosmic birefringence
...etc

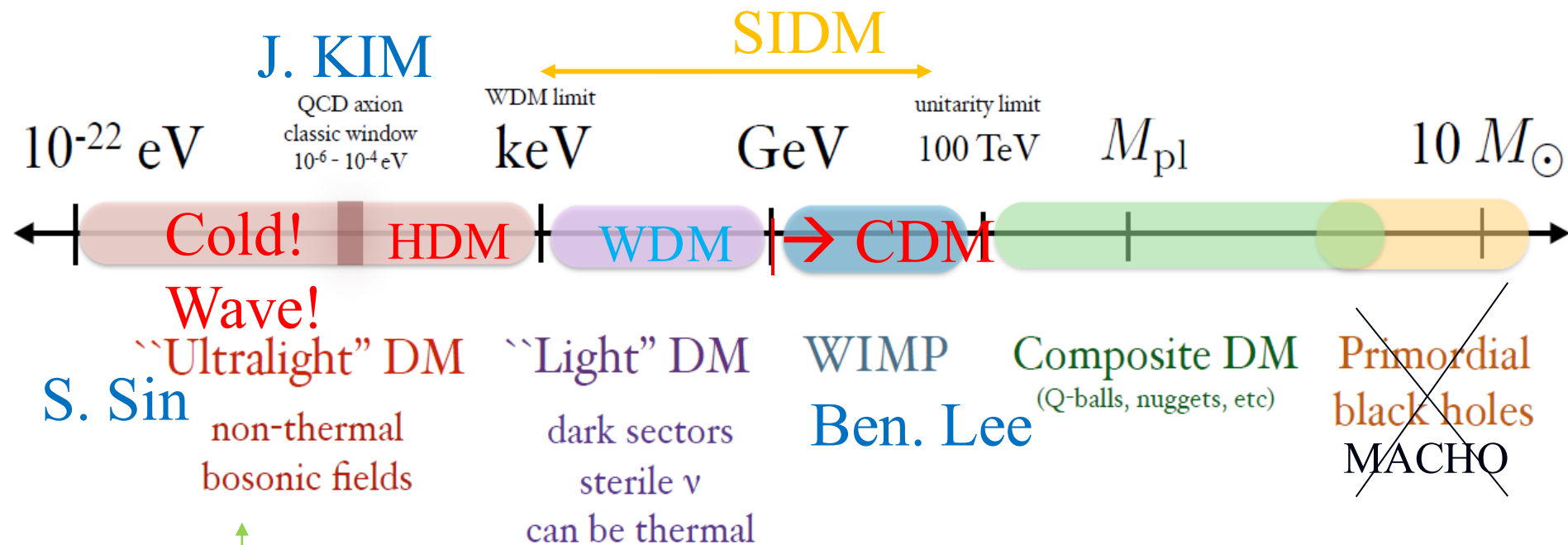
→ Any good DM model should address these tensions



Mass scale of dark matter

(not to scale)

TASI lectures by Lin arXiv:1904.07915

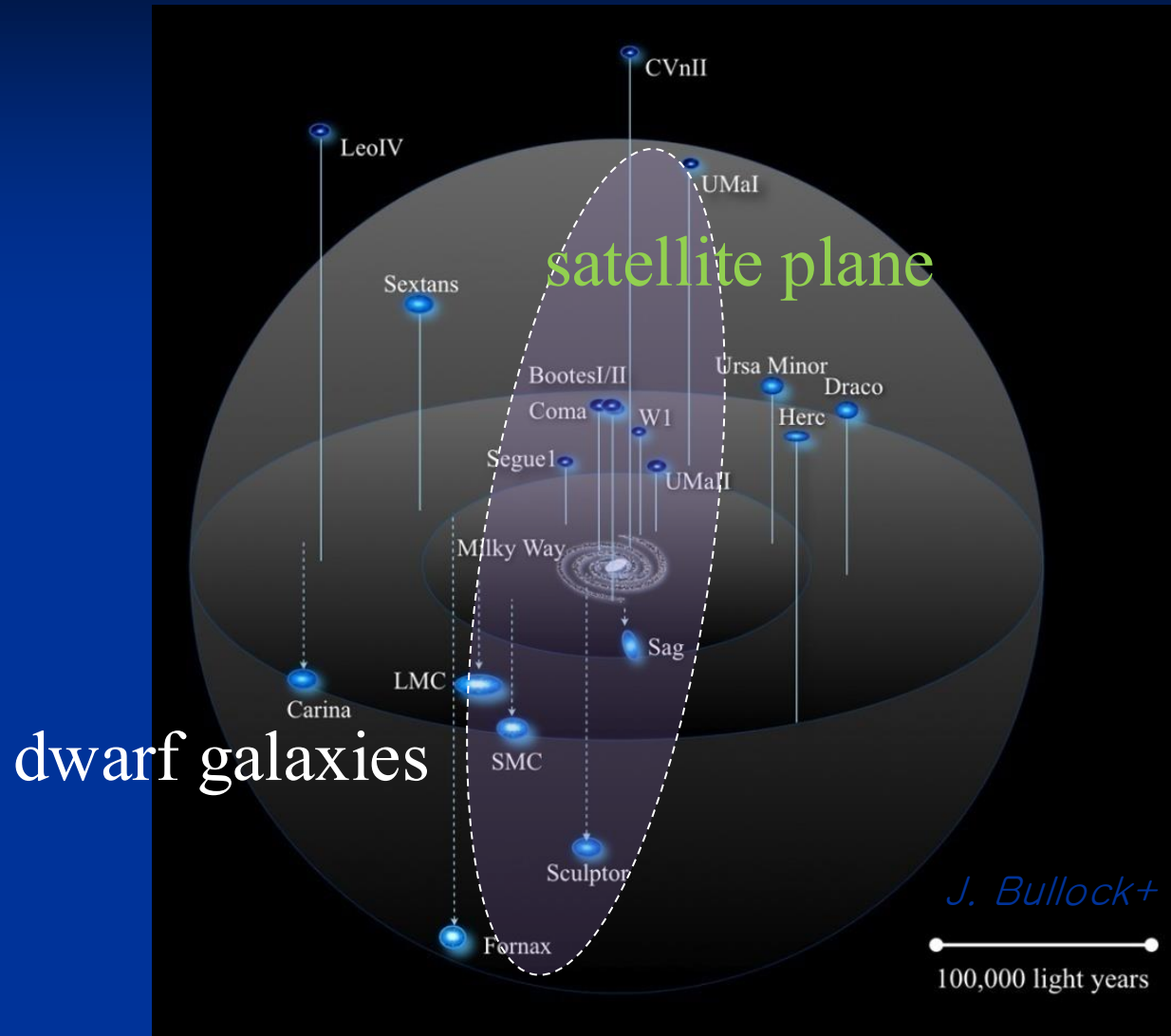


ULDM = Fuzzy, ULA, BEC, Wave, Scalar Field, ψ , Superfluid, quantum ... $\leftarrow$ PNGB, moduli, dilaton, ALP...

Compact objects in the mass range from $1.3 \times 10^{-5} M_{\odot}$ to $860 M_{\odot}$ cannot make up more than 10% of dark matter. (Mroz+ 2403.02386)

$\rightarrow$ No DM star or planet observed so far in galactic halos

Galaxies observed



Minimum mass
 $\sim 10^6 M_{\odot}$

Any good DM model should explain observed galaxies!

Small scale issues with CDM

Sales+ 2206.05295

Λ CDM Tensions with Dwarf Galaxies

No tension

Uncertain

Weak tension

Strong tension

✓ Missing satellites

$M_{\star}-M_{\text{halo}}$ relation

✓ Too big to fail

no DM model solved
Diversity of rotation curves

BTFR,
L catastrophe

✓ Core-cusp

WDM

Diversity of dwarf sizes

✓ Satellite planes

Park+ Jcap 2022

Quiescent fractions

- Key problem is how to suppress small scale structures < dwarf galaxies.
→ we need a new CDM → ULDM with $m \sim 10^{-22}$ eV might solve many of these
- Still unsolved problems seem to be related to Baryon-DM relation
- Can baryon physics + more precise numerical simulation + more observations save CDM?

Solutions to Small scale problems

- CDM : $m \sim \text{GeV}$ (can't solve the problems)

→ need baryon physics (SN, BH jets,...)

* $\sigma/m \approx 10^{-26} \text{ cm}^2/\text{g}$ (100 GeV WIMP)

- SIDM: $\sigma/m \sim 0.5\text{-}1$ (Cluster), 5 (dwarf) [cm^2/g],

→ velocity dependent σ ?

- WDM : $m \sim \text{keV}$

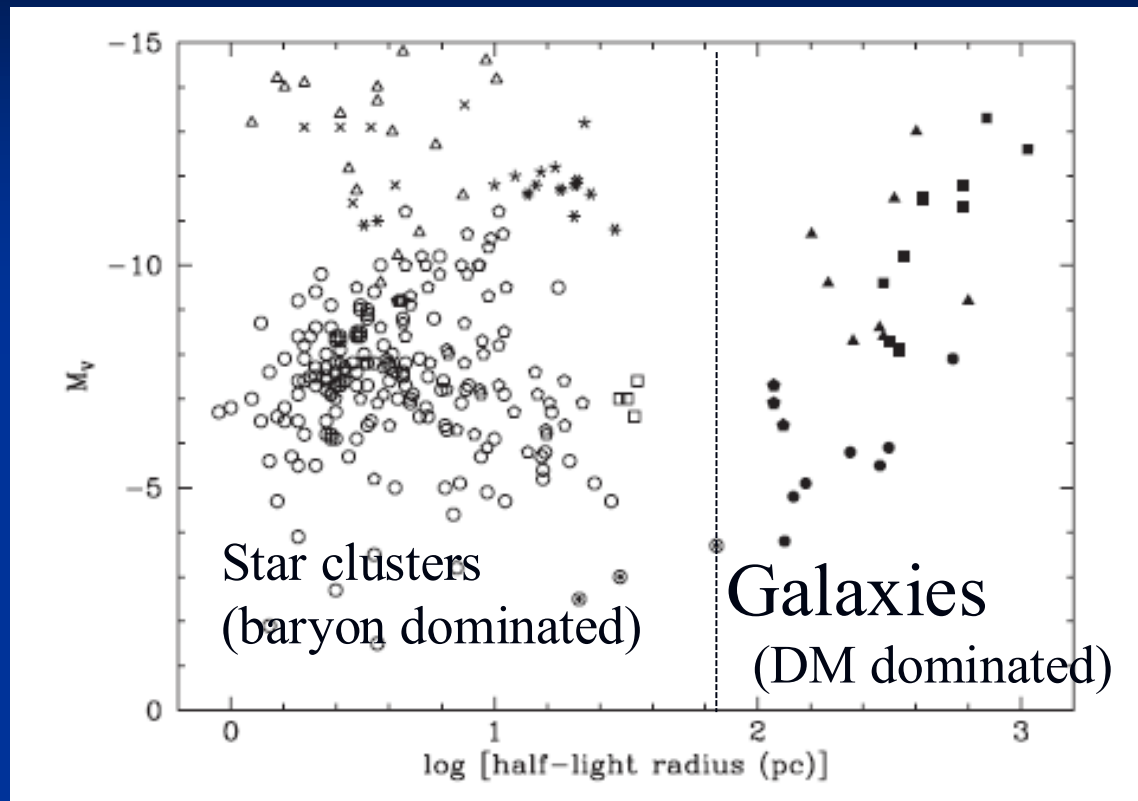
→ Catch 22 problem (cusp)

Suppression of the power spectrum prohibiting the formation of the dwarf galaxy

- ULDM: $m \sim 10^{-22} \text{ eV}$

→ Lyman alpha favors $m > 10^{-21} \text{ eV}$?

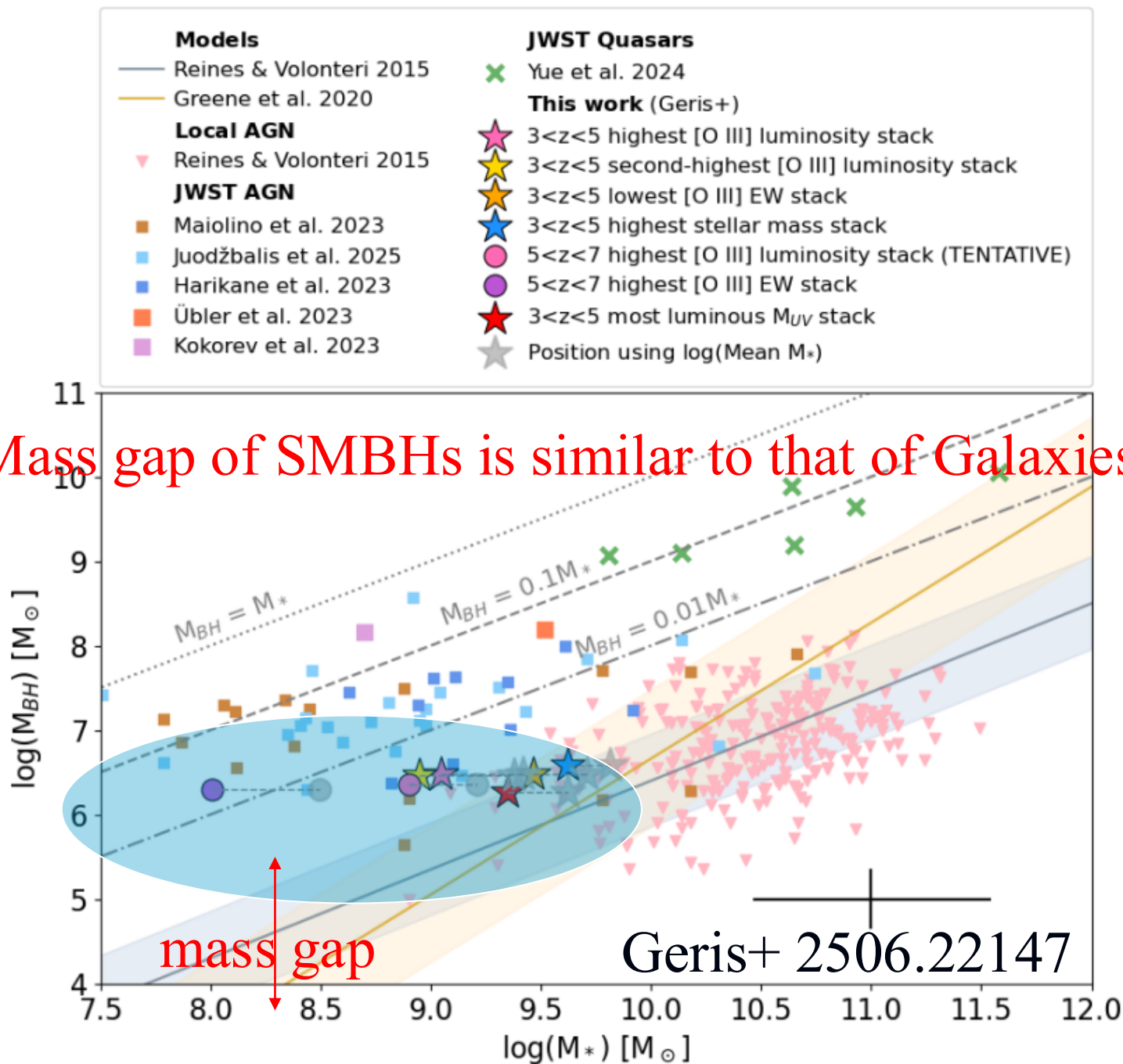
Characteristic size & mass of dwarf galaxies



- There are no known stable galaxies with half-light radius smaller than 120 pc while the maximum size for star clusters deprived of DM is about 30 pc .
- The minimal observed mass of a stable dwarf galaxy is on the order of $10^7 M_\odot$.
- The minimal observed mass of SMBHs is on the order of $10^5 M_\odot$.
- Max mass of SMBHs and galaxies $\sim 10^{11} M_\odot$.

→ *What is the origin of these scales (gaps)?*

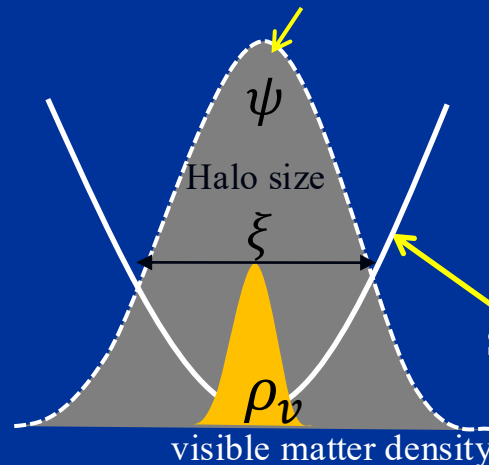
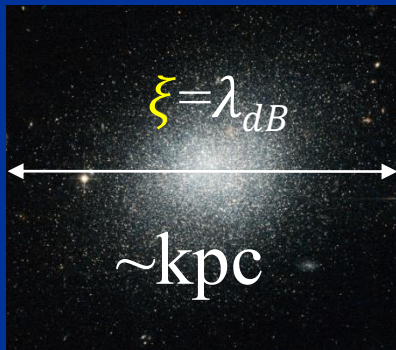
Mass gap of SMBHs is similar to that of Galaxies



Galactic DM halos in ULDM

- Galactic DM halos (cores) are made of **BEC ultra-light** bosons described by **macroscopic wave fn** (similar to solitons).
- Quantum pressure** (from **uncertainty principle**) prevents collapse
- Acts as CDM on super-galactic scales and **solves small scale crisis of CDM**

DM wave fn. of halo



uncertainty principle $\xi \times p > \hbar$
 min. halo size $\xi \approx \lambda_{dB} = \frac{h}{mv} \sim kpc$
 $\rightarrow m \sim 10^{-22} eV$ for $v \sim 10 km/s$

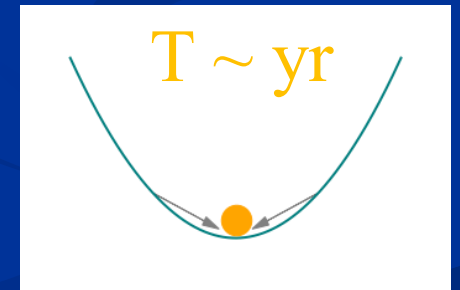
Schroedinger
-Poisson Eq
(SPE), **nonlinear**

$$\begin{cases} i\hbar\partial_t\psi = -\frac{\hbar^2}{2m}\nabla^2\psi + mV\psi & \text{DM density} \\ \nabla^2 V = 4\pi G(\rho_d + \rho_v), & \rho_d = m^2|\psi|^2 \end{cases}$$

Features of ULDM

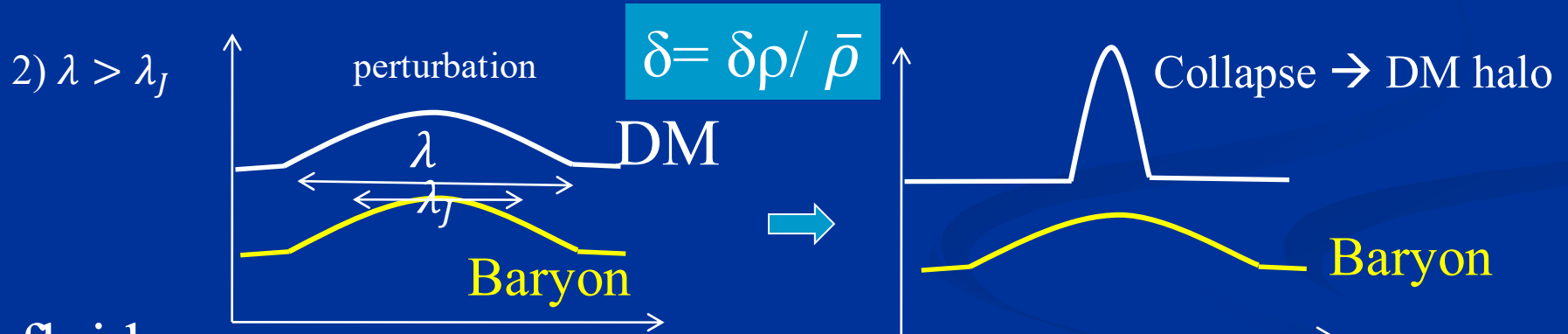
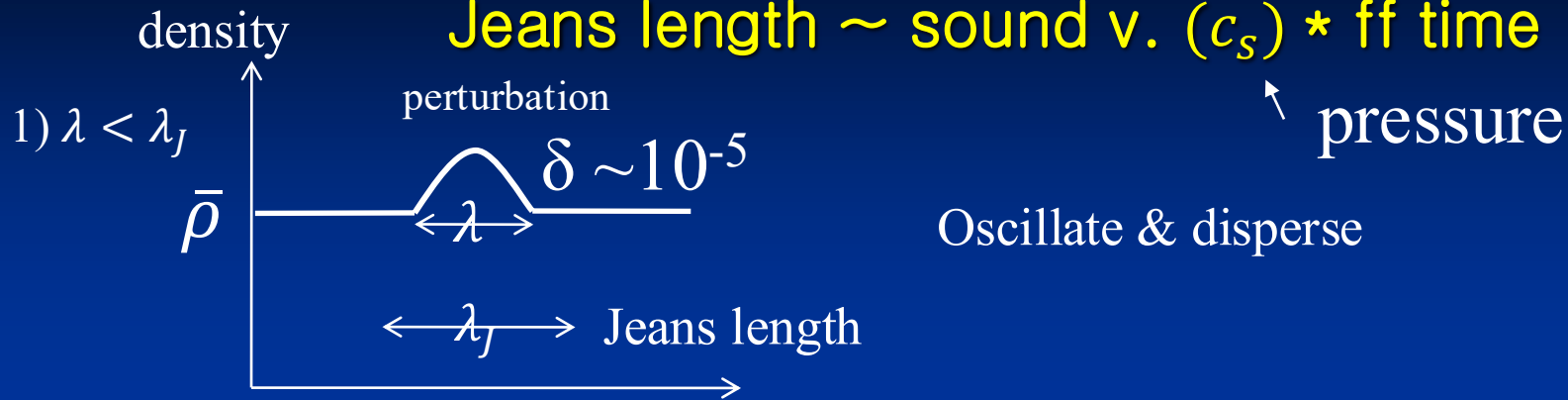
$$\phi(t, x) = \frac{1}{\sqrt{2m}} \left[e^{-imt} \underset{\substack{\text{fast (bg),} \\ T \sim \text{yr}}}{\psi(t, x)} + e^{imt} \underset{\substack{\text{slow (galaxy)} \\ T \sim \text{Myr}}}{\psi^*(t, x)} \right]$$

- Typical galaxy size $\sim \lambda_{dB} \sim \text{kpc}$
 $\rightarrow$ min. mass & max mass of galaxies
- wave nature $\rightarrow$ gravitational cooling, interf.
- small dynamical friction
- bg oscillation with $f \sim m \sim \text{nHz}$
- to explain DM relic density
 $\rightarrow$ GUT scale field value
 $\rightarrow$ can explain many mysteries of galaxies



How did cosmic structures form?:

Jeans length $\sim$ sound v. (c_s) * ff time

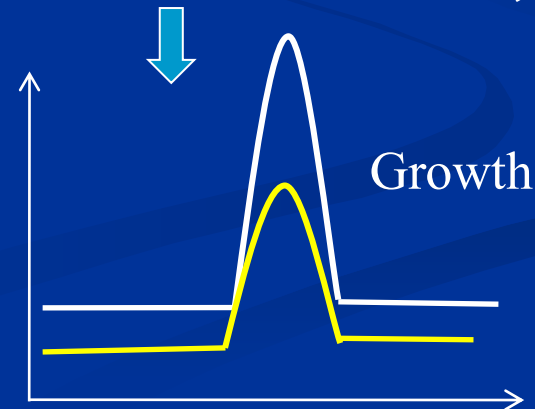


fluid

$$(\partial_t^2 - c_s^2 \nabla_r^2) \delta\rho = 4\pi G \bar{\rho} \delta\rho$$

$$k_J = \frac{\sqrt{4\pi G \bar{\rho}}}{c_s} = 2\pi / \lambda_J$$

Jeans length



Linear pert. Of ULDM

FDM has only 2 parameters m and bg density ρ_0 (or F)
 (+ λ for ϕ^4 self-interacting ULDM)

a =scale factor

Nonrelativistic
 $\longrightarrow$

Madelung
 representation

Density contrast
 (k space)

$$\begin{aligned}
 & i\hbar\left(\frac{\partial\psi}{\partial t} + \frac{3}{2}H\psi\right) = -\frac{\hbar^2}{2ma^2}\Delta\psi + mV\psi + \frac{\lambda|\psi|^2\psi}{2m^2} \quad \text{self-int} \\
 & \text{perturbation with } \psi = \sqrt{\rho}e^{iS}, \quad v \equiv \frac{\hbar}{ma}\nabla S \Rightarrow \\
 & \begin{cases} \partial_t\rho + 3H\rho + \frac{1}{a}\nabla \cdot (\rho v) = 0 \\ \partial_tv + \frac{1}{a}v \cdot \nabla v + Hv + \frac{1}{\rho a}\nabla p + \frac{1}{a}\nabla V + \frac{\hbar^2}{2m^2a^3}\nabla\left(\frac{\Delta\sqrt{\rho}}{\sqrt{\rho}}\right) = 0 \end{cases} \\
 & \text{perturbation } \delta = \delta_k = \delta\rho/\rho_0 \quad \text{Quantum Pressure} \\
 & \Rightarrow \partial_t^2\delta + 2H\partial_t\delta + \left(\left(\frac{\hbar^2k^2}{4m^2a^2} + c_s^2\right)\frac{k^2}{a^2} - 4\pi G\rho_0\right)\delta = 0 \\
 & \quad \text{Hubble drag} \quad \text{gravity}
 \end{aligned}$$

Quantum Jeans length

$$\lambda_{QJ} = \frac{2\pi}{k_J} a = \pi^{3/4} \hbar^{1/2} (G\rho_0 m^2)^{-1/4} \propto 1/\sqrt{mH} \sim 10 \text{ kpc}$$

- CDM-like on super-galactic scale (for a small $k < k_J$)
 - Suppress sub-galactic structure (for a large $k > k_J$)
- $\rightarrow$ ULDM is an ideal alternative to CDM

Galaxy mass from DM particle mass m

SIDM (particle) $M_J = \frac{\pi \bar{\rho}}{6} \left(\frac{\pi c_S^2}{G \bar{\rho}} \right)^{3/2} \leftarrow \lambda_J = c_S \sqrt{\frac{\pi}{G \bar{\rho}}} \leftarrow c_S = \sqrt{\frac{dp}{d\rho}}$

observation $\sigma/m \sim cm^2/g$

Baryon (H) $M_J \equiv \frac{4\pi}{3} \rho_0 \left(\frac{1}{2} \lambda_J \right)^3 = 26 M_\odot \left(\frac{T}{10 \text{ K}} \right)^{3/2} \left(\frac{n}{10^3 \text{ cm}^{-3}} \right)^{-1/2}$

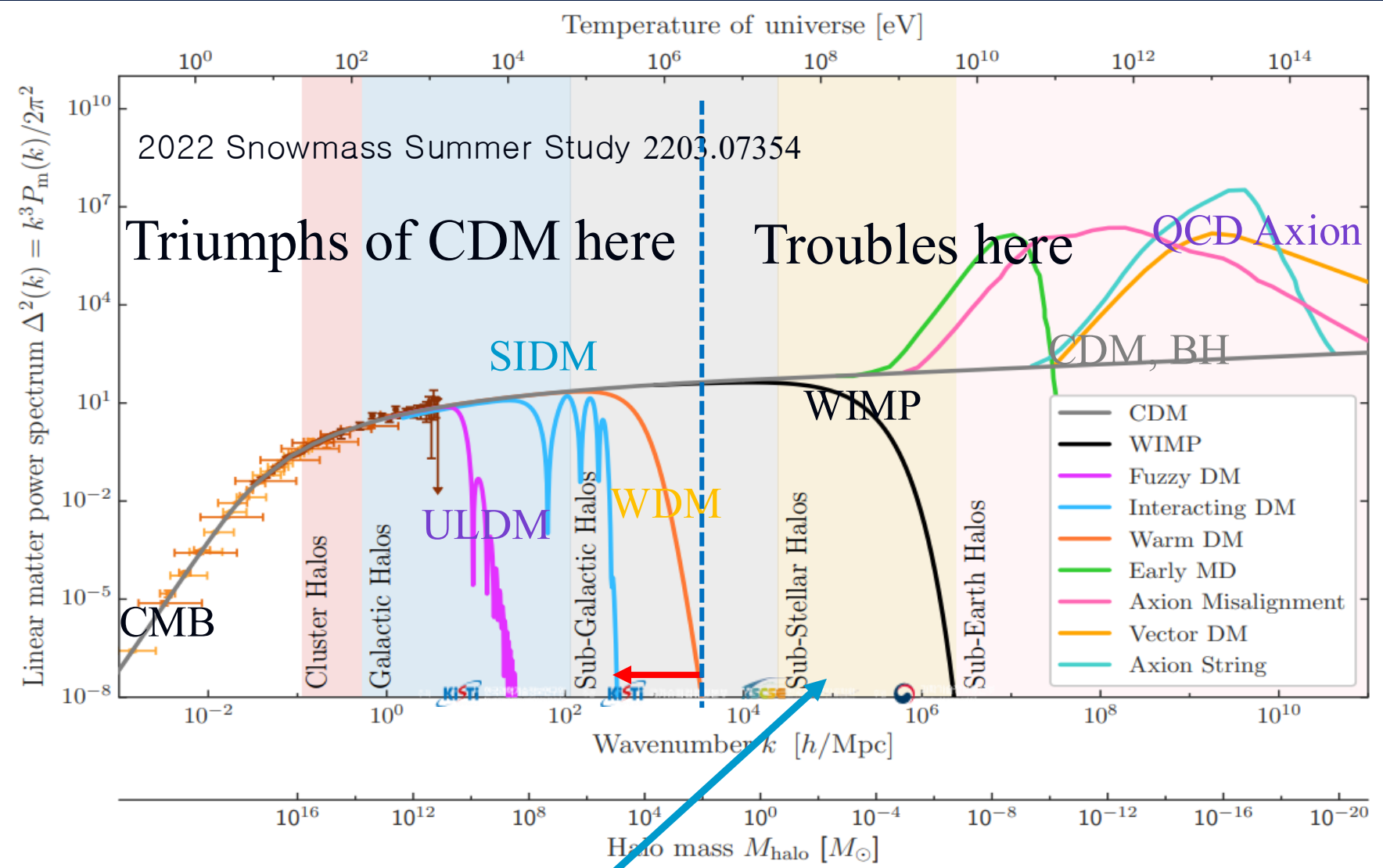
ULDM
(fuzzy) $M_{QJ}^Q = \frac{4\pi}{3} \rho(z) \lambda_{QJ}^3(z)$ Lee & Lim 0812.1342
 $= 1.204 \times 10^8 \left(\frac{1+z}{m_{22}^2} \right)^{3/4} \left(\frac{\Omega_{\text{dm}}}{0.27} \right)^{1/4} \left(\frac{h}{0.7} \right)^{1/2} M_\odot$

CDM, WDM
collisionless

$$M_{fs} = 10^{10} \left(\frac{m}{1 \text{ keV}} \right)^{-3.33} M_s$$

$$m_{22} \equiv \frac{m}{10^{-22} \text{ eV}}$$

Galaxies are DM dominated and seem to have kpc size scale



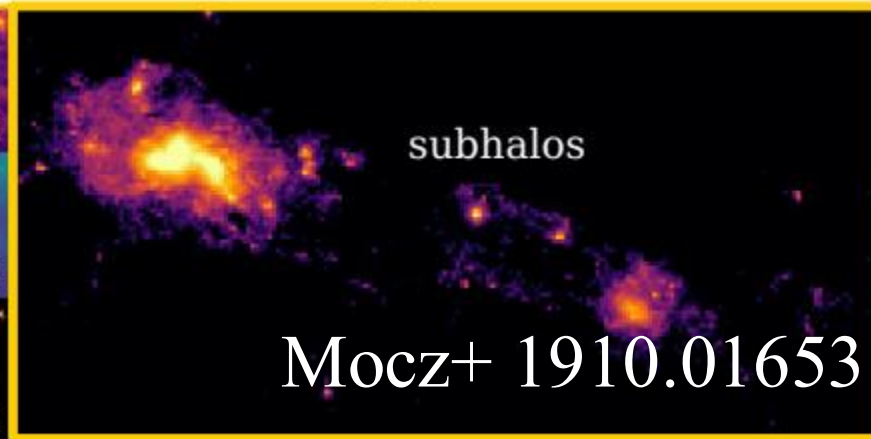
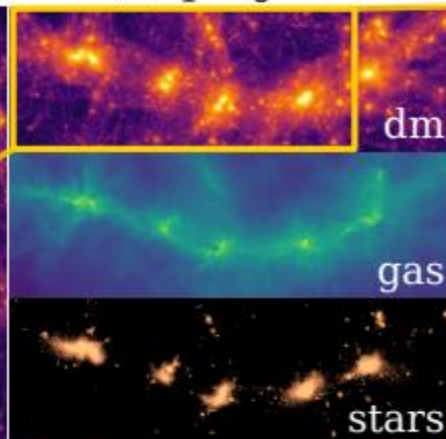
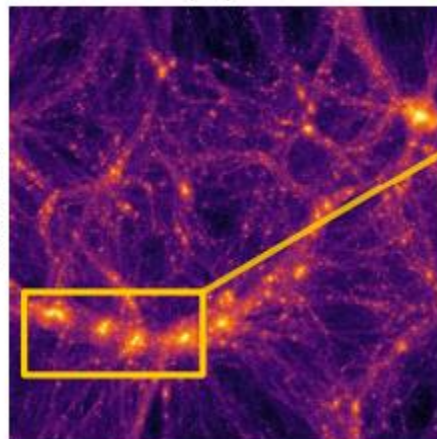
No cuspy subhalo, DM star or DM planet found so far
→ DM has kpc length scale?

(a) box

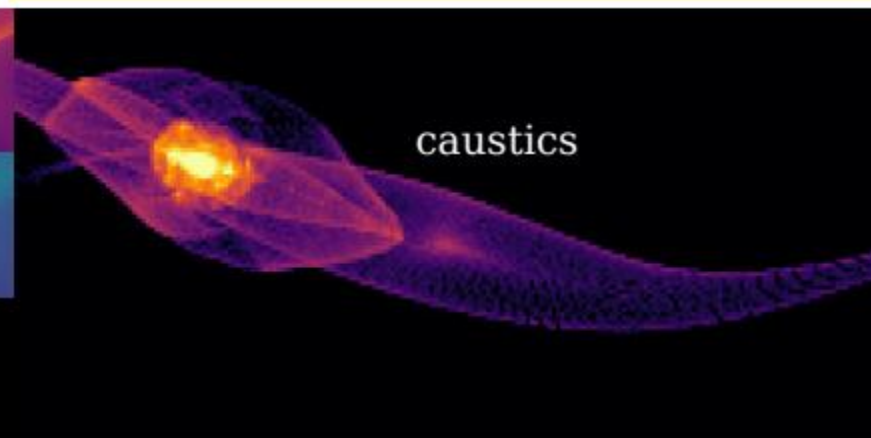
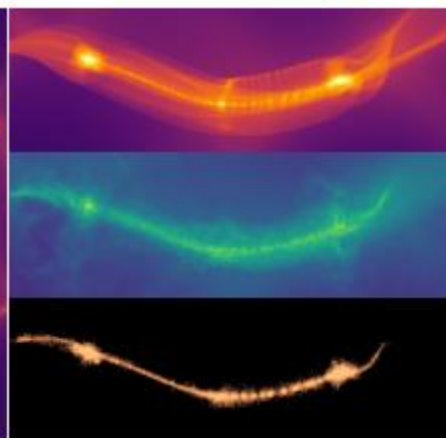
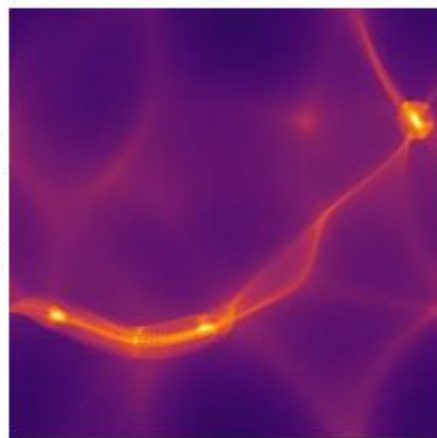
(b) projection

(c) slice

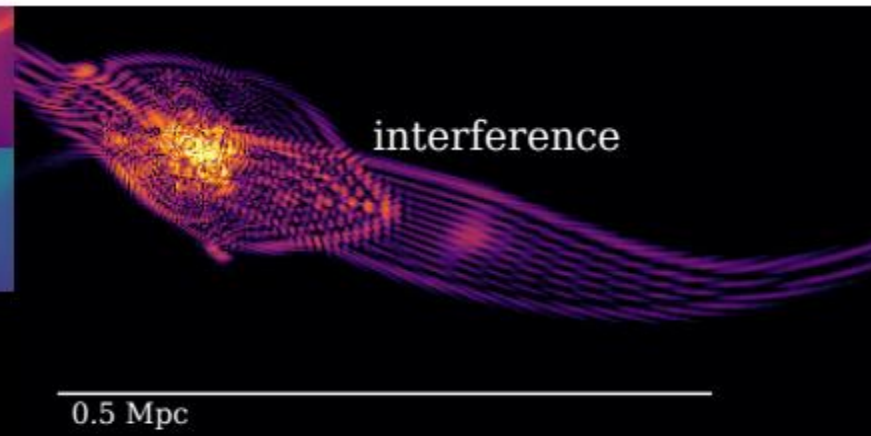
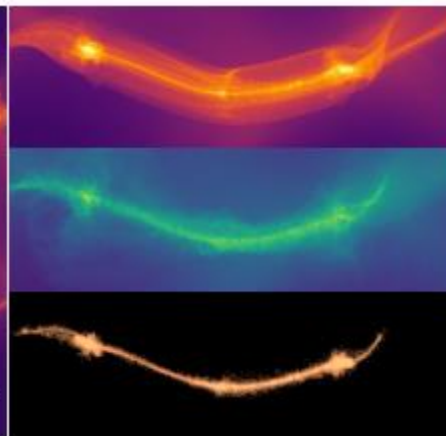
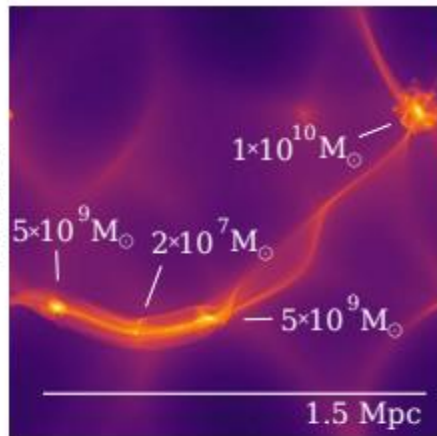
CDM

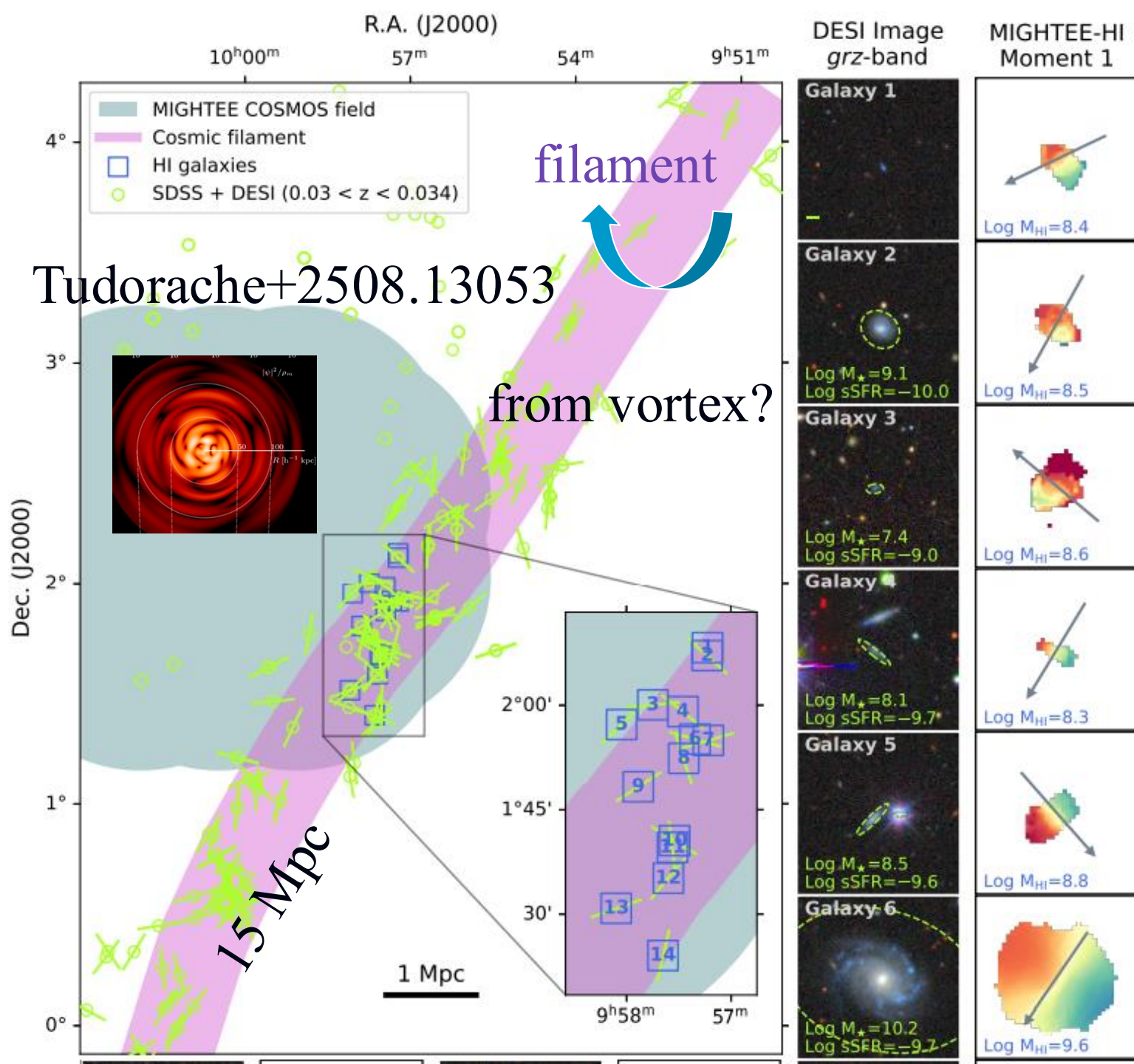


WDM



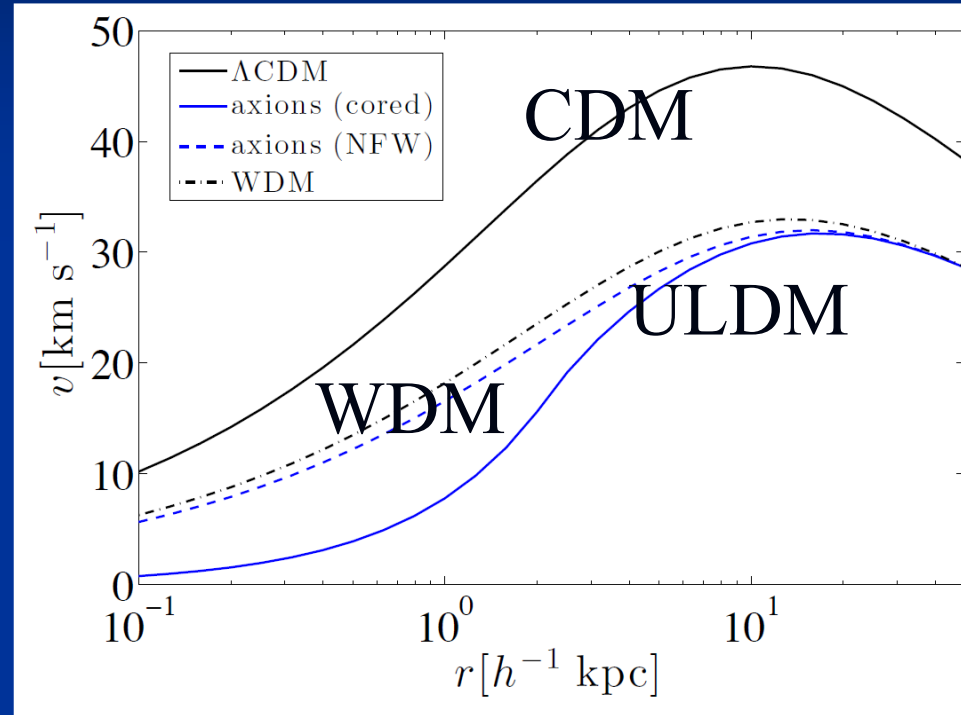
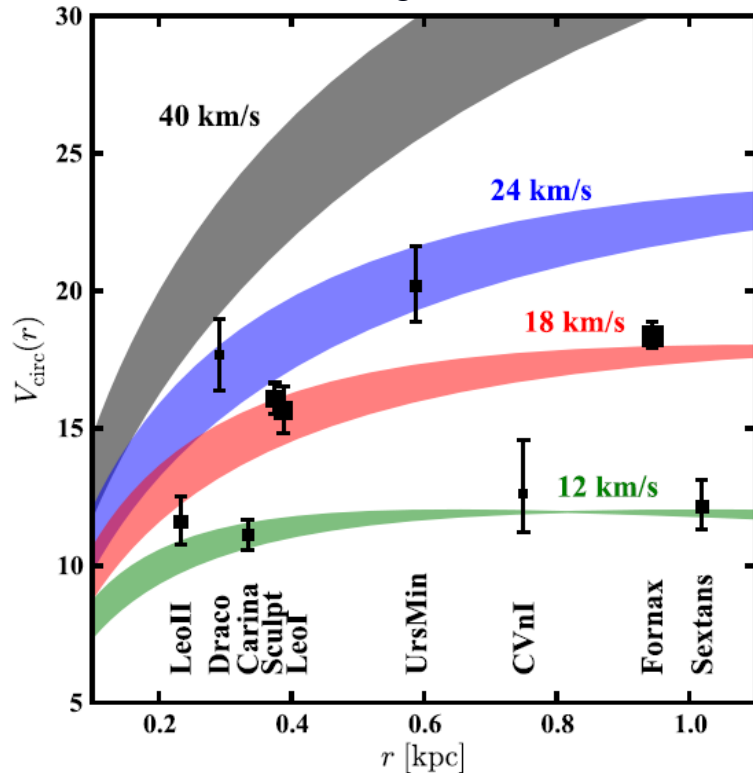
FDM





Too big to fail= no dense satellite

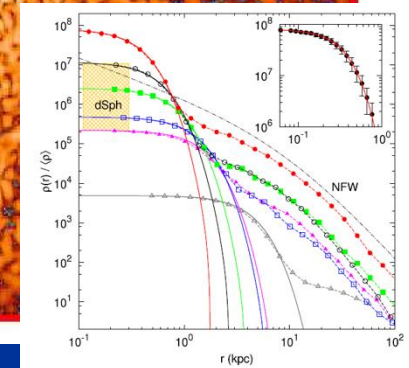
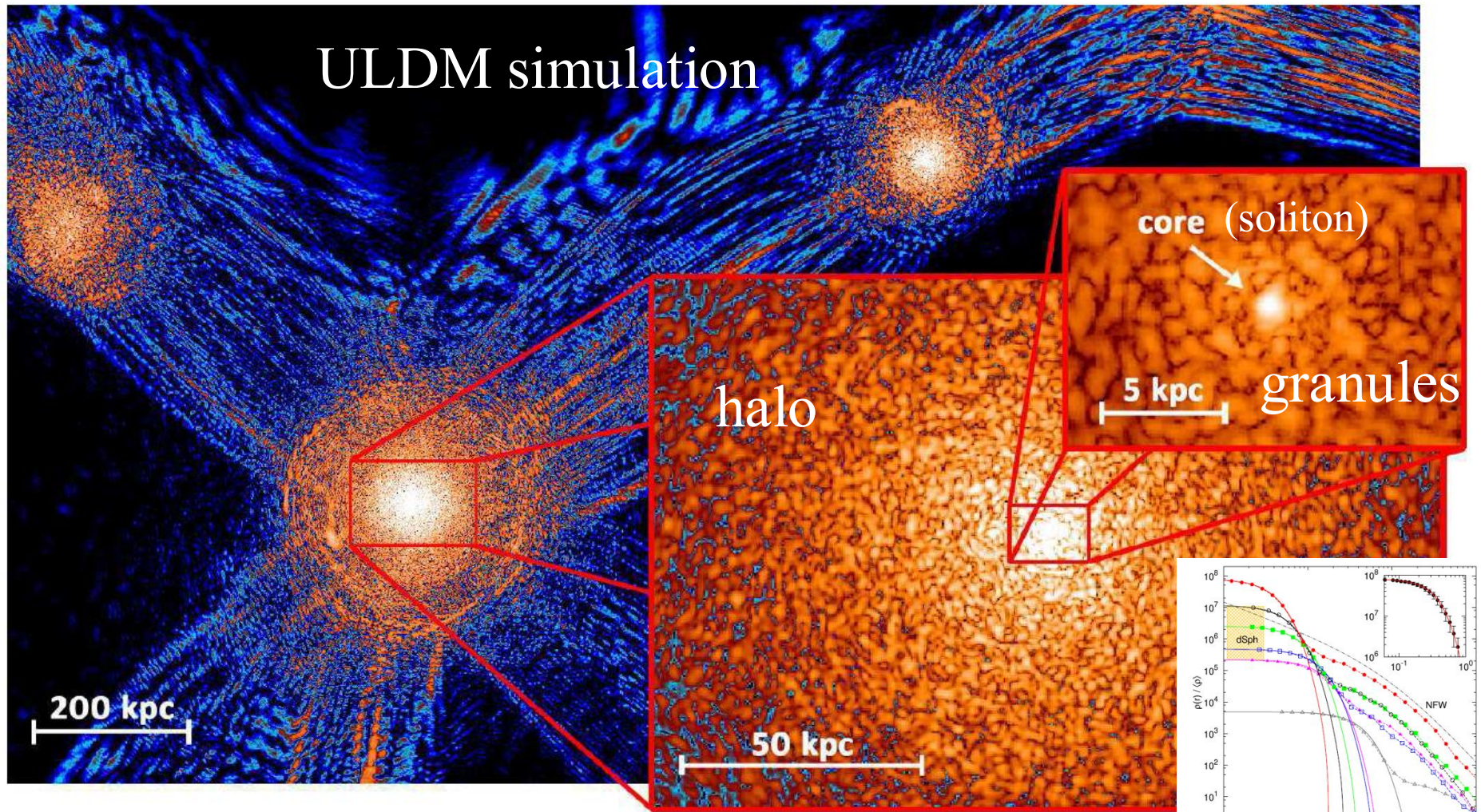
Where are these bright satellites?



Marsh & Silk 2013

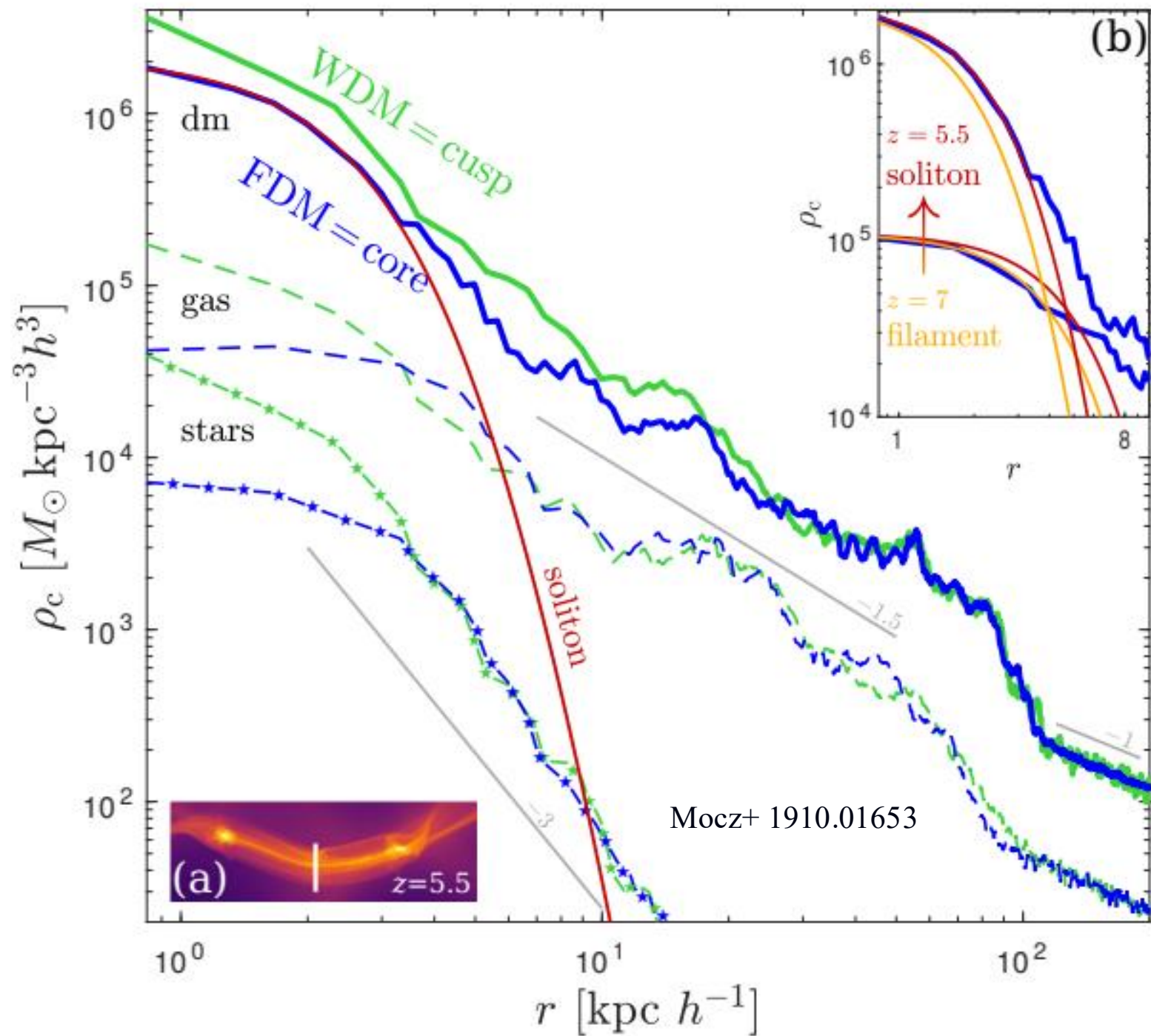
Boylan et al 2012

ULDM simulation



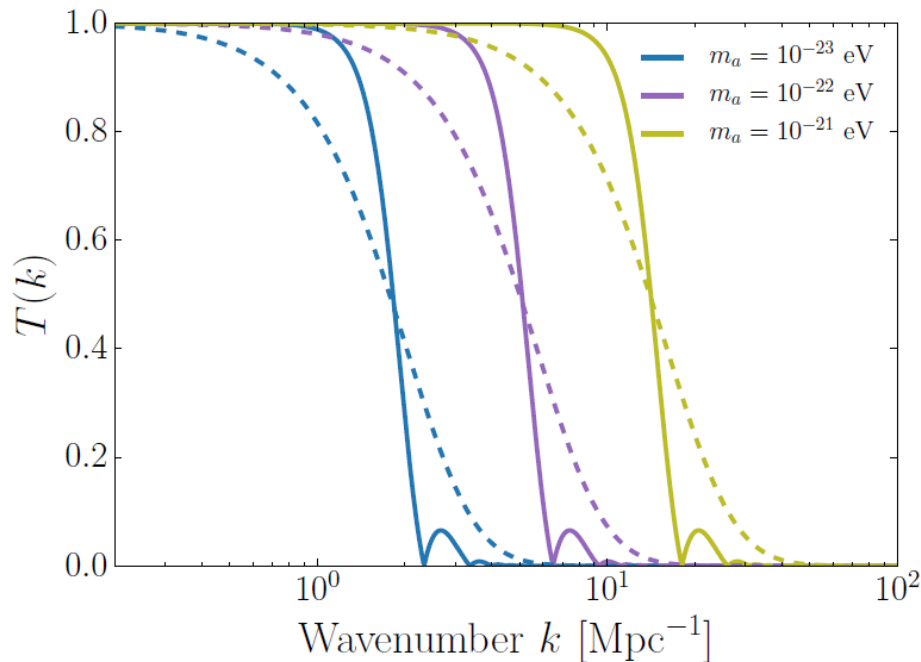
core size $\sim$ granule size $\sim$ typical length
 $\sim$ Q. Jeans length (kpc)

Schive+, Nature physics 2014



Can FDM solve small scale issues?

Hu et al PRL 2000 (Fuzzy DM)



$$P_{FDM} = T_F^2(k) P_{CDM}(k)$$

$$T_F^2(k) \approx \frac{\cos x^3}{1 + x^8}$$

$$x = 1.61 m_{22}^{1/18} k/k_J,$$

$$m = m_{22} \times 10^{-22} \text{ eV}$$

The power drops by a factor of

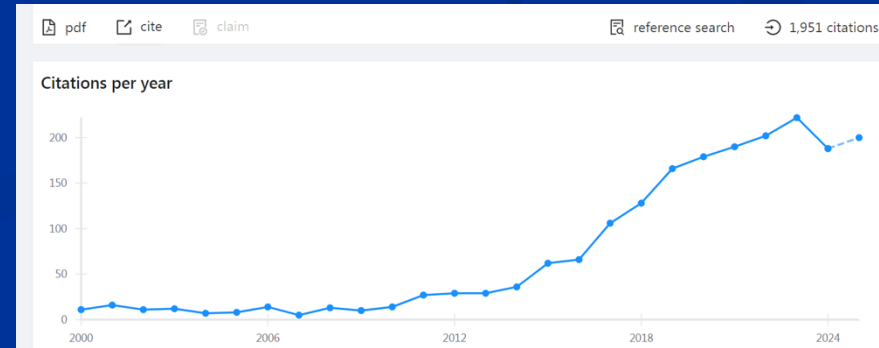
$$2 \text{ at } 4.5 m_{22}^{4/9} \text{ Mpc}^{-1}$$

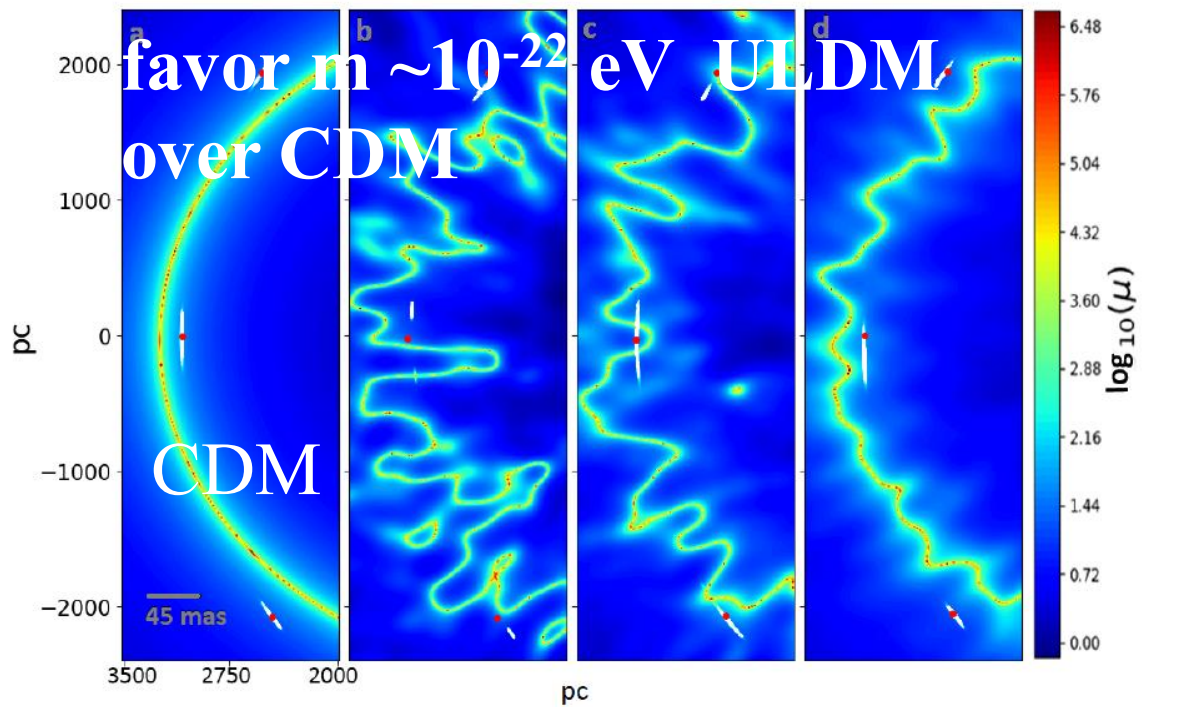
cutoff at $k \sim 4.5 h \text{ Mpc}^{-1}$

$$\Rightarrow m \approx 10^{-22} \text{ eV!}$$

Marsh

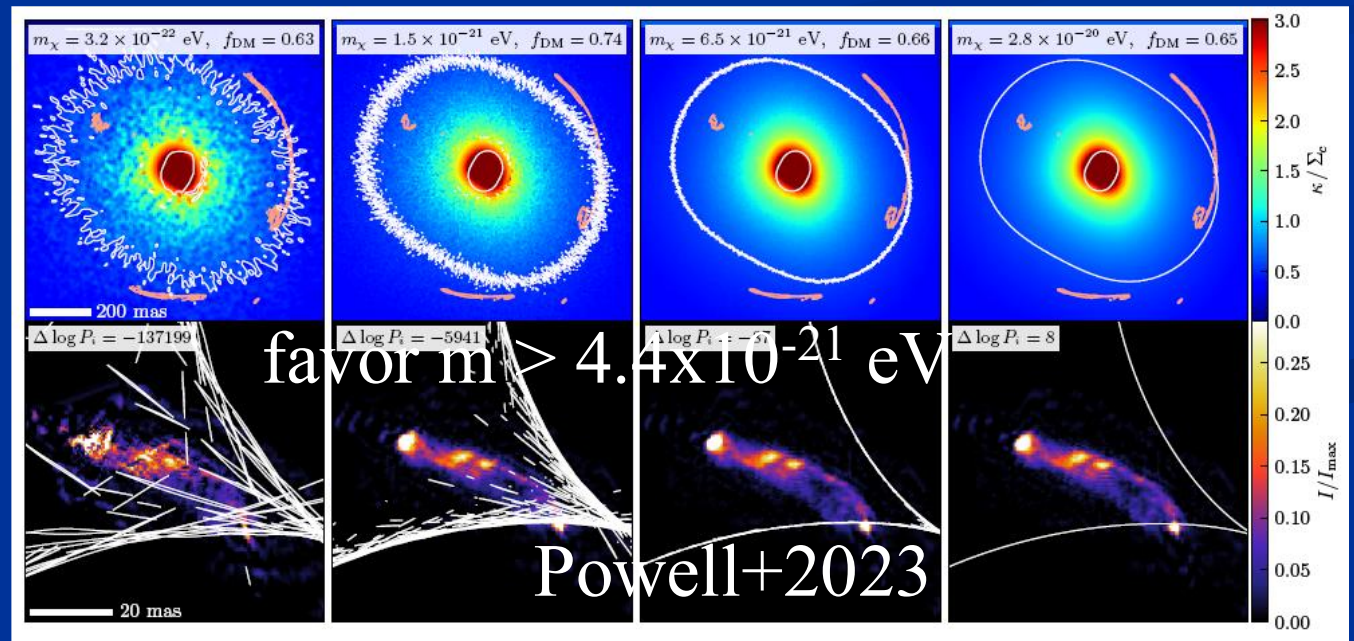
To solve the cusp and missing satellites suppression at small scale needed





ULLDM well reproduce
lens of radio objects

Armurth+2023

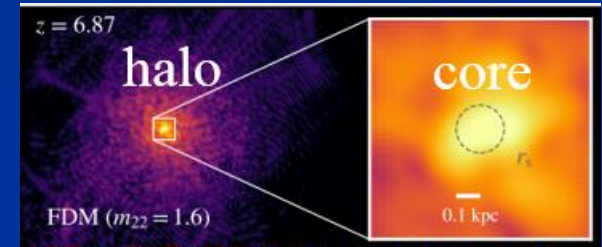


core-halo relation of ULDM

Q. Jeans mass

$$M_{Q_J} = \frac{4\pi}{3} \rho(z) \lambda_{Q_J}^3(z) = 1.204 \times 10^8 \left(\frac{1+z}{m_{22}^2} \right)^{3/4} \left(\frac{\Omega_{\text{dm}}}{0.27} \right)^{1/4} \left(\frac{h}{0.7} \right)^{1/2} M_{\odot}$$

$\sim M_h$ *halo mass*



DM core size

$$r_s = 0.135 m_{22}^{-1} \left(\frac{\zeta'(z)}{\zeta'(7)} \right)^{-1} \left(\frac{M_h}{10^{11} M_{\odot}} \right)^{-1/3} \text{ kpc}$$

DM core mass

$$M_s = 1.68 \times 10^9 m_{22}^{-1} \left(\frac{\zeta'(z)}{\zeta'(7)} \right) \left(\frac{M_h}{10^{11} M_{\odot}} \right)^{1/3} M_{\odot} \sim \text{Mass gap}$$

where $m_{22} \equiv \frac{m}{10^{-22} \text{ eV}}$, $\zeta' = (1+z)^{1/2} \zeta(z)^{1/6}$

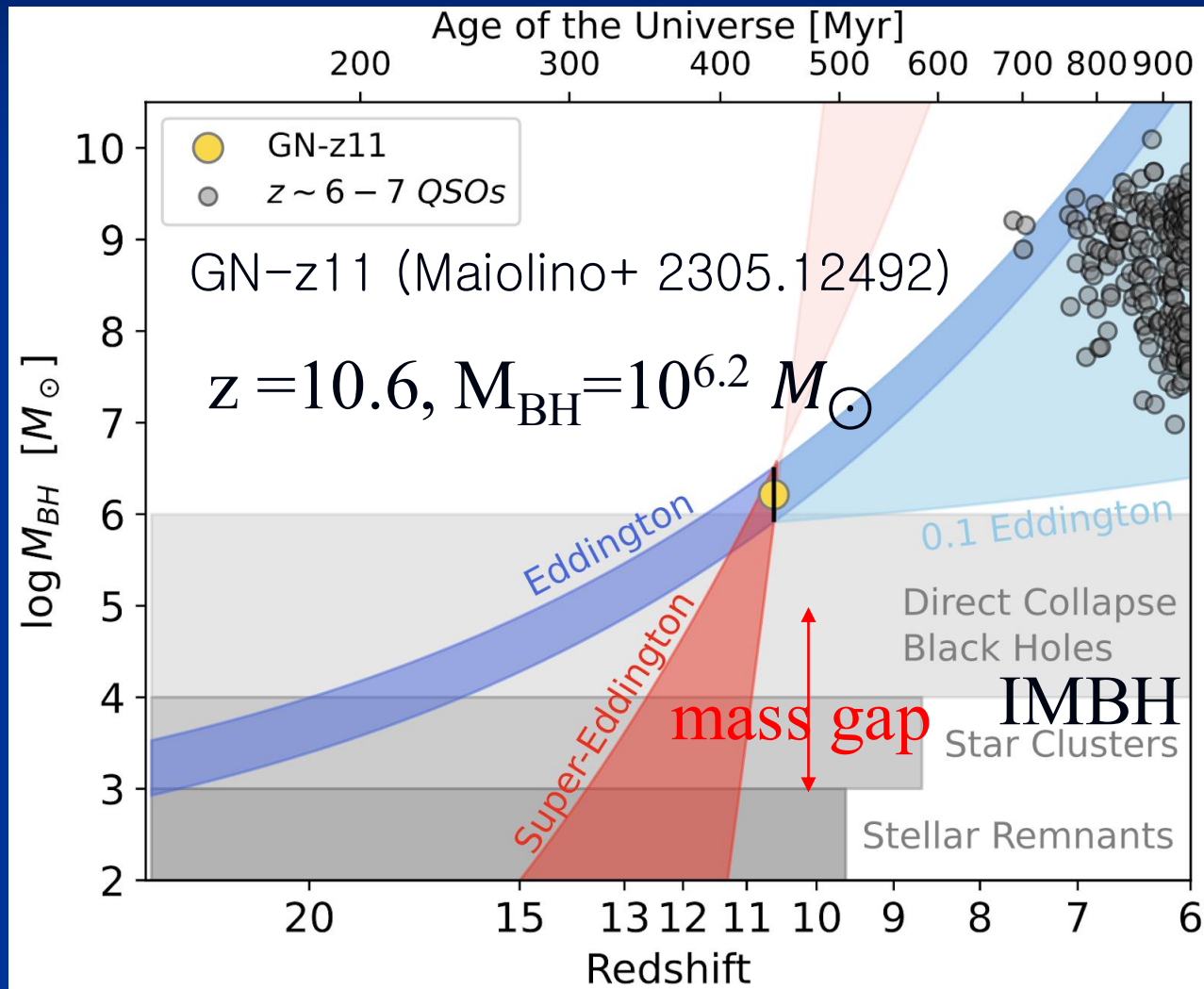
with $\zeta(z) = (18\pi^2 + 82(\Omega_m(z) - 1) - 39(\Omega_m(z) - 1)^2) / \Omega_m(z) \sim 180$ at $z \geq 1$

Impossibly early SMBH?

SMBHs (Quasar) with $M \sim 10^6 M_\odot$ were already mature at around $z=11$ (~ 0.45 Gyr)

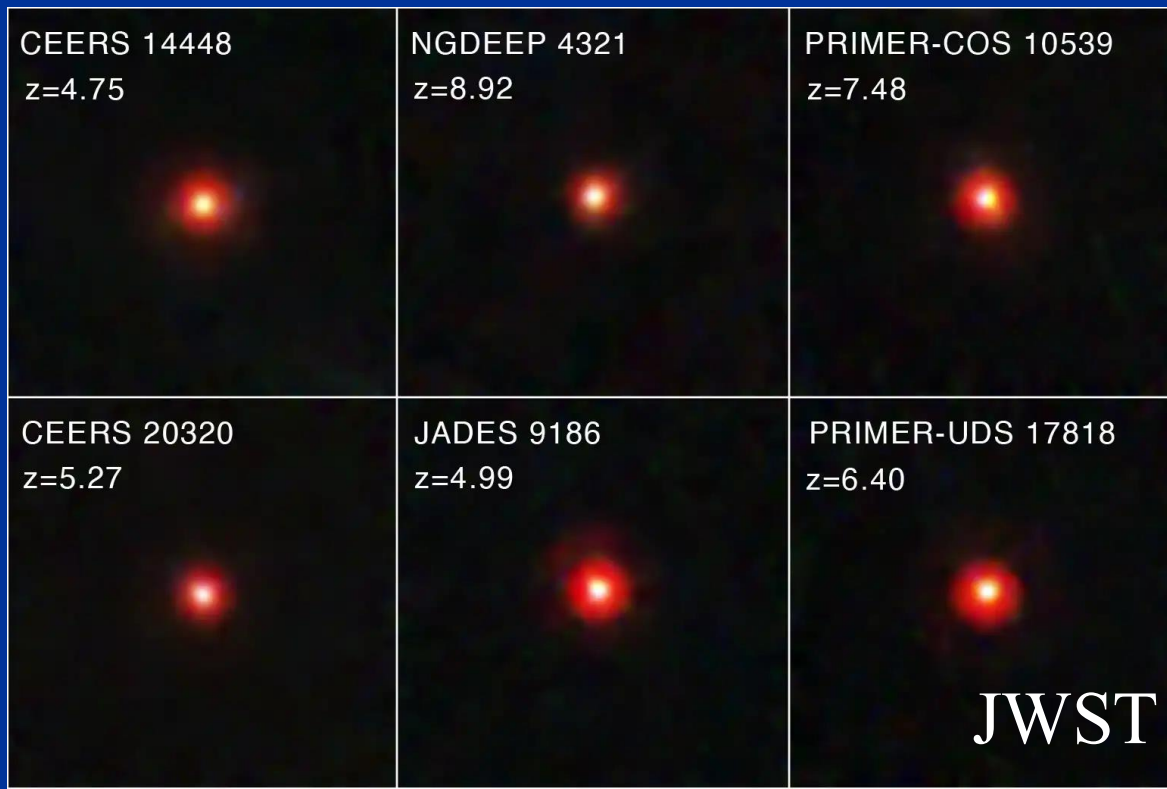


Cosmic
chicken-egg
problem



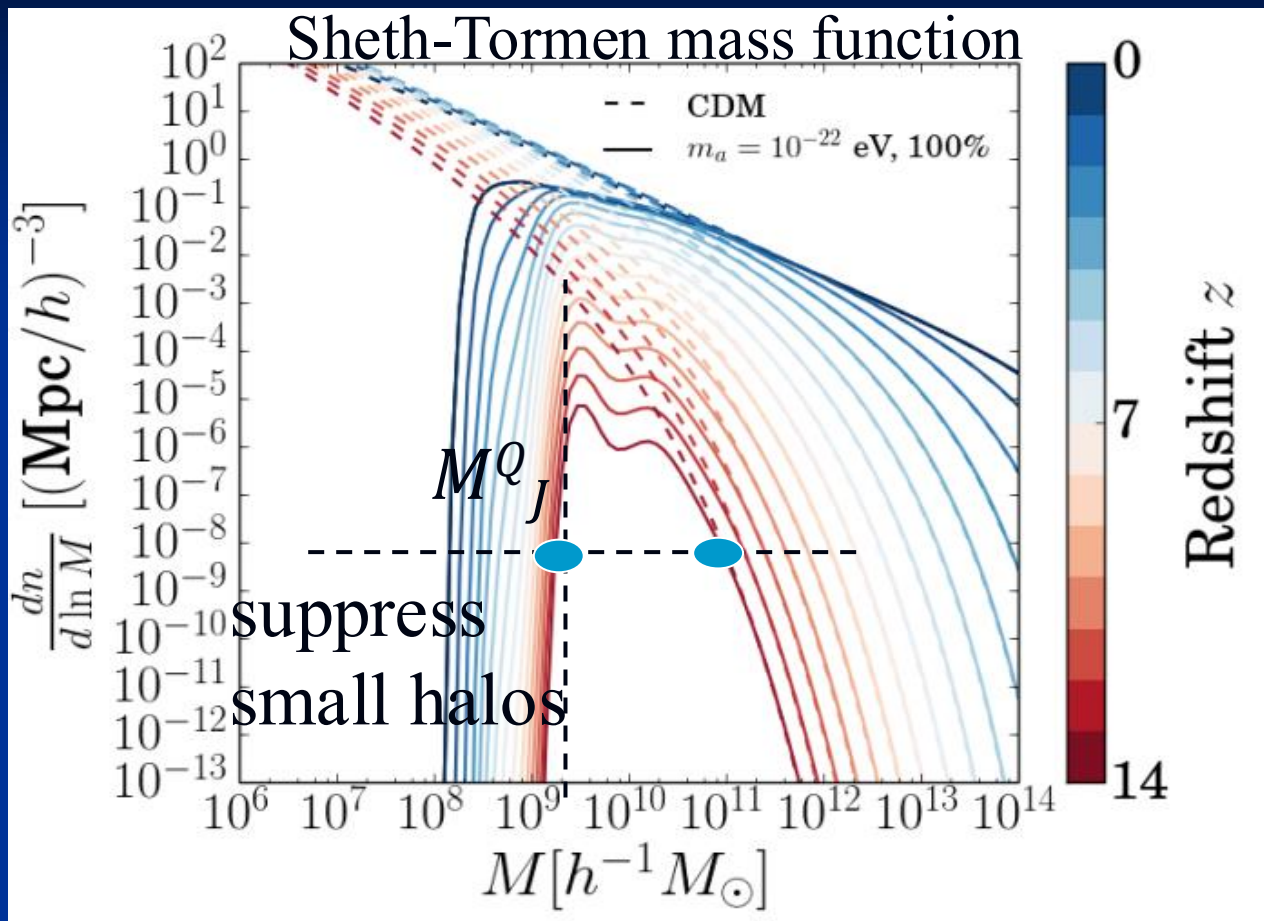
Little Red Dots

- Size: less than a few hundred light-years (smaller, older than quasar hosts)
- Color: Red; influenced by redshift and dust absorption?
- Epoch of Formation: Approximately 0.6–1.5 billion years after the Big Bang
- Central Structure: Rapid gas rotation → high likelihood of hosting an AGN
- Observational Features: Weak X-ray emission, no photometric variability
- Possible Models: Obscured AGN?, supermassive stars? Black hole stars?



Halo mass function

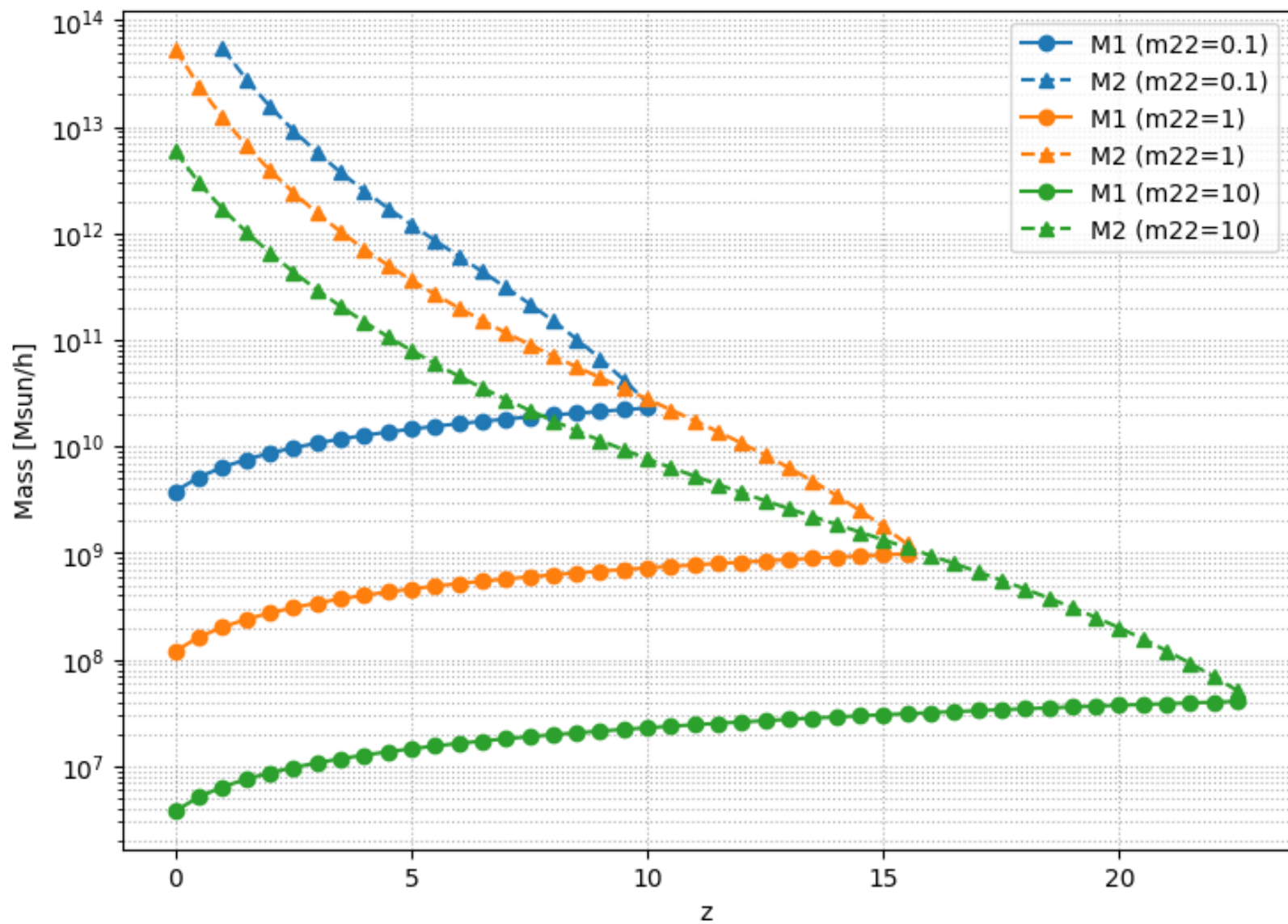
Bozek+ 1409.3544



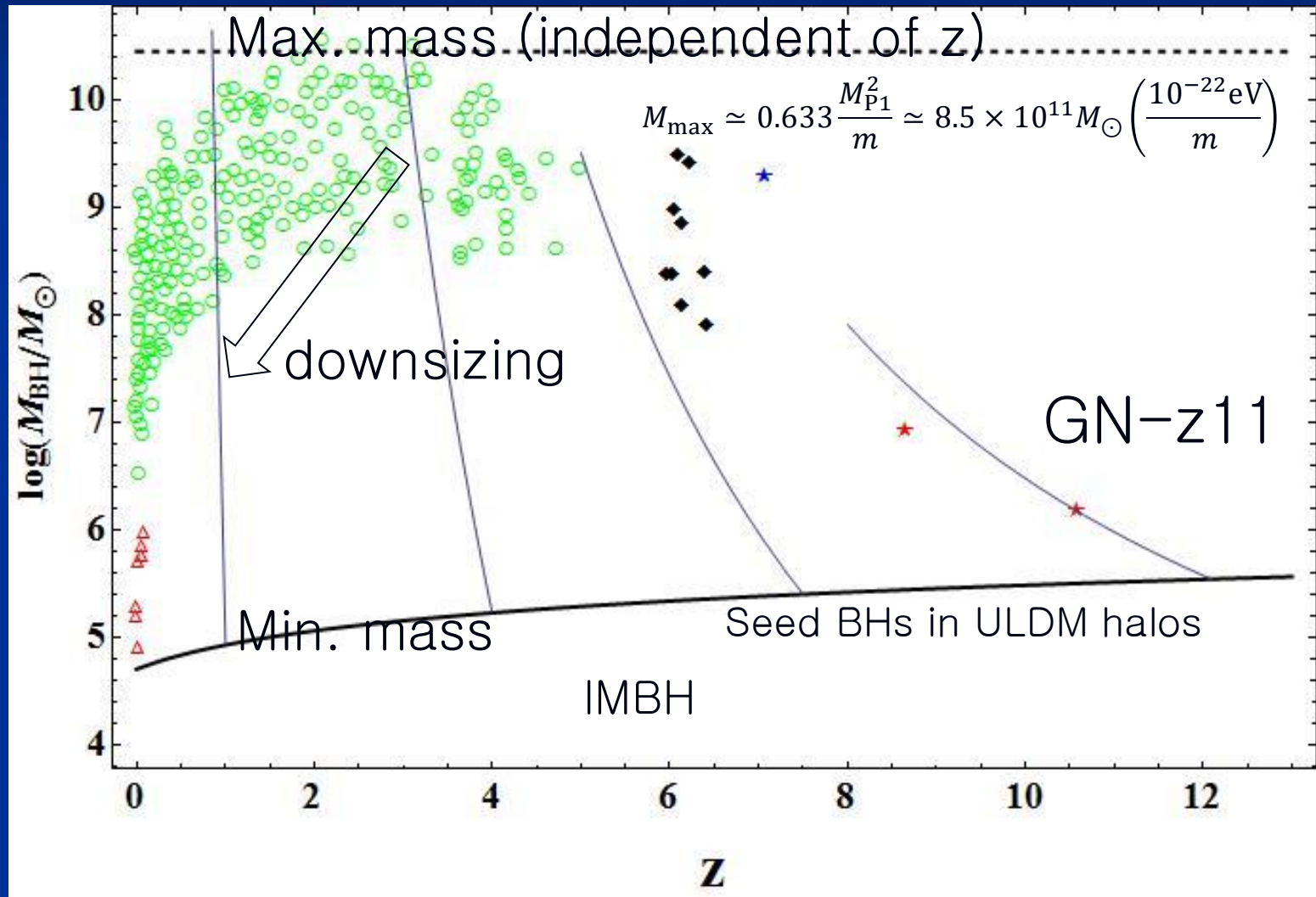
$$M_h = \alpha^3 M_J^Q$$

$$\left. \frac{dn}{d \ln M} \right|_{\text{FDM}} = \left. \frac{dn}{d \ln M} \right|_{\text{CDM}} \left[1 + \left(\frac{M_h}{M_0} \right)^{-1.1} \right]^{-2.2} \quad M_0 = 1.6 \times 10^{10} m_{22}^{-4/3} M_{\odot}$$

$$\left. \frac{dn}{d \ln M} \right|_{\text{CDM}} \propto f(v) = A \sqrt{\frac{1}{2\pi}} \sqrt{q} v [1 + (\sqrt{q} v)^{-2p}] \exp\left(-\frac{q v^2}{2}\right) \frac{1}{v} \quad v = \frac{1.69}{\sigma(M)}$$

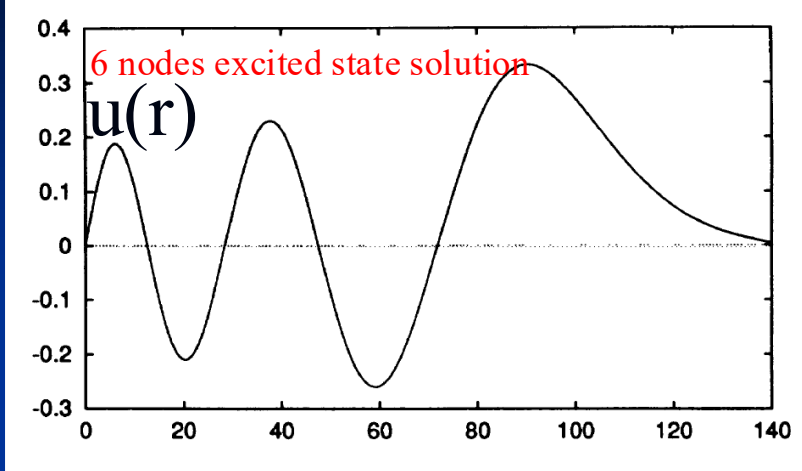


Evolution of SMBHs from ULDM



BEC DM

S. Sin PRD 1994

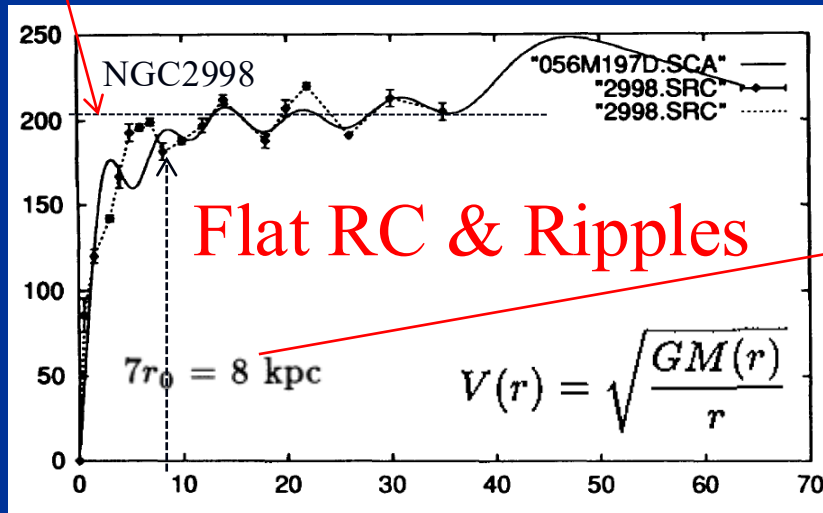


$$\hat{\psi}(\hat{r}) = \frac{1}{\sqrt{4\pi}} \frac{u(\hat{r})}{\hat{r}} \quad \text{rescaled}$$

$$v_{\text{rotation}} \lambda_{\text{ripple}} = \frac{1}{\sqrt{2}} \frac{1}{m} f(\hat{r})$$

$$f(\hat{r}) = \left(\int_0^{\hat{r}} d\hat{r} u^2(\hat{r}) / \hat{r} [\hat{V}(\hat{r}_t) - \hat{V}(\hat{r})] \right)^{\frac{1}{2}}$$

$$0.32v_0 = 204 \text{ km/sec.}$$



$$\rho_{\text{dark}} = GM_0 m |\psi|^2$$

Gravitational Bohr radius

$$r_0 = \hbar^2 / 2GM_0 m^2 \rightarrow$$

$$m = 3.3 \times 10^{-23} \text{ eV}$$

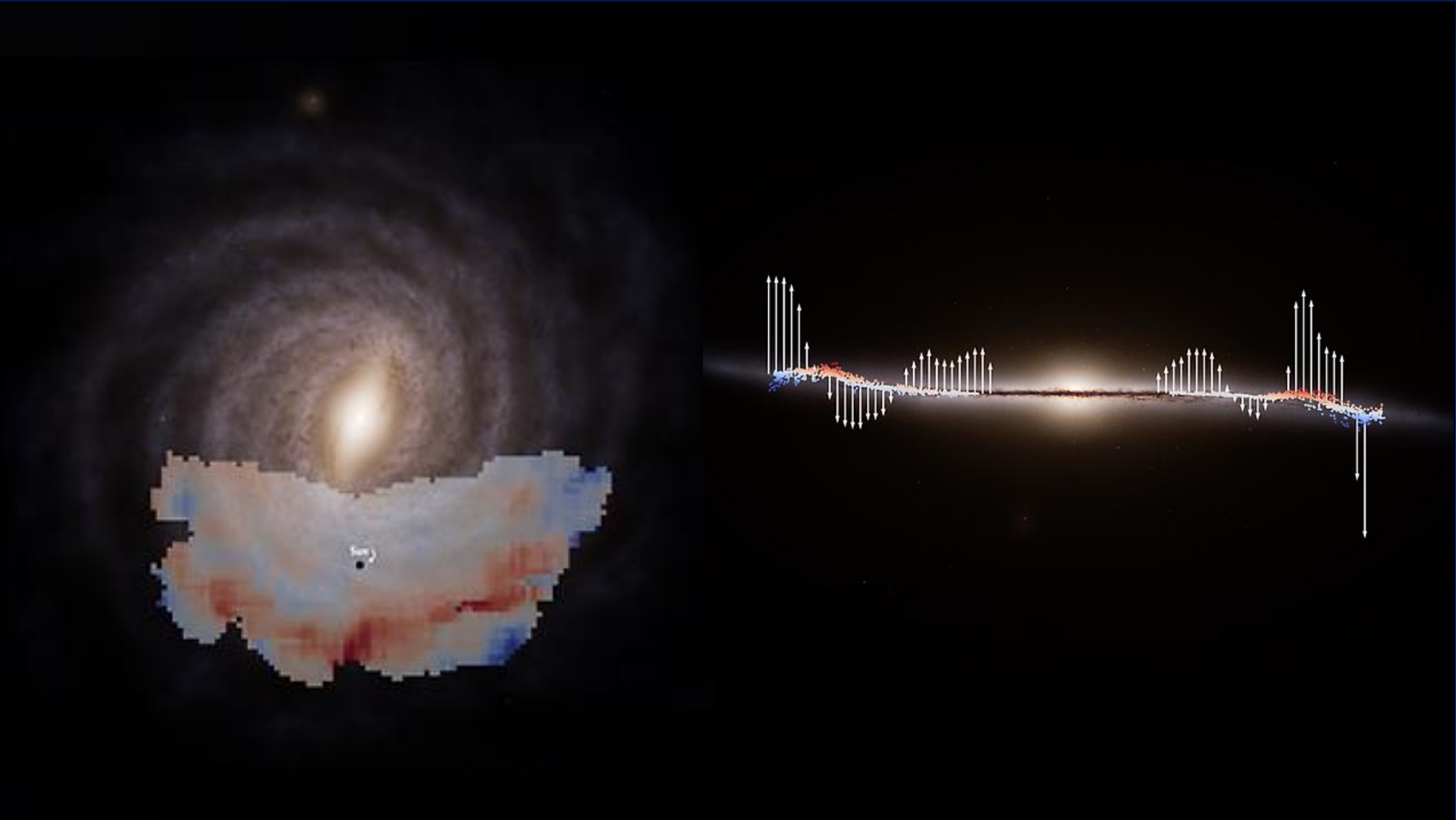
$$M_0 = 0.69 \times 10^{12} M_{\odot}$$

$$v_0 = \sqrt{GM_0/r_0}$$

Total mass

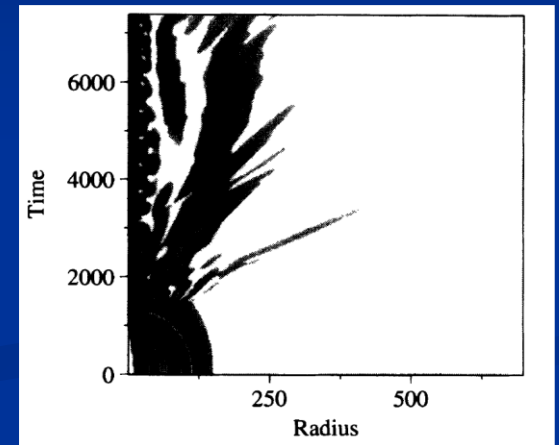
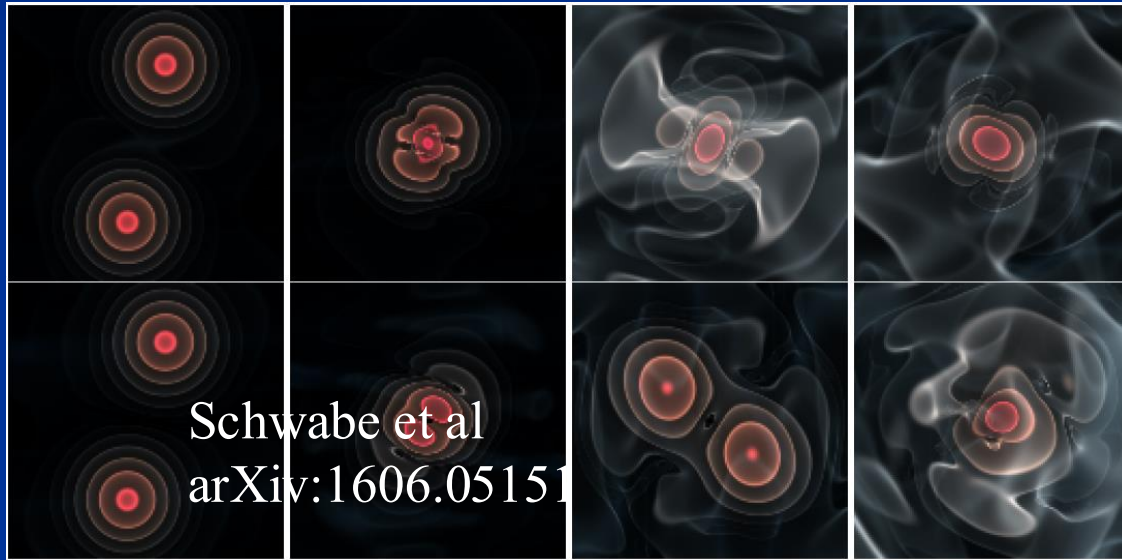
$$M = \hat{M} M_0 = 5.9 \times 10^{12} M_{\odot}$$

Warp, Ripple, Wave



Gravitational cooling of FDM

Unique and efficient dissipation mechanism of FDM



Seidel, Suen PRL 72

DM halo (wave) collision $\rightarrow$ nonlinear interference
 $\rightarrow$ creation of high momentum (k) modes
 $\rightarrow$ escape of waves carrying E , p , L a way

Oscillation of ULDM

$$\phi(\mathbf{x}, l) = \bar{\phi}(l) + \delta\phi(\mathbf{x}, l)$$

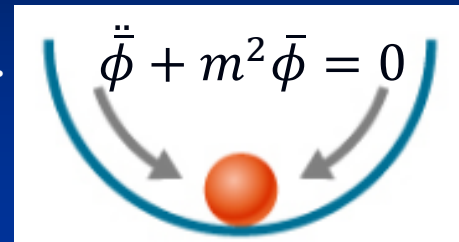
where

- $\bar{\phi}(l)$: spatially homogeneous cosmic background field,
- $\delta\phi(\mathbf{x}, l)$: inhomogeneous part forming nonlinear galactic halos.

For a scalar potential

$$V(\phi) = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4}\phi^4$$

$$\omega = \frac{1}{2.5\text{months}} \frac{m}{10^{-22}\text{eV}}$$



$$\ddot{\phi} + 3H\dot{\phi} - \frac{1}{a^2}\nabla^2\phi + m^2\phi + \lambda\phi^3 = 0$$

Inserting $\phi = \bar{\phi} + \delta\phi$ and spatially averaging, the background satisfies

$$\ddot{\bar{\phi}} + 3H\dot{\bar{\phi}} + m^2\bar{\phi} + \lambda\bar{\phi}^3 + \lambda\bar{\phi}\langle\delta\phi^2\rangle + \dots = 0.$$

The local perturbation $\phi_h \equiv \delta\phi$ obeys

$$\ddot{\phi}_h + 3H\dot{\phi}_h - \frac{1}{a^2}\nabla^2\phi_h + m^2\phi_h + 3\lambda\bar{\phi}^2\phi_h + \dots = 0$$

this simplifies to

$$\ddot{\phi}_h - \nabla^2\phi_h + m^2\phi_h + \lambda\phi_h^3 = 0.$$

ULA miracle

Hui et al 2017

$$I = \int d^4x \sqrt{g} \left[\frac{1}{2} F^2 g^{\mu\nu} \partial_\mu a \partial_\nu a - \mu^4 (1 - \cos a) \right]$$

$$m = \frac{\mu^2}{F}, \quad F = \text{typical field value}$$

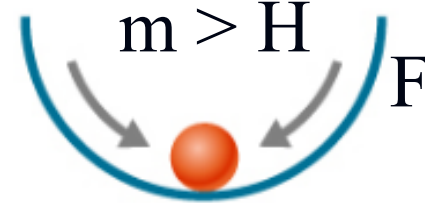
$$\ddot{a} + 3H\dot{a} + m^2 \sin a = 0$$

$$\text{oscillation starts at } H \sim \frac{T_{\text{osc}}^2}{M_P} = m$$

$$\text{MDE starts at } T_1 \sim 1 \text{ eV} \rightarrow \frac{\mu^4(\text{DM})}{T_{\text{osc}}^4(\text{rad})} \rightarrow \frac{\mu^4 T_{\text{osc}}}{T_{\text{osc}}^4 T_1} \sim 1$$

$$F = \frac{\mu^2}{m} \sim \frac{M_P^{3/4} T_1^{1/2}}{m^{1/4}} \sim 10^{17} \text{ GeV} \quad \text{typical field value}$$

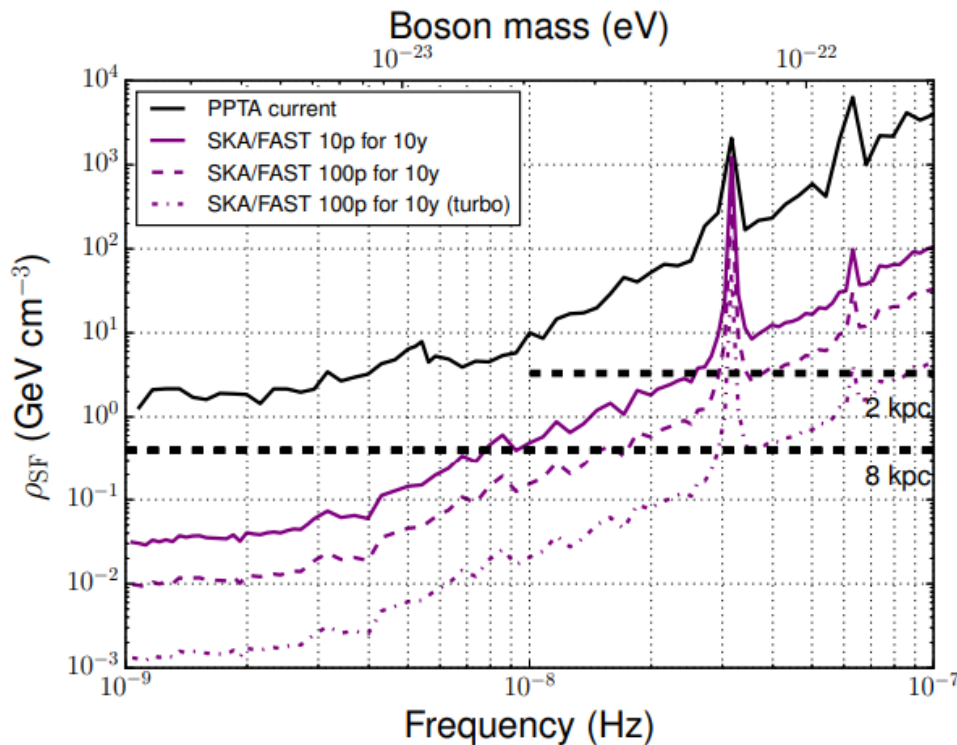
$$\Omega_a \sim 0.1 \left(\frac{F}{10^{17} \text{ GeV}} \right)^2 \left(\frac{m}{10^{-22} \text{ eV}} \right)^{1/2} \quad \text{ULA miracle?}$$



ULDM naturally explains DM density with GUT scale.
 This holds for ULDM with a quadratic pot.

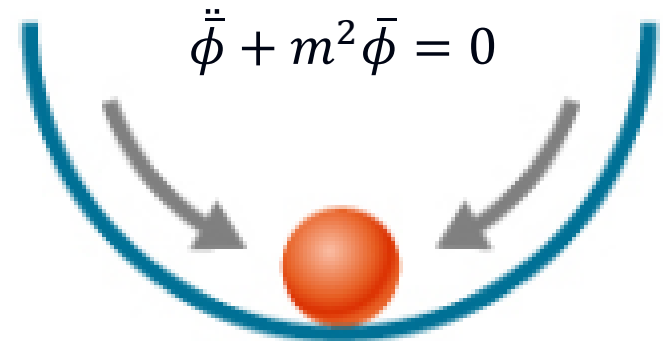
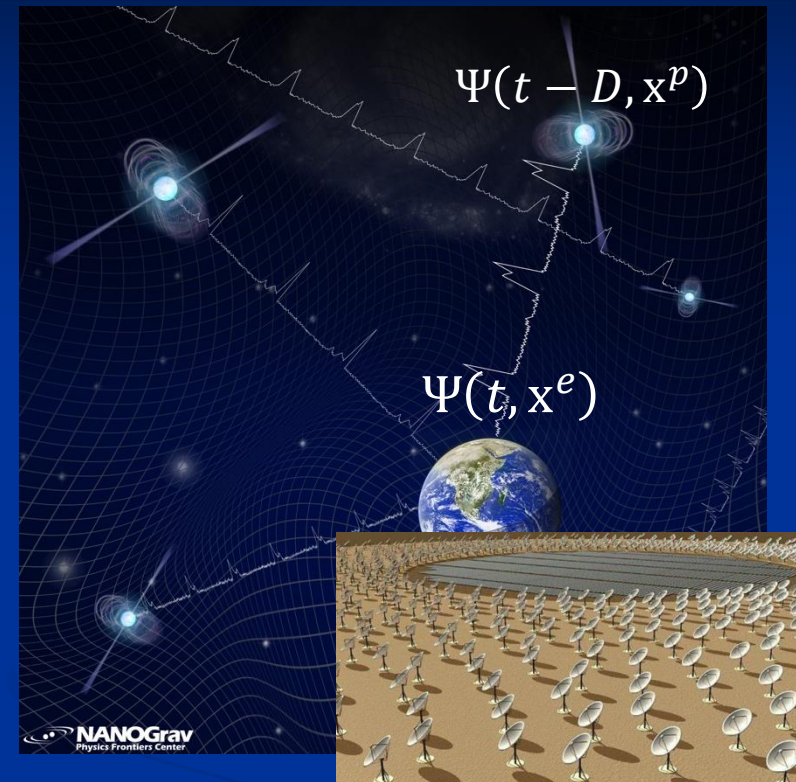
GW background detected by pulsar timing array

1810.03227



ULDM has
intrinsic osc time scale
 $1/m \sim \text{yrs}$

$$\omega = \frac{1}{2.5 \text{ months}} \frac{m}{10^{-22} \text{ eV}}$$



Gravitational atom



$$(g^{\alpha\beta} \nabla_{\alpha} \nabla_{\beta} - \mu^2) \Phi(t, r) = 0$$

ansatz

$$\Phi(t, r) = \frac{1}{\sqrt{2\mu}} \left[\psi(t, r) e^{-i\mu t} + \psi^*(t, r) e^{i\mu t} \right]$$

Schroedinger equation
with a Coulomb-like
central potential

$$i\partial_t \psi(t, r) = \left(-\frac{1}{2\mu} \nabla^2 - \frac{\alpha}{r} + O(\alpha^2) \right) \psi(t, r)$$

hydrogen atom like
solution

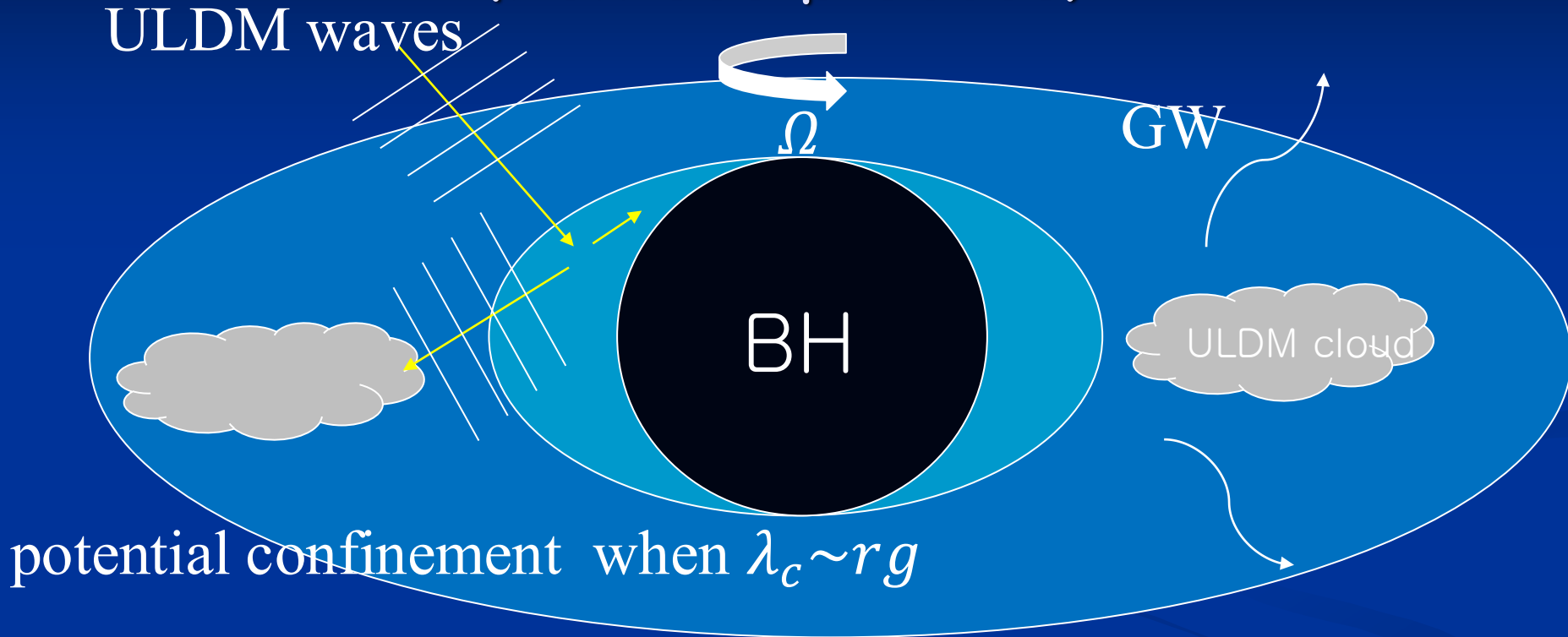
$$\psi_{nlm}(t, r) = R_{nl}(r) Y_{lm}(\theta, \phi) e^{i(\mu - \omega_{nlm})t}$$

gravitational
fine structure constant

$$\alpha \equiv \frac{r_g}{\lambda_c} = \frac{GM\mu}{\hbar c} \approx 0.1 \left(\frac{M}{10^{11} M_S} \right) \left(\frac{\mu}{10^{-22} \text{eV}} \right) \ll 1$$

Occupation # is huge, presence of horizon

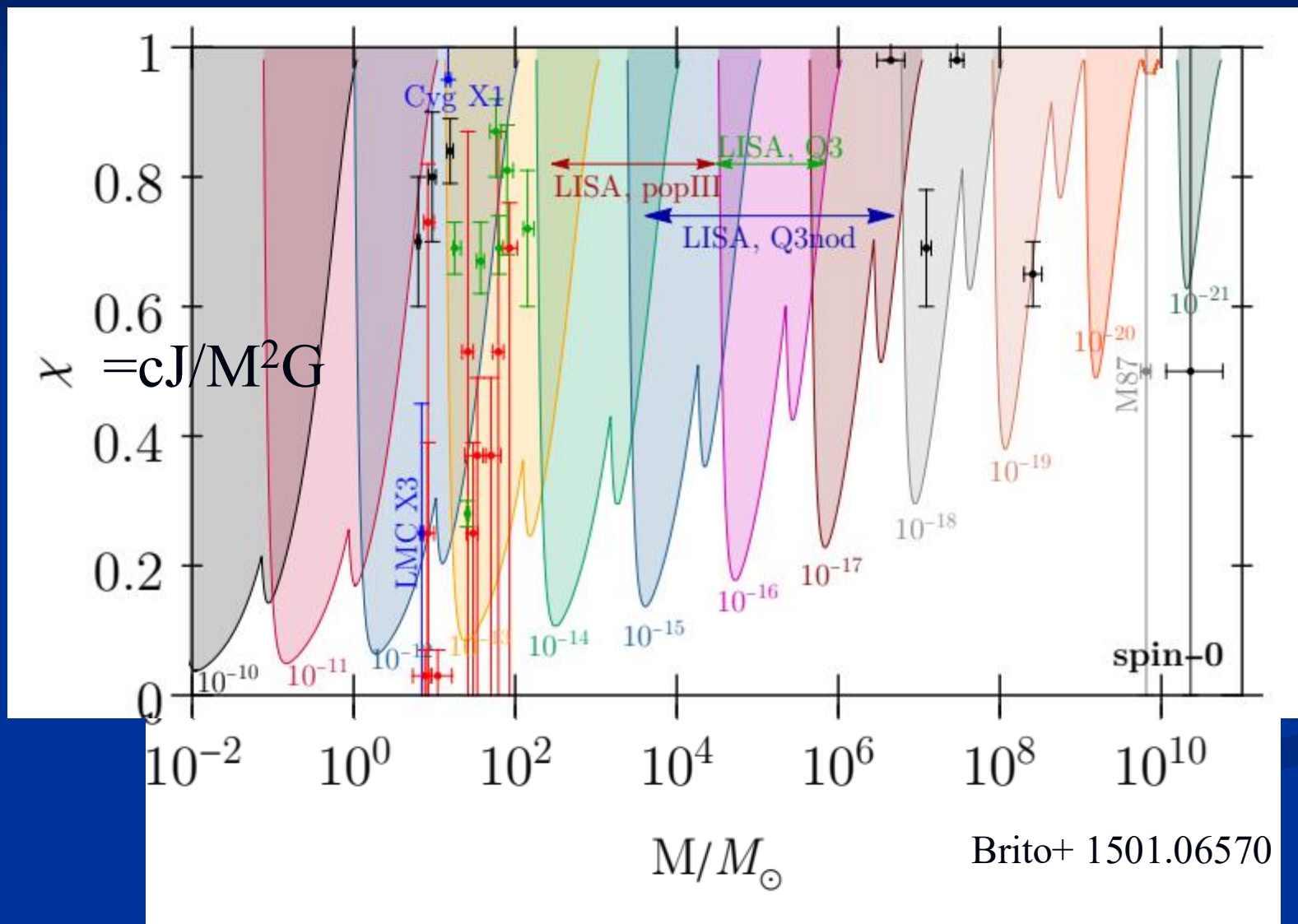
Superradiance (Penrose process)



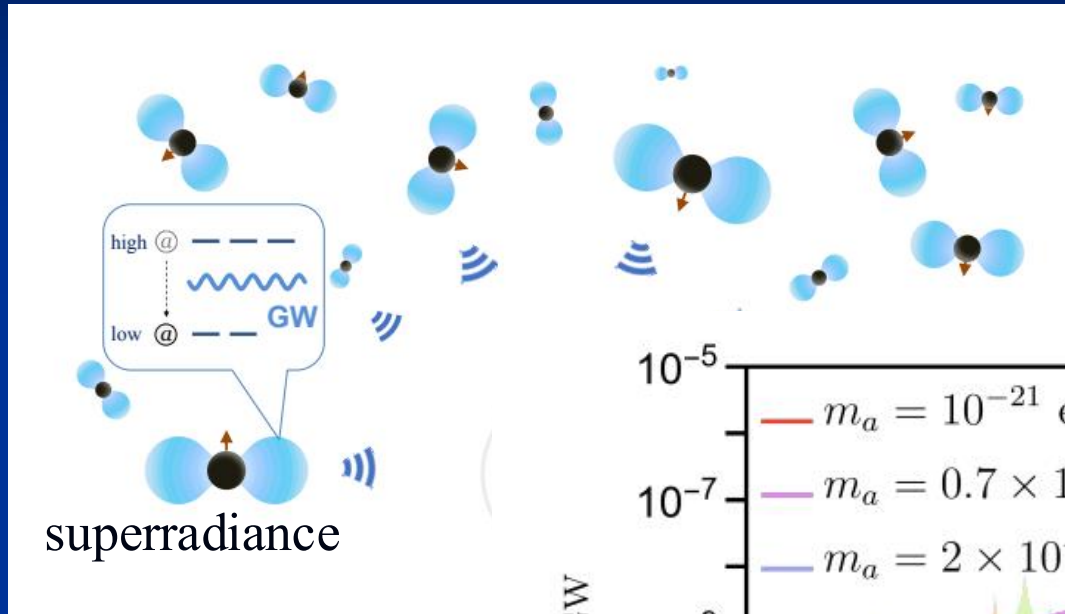
- growing mode if $\omega_{nlm} < m\Omega \rightarrow$ BH spin decreases
magnetic q. number
- can change GW patterns from a BH and BH binary

See Zhang & Yang 2018, 1908.10370, 2208.06408

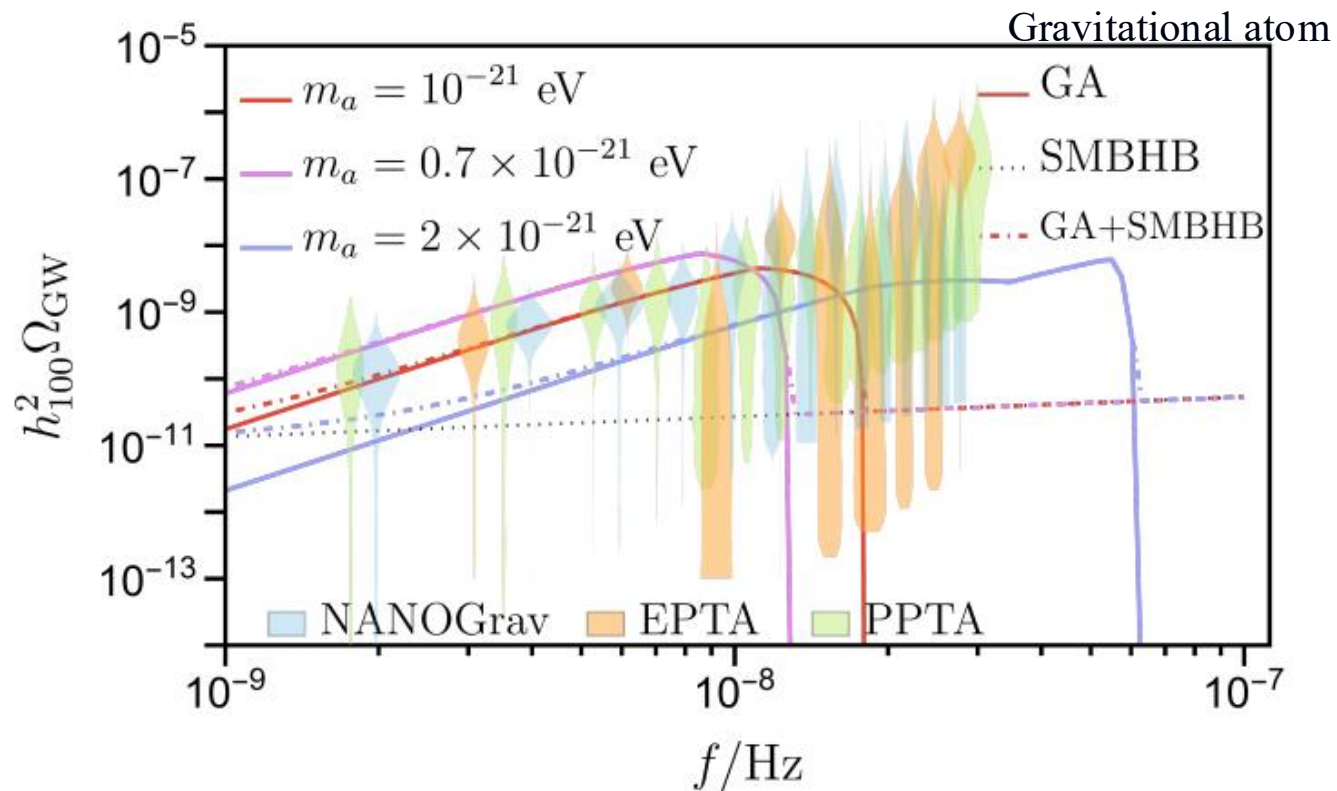
Bounds from BH-spin measurements (Regge plane)

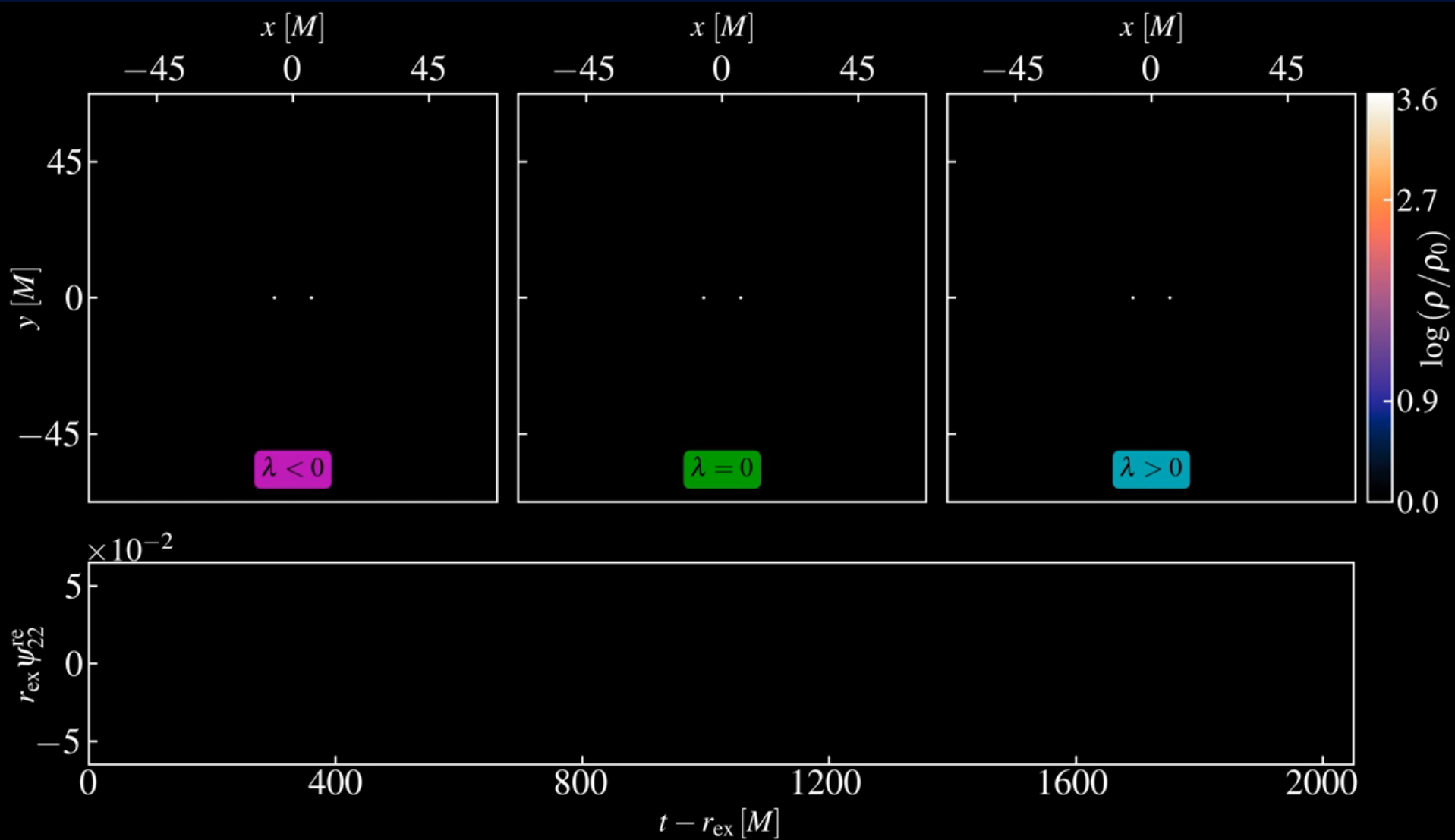


nano-Hertz stochastic GW background



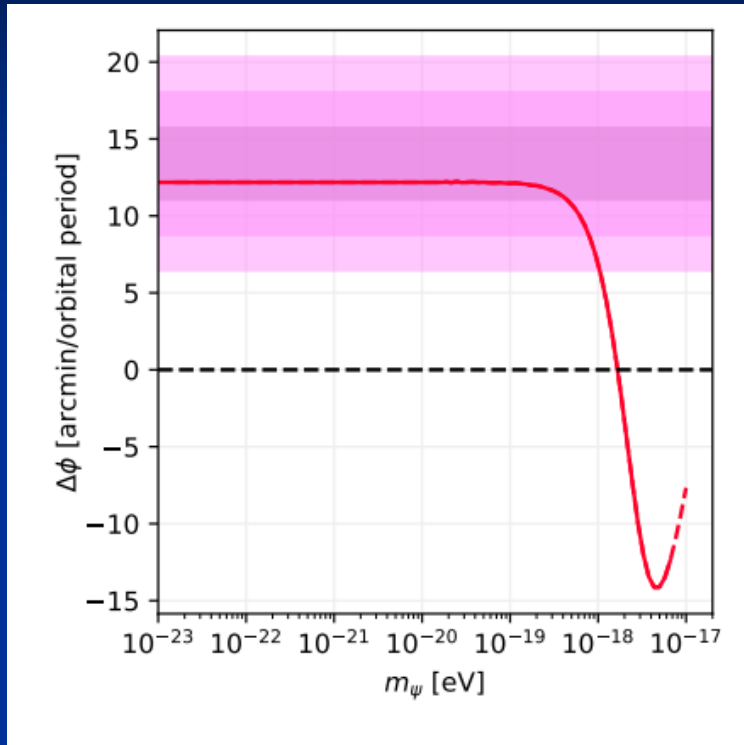
Yang+ 2306.17113





S2

Monica+ 2305.10242

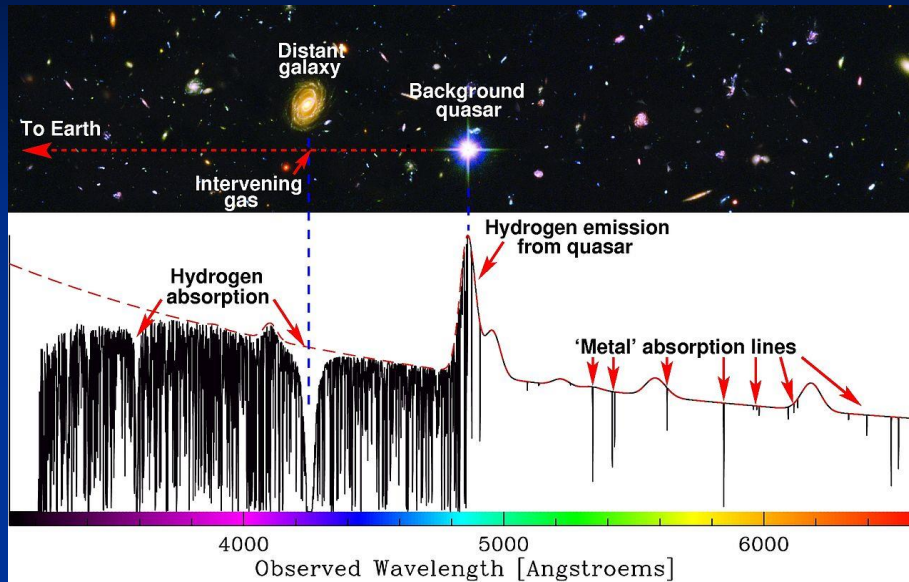


$$ds^2 = -F(r)dt^2 + \frac{1}{F(r)}dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

$$m_\psi \lesssim 3.2 \times 10^{-19} \text{ eV}$$

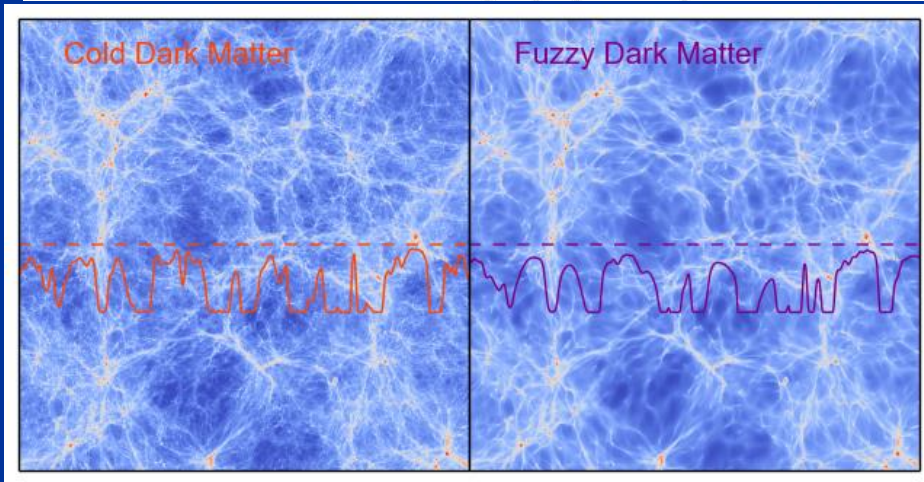
$$F(r) = \exp\left(\frac{4\pi\rho_c r_c^3}{91\alpha^8 r} \left(\left(\frac{r_c}{r}\right)^{13} \mathcal{F} - \frac{3003\pi\alpha^{\frac{13}{2}}}{2048}\right)\right) - \frac{2M}{r},$$

Lyman alpha tension?



$$m > 10^{-21} \text{eV}$$

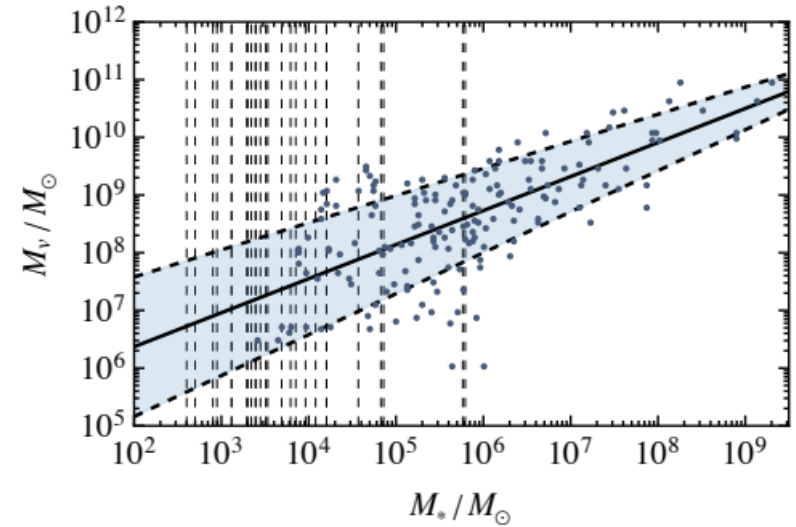
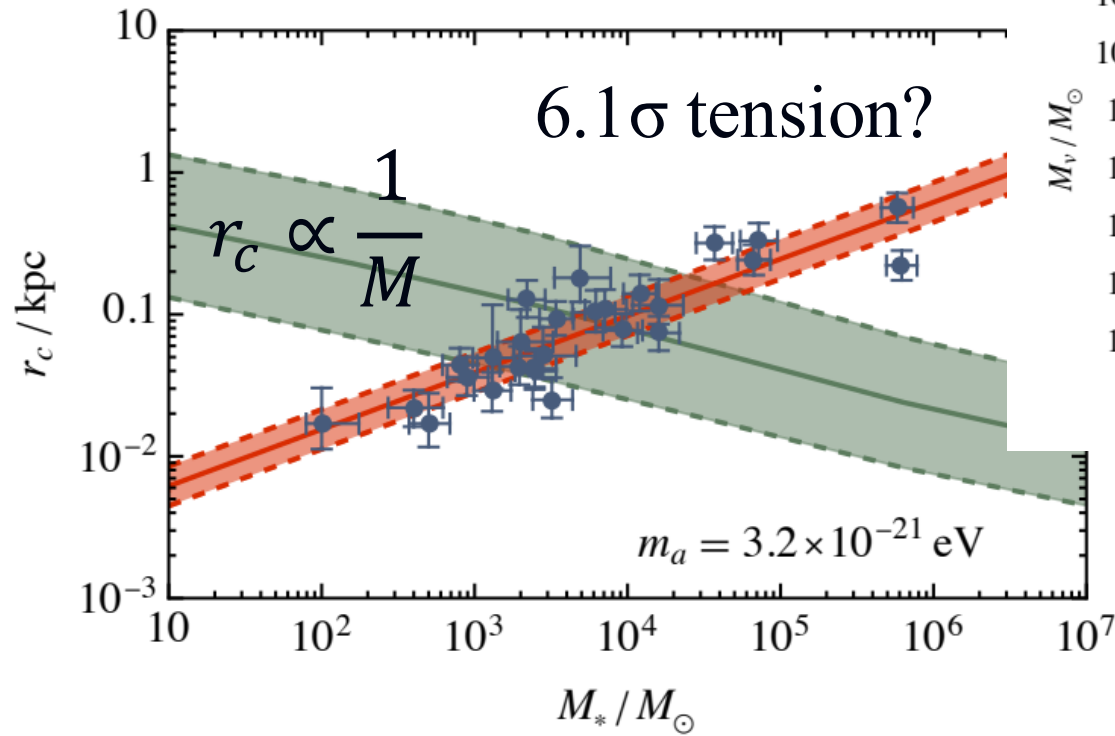
PRL 2017 (Irsic et al)



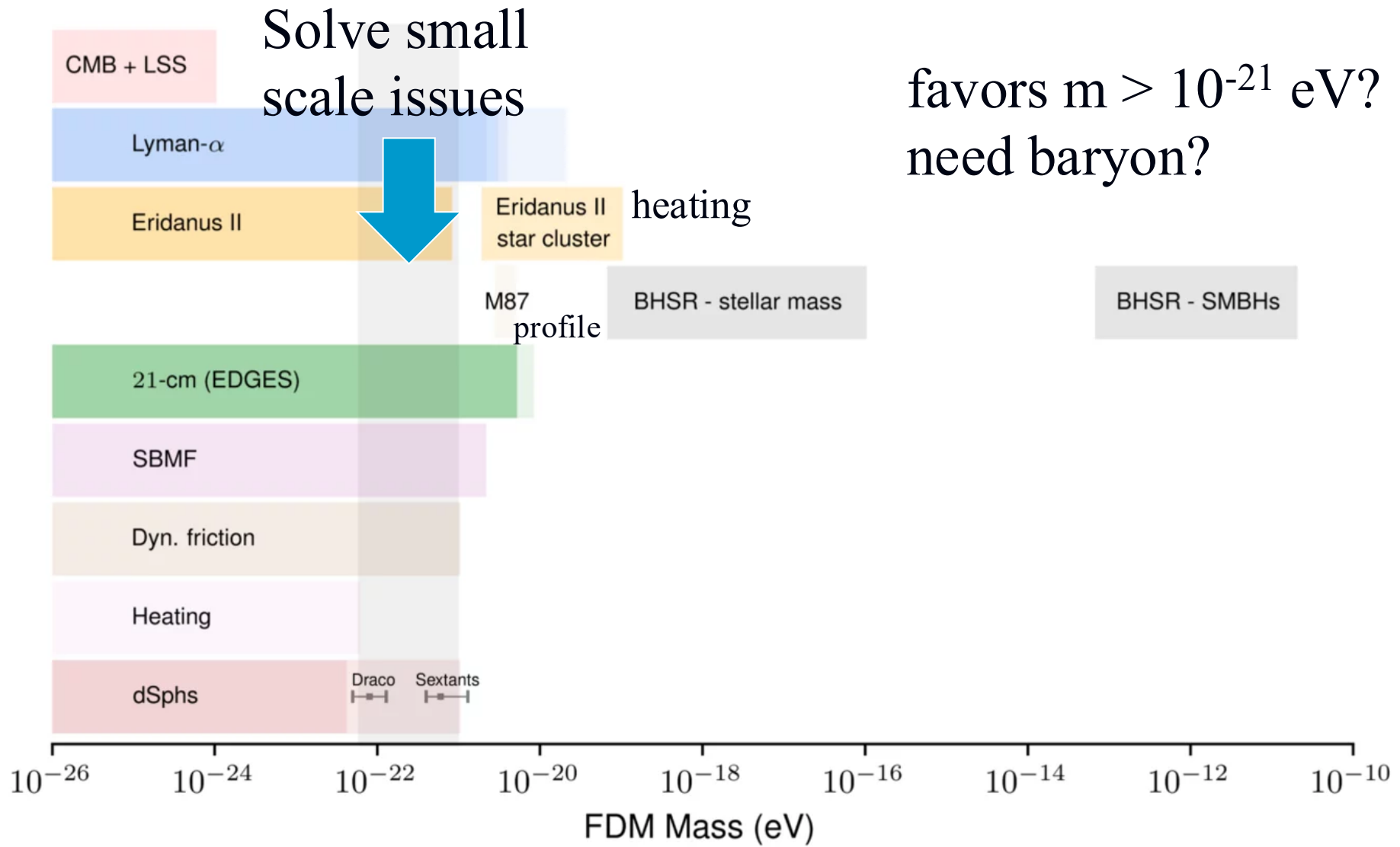
Hydrosimulation uncertainty
is large

$$\text{WDM (1, 3.3) keV} \sim \text{FDM (1, 20) } \times 10^{-22} \text{eV}$$

UFD constraints



Constraints on FDM (free ULDM) mass



Merits of studying self-interacting ULDM

We can

- allow wider mass range
→ avoid some tensions of FDM
- study direct, or indirect detection of ULDM
- calculate abundance
- understand particle model
- explain other mysteries like Hubble tension & EW scale ...

Constraints on **particle** SIDM

Positive observations	σ/m	v_{rel}	Observation	Refs.
Cores in spiral galaxies (dwarf/LSB galaxies)	$\gtrsim 1 \text{ cm}^2/\text{g}$	30 – 200 km/s	Rotation curves	[102, 116]
Too-big-to-fail problem				
Milky Way	$\gtrsim 0.6 \text{ cm}^2/\text{g}$	50 km/s	Stellar dispersion	[110]
Local Group	$\gtrsim 0.5 \text{ cm}^2/\text{g}$	50 km/s	Stellar dispersion	[111]
Cores in clusters	$\sim 0.1 \text{ cm}^2/\text{g}$	1500 km/s	Stellar dispersion, lensing	[116, 126]
<i>Abell 3827 subhalo merger</i>	$\sim 1.5 \text{ cm}^2/\text{g}$	1500 km/s	DM-galaxy offset	[127]
<i>Abell 520 cluster merger</i>	$\sim 1 \text{ cm}^2/\text{g}$	2000 – 3000 km/s	DM-galaxy offset	[128, 129, 130]
Constraints				
Halo shapes/ellipticity	$\lesssim 1 \text{ cm}^2/\text{g}$	1300 km/s	Cluster lensing surveys	[95]
Substructure mergers	$\lesssim 2 \text{ cm}^2/\text{g}$	$\sim 500 - 4000 \text{ km/s}$	DM-galaxy offset	[115, 131]
Merging clusters	$\lesssim \text{few cm}^2/\text{g}$	2000 – 4000 km/s	Post-merger halo survival (Scattering depth $\tau < 1$)	Table II
<i>Bullet Cluster</i>	$\lesssim 0.7 \text{ cm}^2/\text{g}$	4000 km/s	Mass-to-light ratio	[106]

Self-Interacting ULDM (SIULDM, SFDM)

Lee and Koh (PRD 53, hep-ph/9507385)

Galactic DM halo is described by **coherent scalar field with self-interaction**

Action
$$S = \int \sqrt{-g} d^4x \left[\frac{-R}{16\pi G} - \frac{g^{\mu\nu}}{2} \phi_{;\mu}^* \phi_{;\nu} - \frac{m^2}{2} |\phi|^2 - \frac{\lambda}{4} |\phi|^4 \right]$$

Metric
$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2 d\Omega \text{ Spherical.}$$

Field
$$\phi(r, t) = (4\pi G)^{-\frac{1}{2}} \sigma(r) e^{-i\omega t}$$

$$\tilde{m} \equiv \frac{m}{\lambda^{\frac{1}{4}}}$$

Exact
ground
state
(soliton)

$$\sigma_* = \sqrt{\frac{\gamma_0 \sin(\sqrt{2}r_*)}{\sqrt{2}r_*}}$$

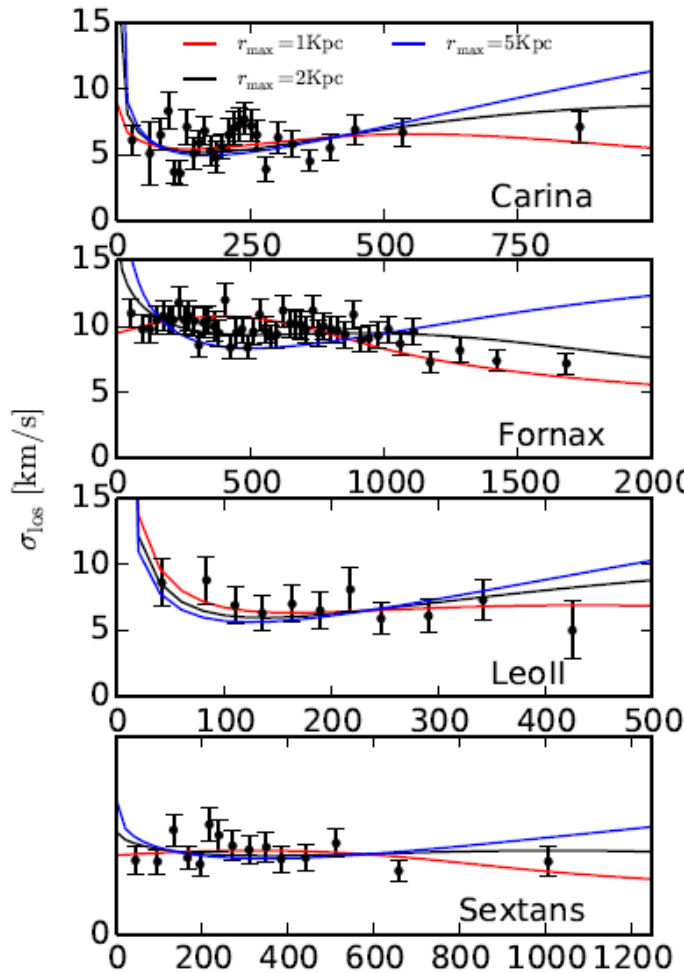
$$r_* = r\Lambda^{-1/2}$$

$$\Lambda \equiv \frac{\lambda m_P^2}{4\pi m^2}, \quad \Lambda \gg 1 \quad (\text{Newtonian \& TF limit})$$

$$\text{New constant length scale } RTF \approx \frac{\sqrt{\Lambda}}{m} = \sqrt{\frac{\pi \hbar^3 \lambda}{8cGm^4}} \sim \tilde{m}^{-2}$$

$$\& \text{ mass scale } M_{\max} = \sqrt{\Lambda} \frac{m_P^2}{m}$$

- Even tiny self-interaction drastically changes the scales!
- allows wider range for m to fit observations
- $\tilde{m}$ determines all scales



$$\lambda^{1/2} \left(\frac{m_P}{m} \right)^2 \geq 10^{50} \rightarrow \left(\frac{m_P}{10^{25}} \right) \approx 1221 eV \geq \frac{m}{\lambda^{1/4}}$$

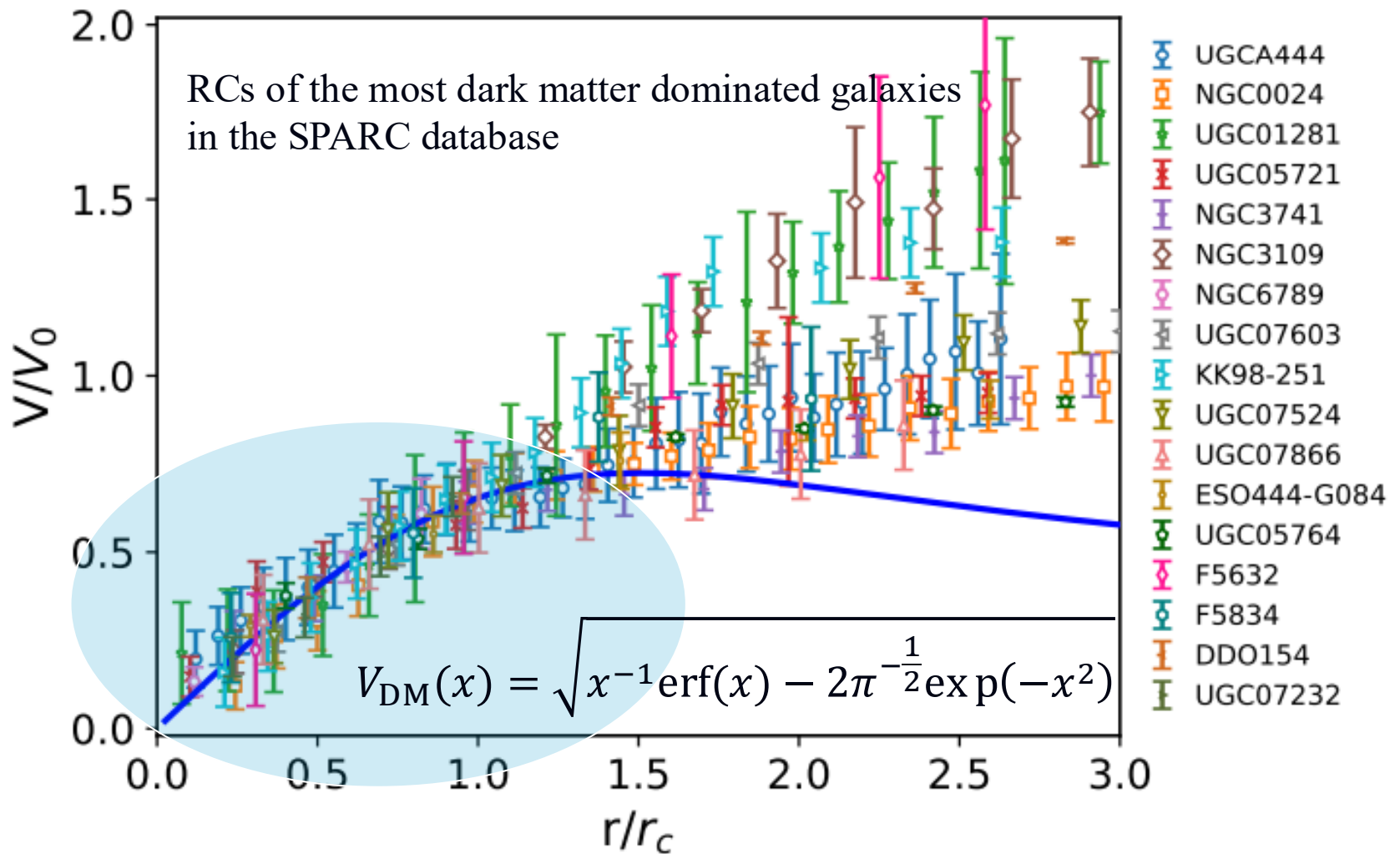
$$\rho = \sigma^2 = \rho_0 \frac{\sin(\pi r / r_{\max})}{\pi r / r_{\max}}$$

$$r_{\max} \equiv \sqrt{\frac{\pi^2 \Lambda}{2}} \square 49 \left(\frac{\lambda^{1/4}}{m / eV} \right)^2 kpc$$

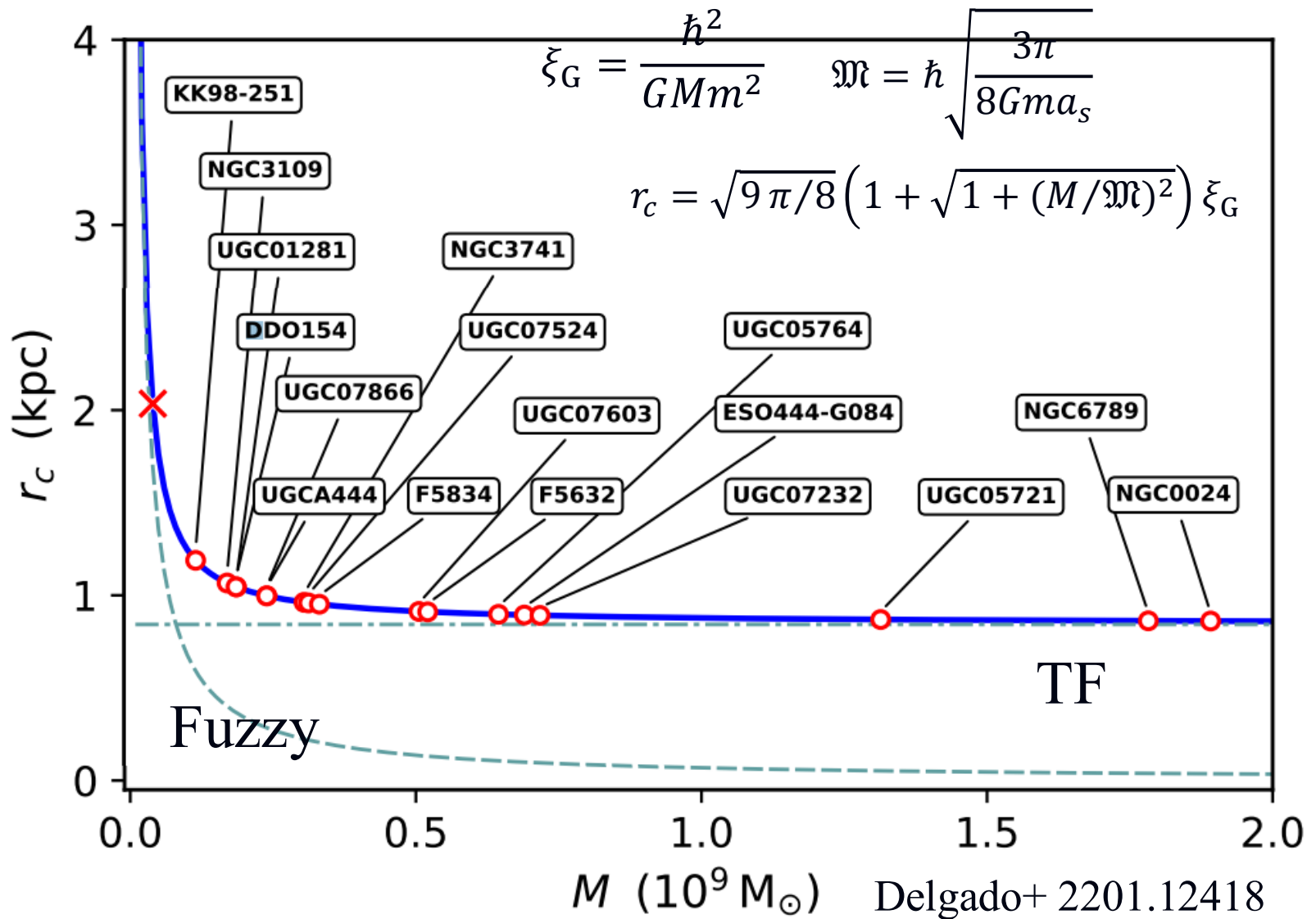
$$\text{dSphs} \quad r_{\max} \square 1 kpc \rightarrow m / \lambda^{1/4} \square 7 eV$$

$$3 \text{ eV} < \tilde{m} = \frac{m}{\lambda^{1/4}} < 7 \text{ eV}$$

Universal RC with SIULDM



$$x = r/r_c$$



best fit $m \simeq 2.2 \times 10^{-22} \text{ eV}$, $a_s = \frac{\lambda}{8\pi m} \simeq 7.8 \times 10^{-77} \text{ meter}$

$\rightarrow \tilde{m} = 5.72 \text{ eV}$

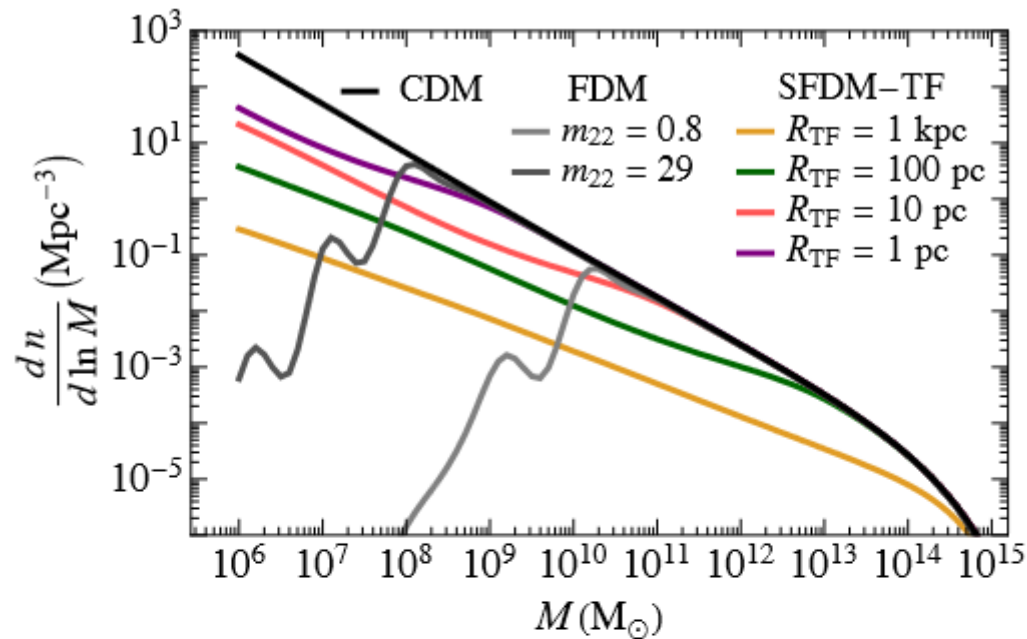
Timing problem of GC

$$\bar{\tau} = \frac{3}{4\pi\rho_c G^2 M_{GC}} \times \begin{cases} c_s^3 & (\text{Thomas Fermi limit}) \\ \frac{1}{\sqrt{2}} \left(\frac{\hbar v_0}{m_\phi r_0} \right)^{3/2} & (\text{FDM limit}). \end{cases}$$

Koo & Lee
JCAP 2026

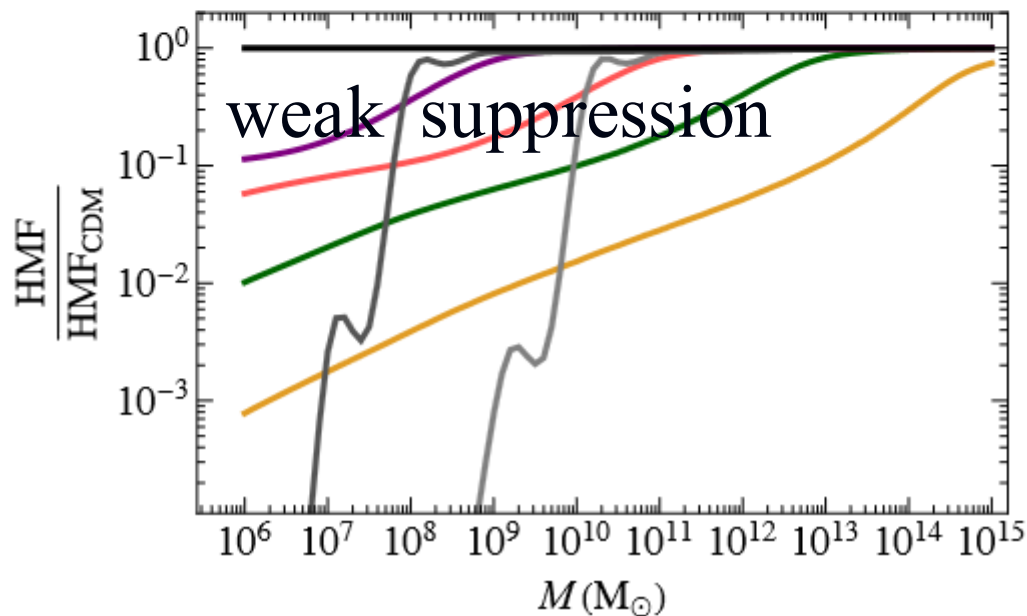
$$1.17 \times 10^{-22} \text{eV} \leq m_\phi \leq 4.68 \times 10^{-22} \text{eV} \quad (\text{FDM})$$

$$3.75 \text{eV} \leq \frac{m_\phi}{\lambda^{1/4}} \leq 5.57 \text{eV} \quad (\text{GC3, TF})$$



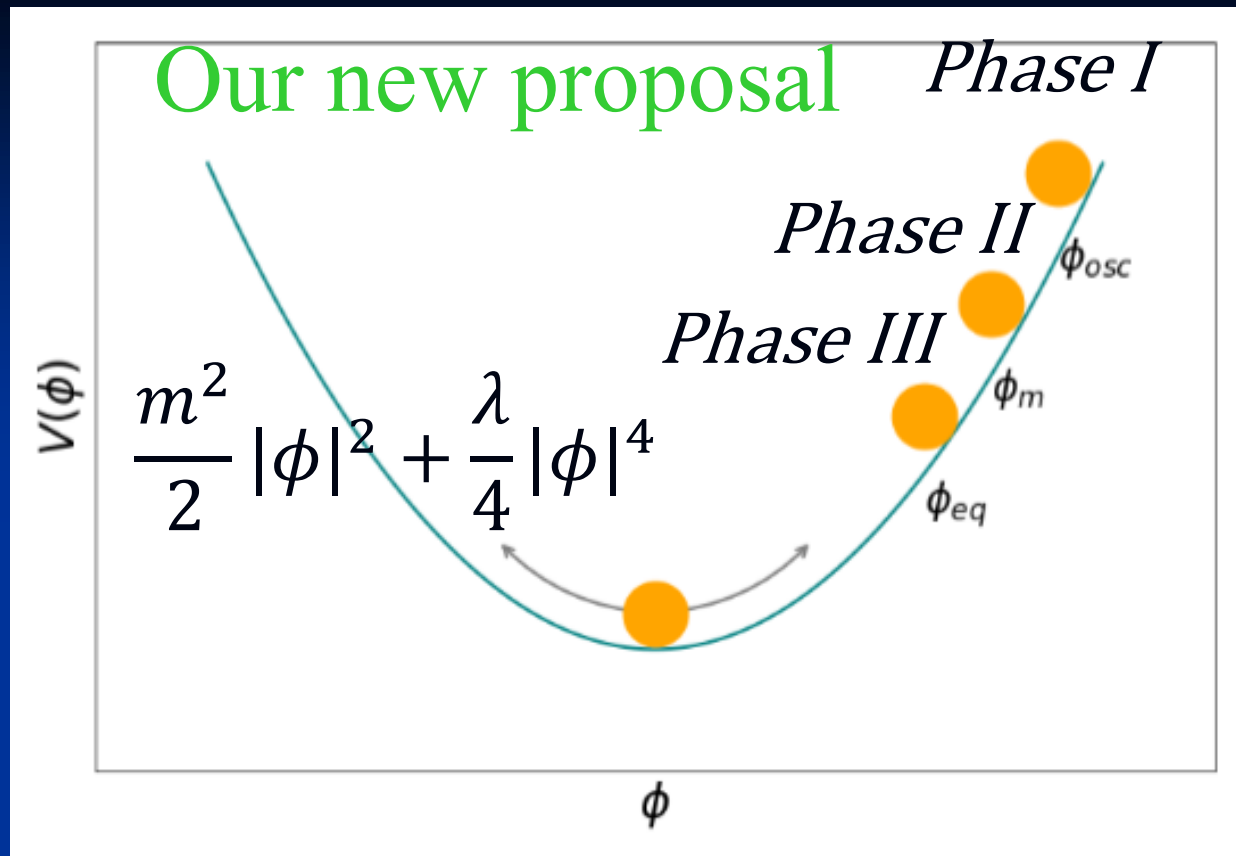
λ_J fixed

$$M_J \propto a^{-3}$$



weak suppression
 $\rightarrow$ weak missing satellites

$\rightarrow$ may avoid Lyman alpha tension

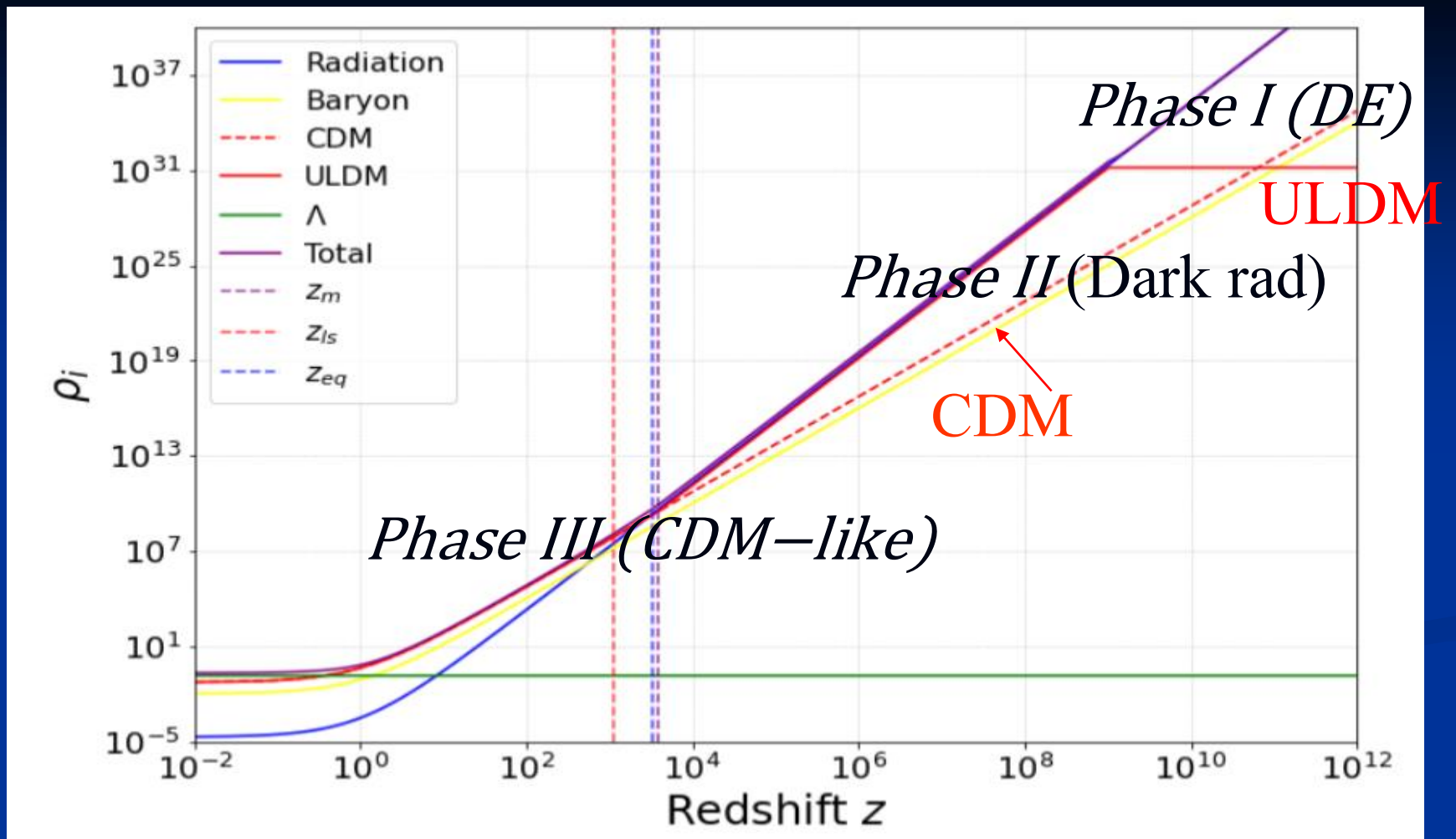


JLee
2502.11568

$V(\phi) \propto \phi^n$ decays with an equation of state $w = \frac{(n-2)}{(n+2)}$

$\left\{ \begin{array}{l} \text{DE-like (slow-roll) for } z_{osc} < z \text{ (Phase I)} \\ \text{Dark rad-like } (\phi^4 > \phi^2) \text{ or } z_m < z < z_{osc} \text{ (Phase II)} \\ \text{CDM } (\phi^4 < \phi^2) \text{ or } z < z_m \text{ (Phase III)} \end{array} \right.$ or stiff

SIULDM alone is enough!



- SIULDM acts as a natural extra radiation comp. just before ls
- We don't need exotic matter or fine tuning?

Cosmological evolution

$$V(\psi) = \frac{1}{2}mc^2|\psi|^2 + \frac{\lambda}{2}|\psi|^4$$

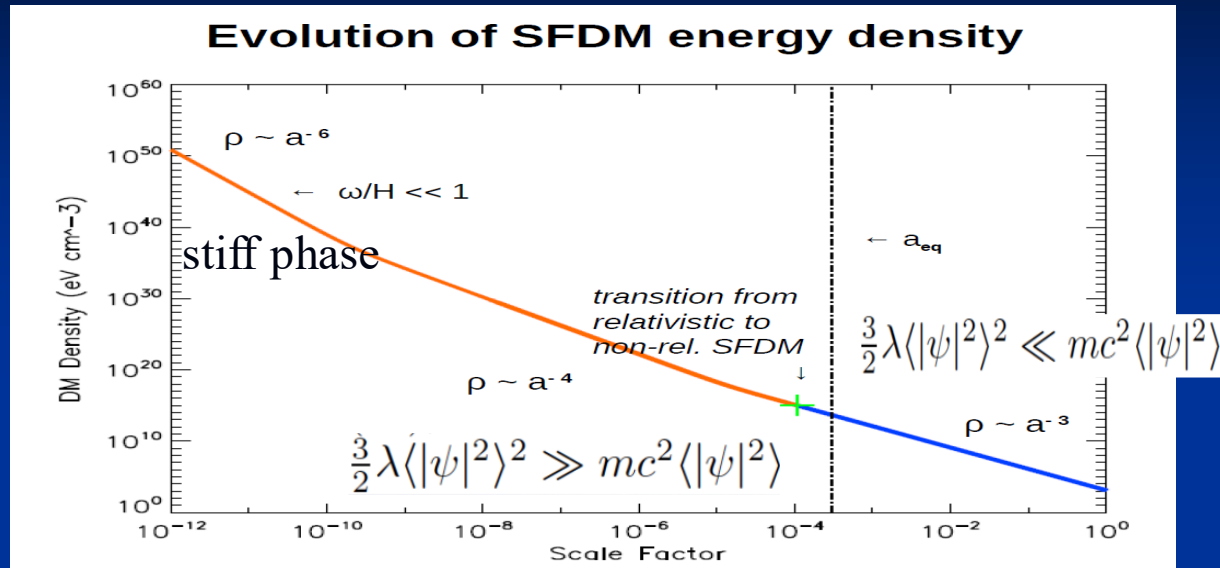
$$\langle \bar{w} \rangle \equiv \frac{\langle \bar{p} \rangle}{\langle \bar{\rho} \rangle} = \frac{1}{3} \left[\frac{1}{1 + \frac{2mc^2}{3\lambda \langle |\psi|^2 \rangle}} \right]$$

Time varying Eq. of state

$$1) z_{eq}=3390$$

$$\frac{\lambda}{(mc^2)^2} \leq 4 \times 10^{-17} \text{ eV}^{-1} \text{ cm}^3.$$

2) Big Bang
Nucleosynthesis



$$\Rightarrow m^4/\lambda > 4.7 \times 10^{-7} \text{ eV}^4$$

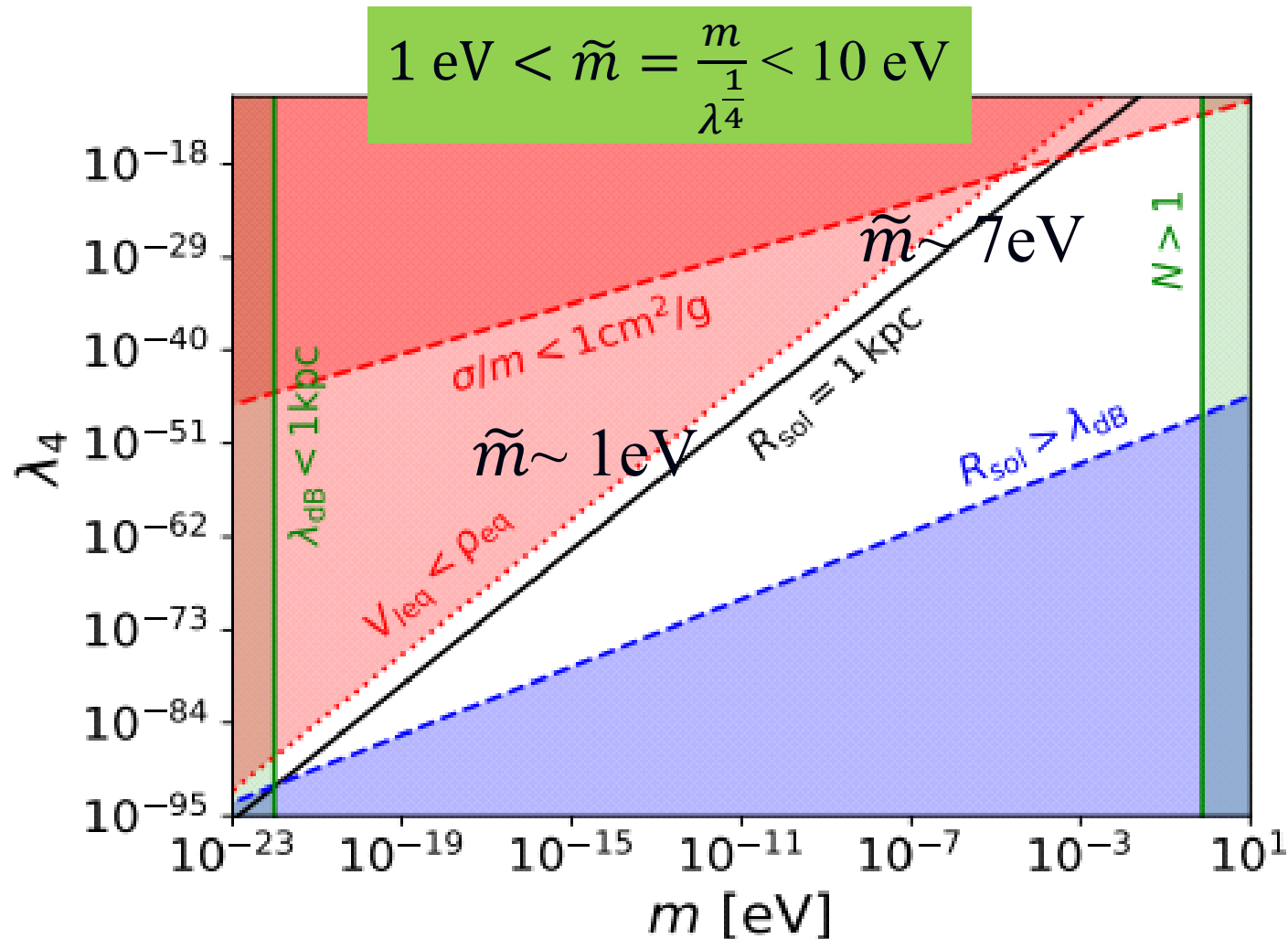
$$\frac{\Delta N_{eff}}{N_{eff,standard}} = \frac{3.71^{+0.47}_{-0.45} - 3.046}{3.046} = \frac{\rho_{SFDM}}{\rho_\nu}$$

$$9.5 \times 10^{-19} \text{ eV}^{-1} \text{ cm}^3 \leq \lambda/(mc^2)^2 \leq 1.5 \times 10^{-16} \text{ eV}^{-1} \text{ cm}^3$$

$$\rightarrow 2.6 \text{ eV} \leq \frac{m}{\lambda^{1/4}} \leq 9.5 \text{ eV}$$

cosmological constraints

Garcia + 2304.10221



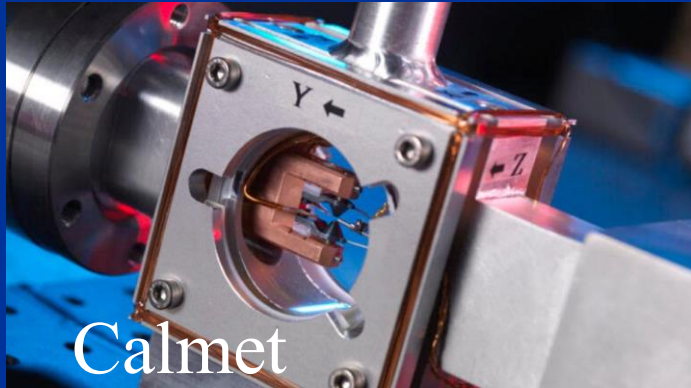
Detection of interacting ULDM

dilatonic coupling Using atomic clocks to detect ULDM by mimicking time variations of fundamental constants

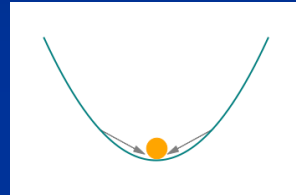
$$\mathcal{L} \supset \varphi \frac{d_e}{4\mu_0} F_{\mu\nu} F^{\mu\nu},$$



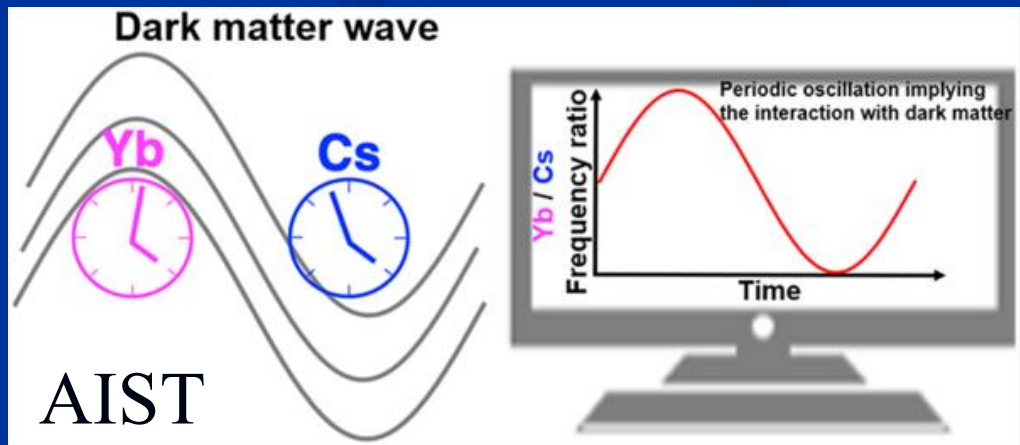
$$\alpha(t) \approx \alpha [1 + d_e \varphi_0 \cos(\omega t + \delta)]$$



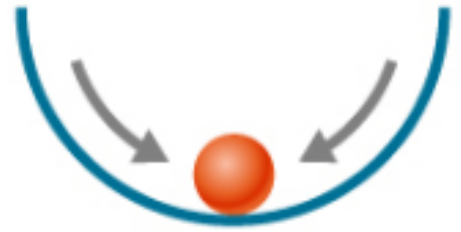
Oscillation of fine structure constant



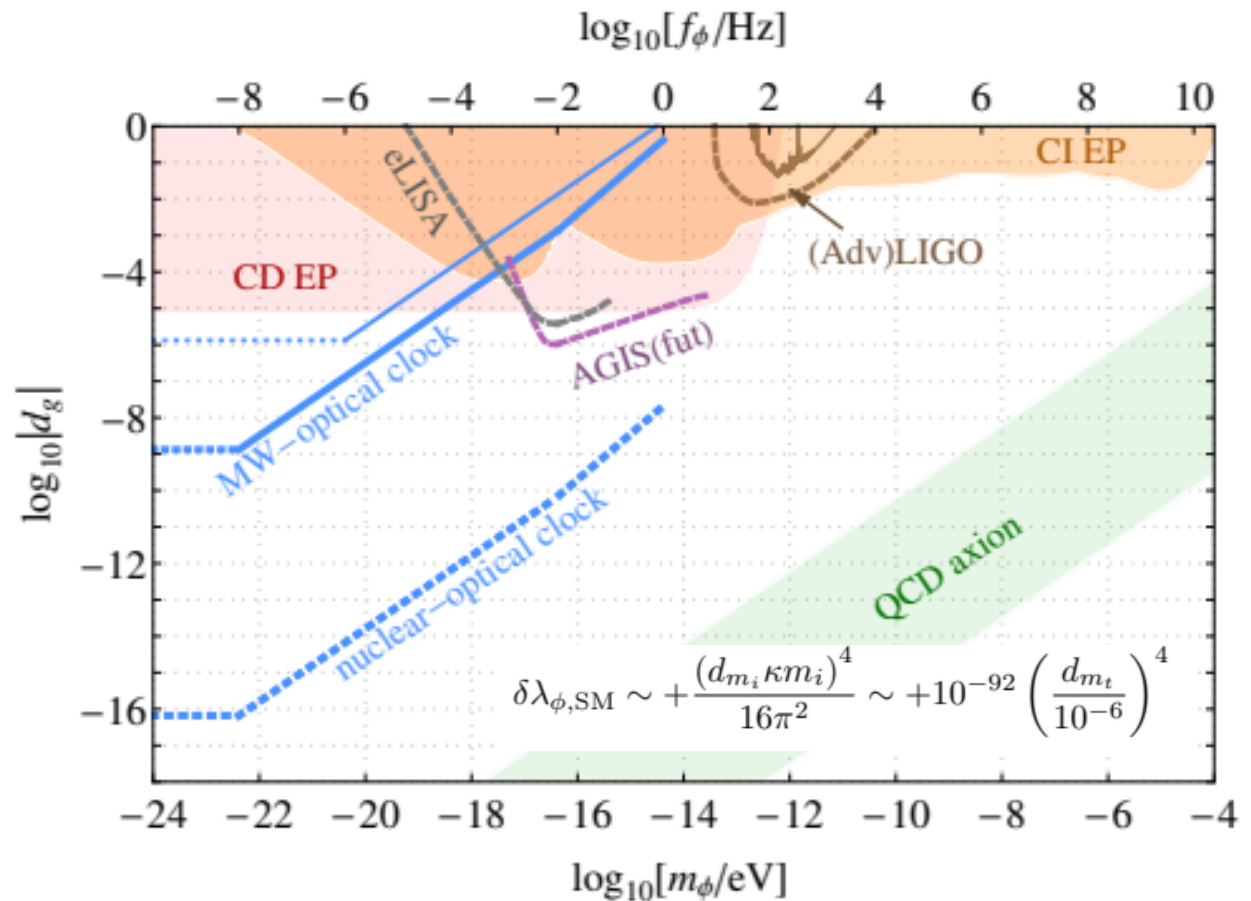
Due to nuclear and atomic structure Yb and Cs have different frequency dependency on α (special relativistic effects)



$$\omega = \frac{1}{2.5 \text{ months}} \frac{m}{10^{-22} \text{ eV}}$$



$$\frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu} + \frac{\varphi d_e}{4g^2} F_{\mu\nu} F^{\mu\nu} \Rightarrow \alpha(t) \approx \alpha [1 + d_e \varphi_0 \cos(\omega t + \delta)]$$



Direct detection

* The common low-energy effective interaction is

$$\mathcal{L}_\phi = \kappa \phi \mathcal{O}_{\text{SM}},$$

where $\kappa \equiv \sqrt{4\pi}/M_{\text{Pl}}$ and

$$\mathcal{O}_{\text{SM}} = d_e \frac{F_{\mu\nu} F^{\mu\nu}}{4} - d_g \frac{\beta_3}{2g_3} G_{\mu\nu}^A G^{A\mu\nu} - d_{me} m_e \bar{e}e - \sum_i d_{m_i} m_i \bar{\psi}_i \psi_i$$

1) Atomic, Molecular, Nuclear Clocks

Detect tiny oscillations of the fine-structure constant α and electron-proton mass ratio.

Frequency ratios of different clocks vary as $\phi(t) \propto \cos(m_\phi t)$.

Effective Lagrangian:

$$\mathcal{L} \supset \kappa \phi \left[d_e \frac{F_{\mu\nu} F^{\mu\nu}}{4} + d_{me} m_e \bar{e}e + d_{\hat{m}} (m_u \bar{u}u + m_d \bar{d}d) \right]$$

2) Optical Interferometers (including GW detectors)

- Sensitive to oscillations of α that change optical path length or refractive index.
- Effective Lagrangian:

$$\mathcal{L} \supset \kappa \phi d_e \frac{F_{\mu\nu} F^{\mu\nu}}{4}.$$

Atom Interferometers

- Detect variations of atomic mass m_A and gravitational potential.

- Effective Lagrangian:
$$\mathcal{L} \supset \kappa \phi \left[d_{\hat{m}} \sum_q m_q \bar{q}q + d_g \frac{\beta_3}{2g_3} G_{\mu\nu}^A G^{A\mu\nu} \right].$$

Torsion Balances and Equivalence Principle Tests

- Search for fifth forces and tiny composition-dependent accelerations.
- Effective Lagrangian:

$$\mathcal{L} \supset \kappa \phi d_g \frac{\beta_3}{2g_3} G_{\mu\nu}^A G^{A\mu\nu}$$

LC Oscillators and Mechanical Resonators

- Look for oscillations of electrical or mechanical resonance frequencies.
- Effective Lagrangian:

$$\mathcal{L} \supset \kappa \phi d_e \frac{F_{\mu\nu} F^{\mu\nu}}{4}$$

Haloscopes (cavity, plasma, dielectric)

- Probe scalar or axion-like couplings to photons.
- Scalar case:

$$\mathcal{L} \supset \kappa \phi d_e \frac{F_{\mu\nu} F^{\mu\nu}}{4}$$

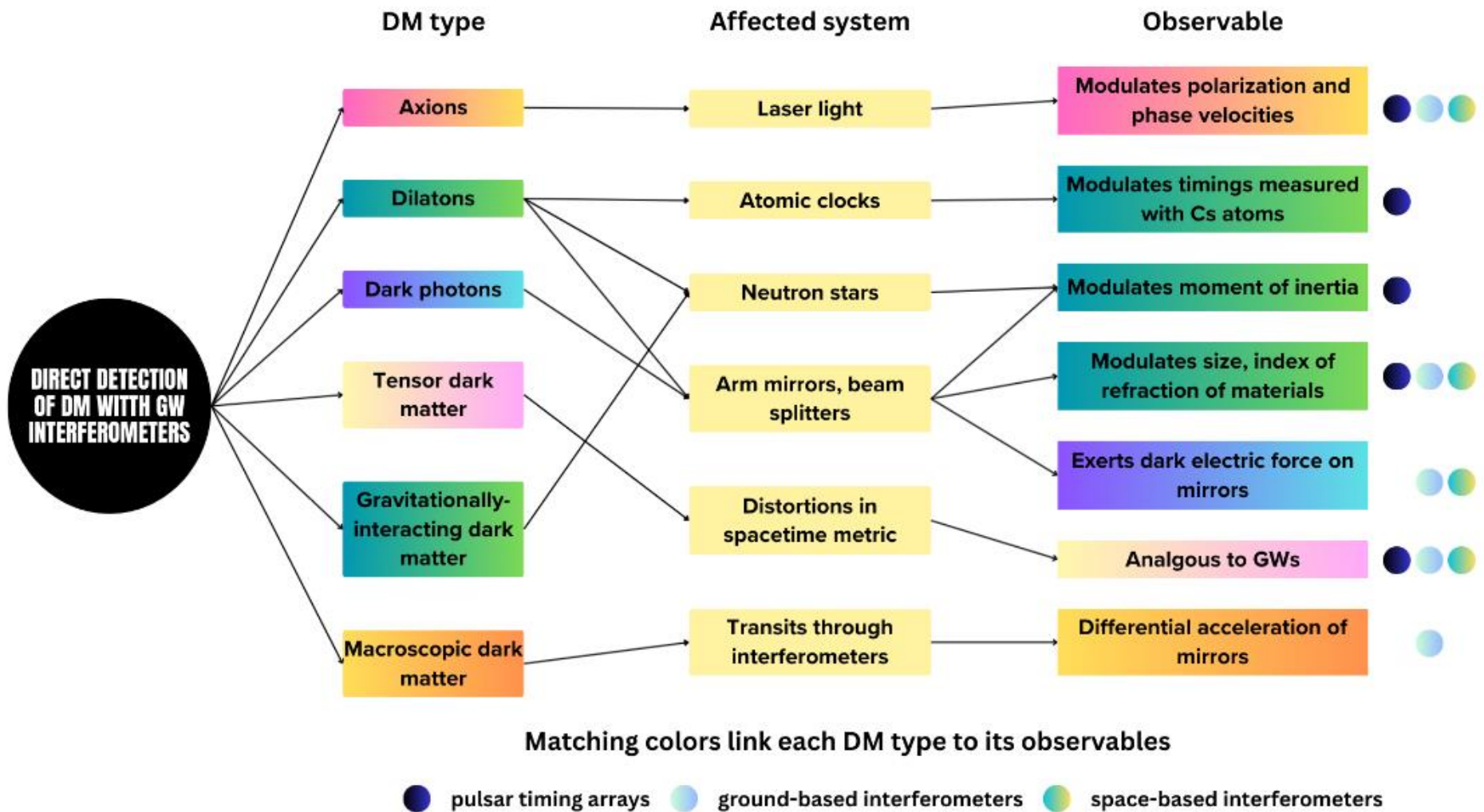
- Pseudoscalar (axion-like) case:

$$\mathcal{L} \supset \frac{g_{\phi\gamma\gamma}}{4} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}.$$

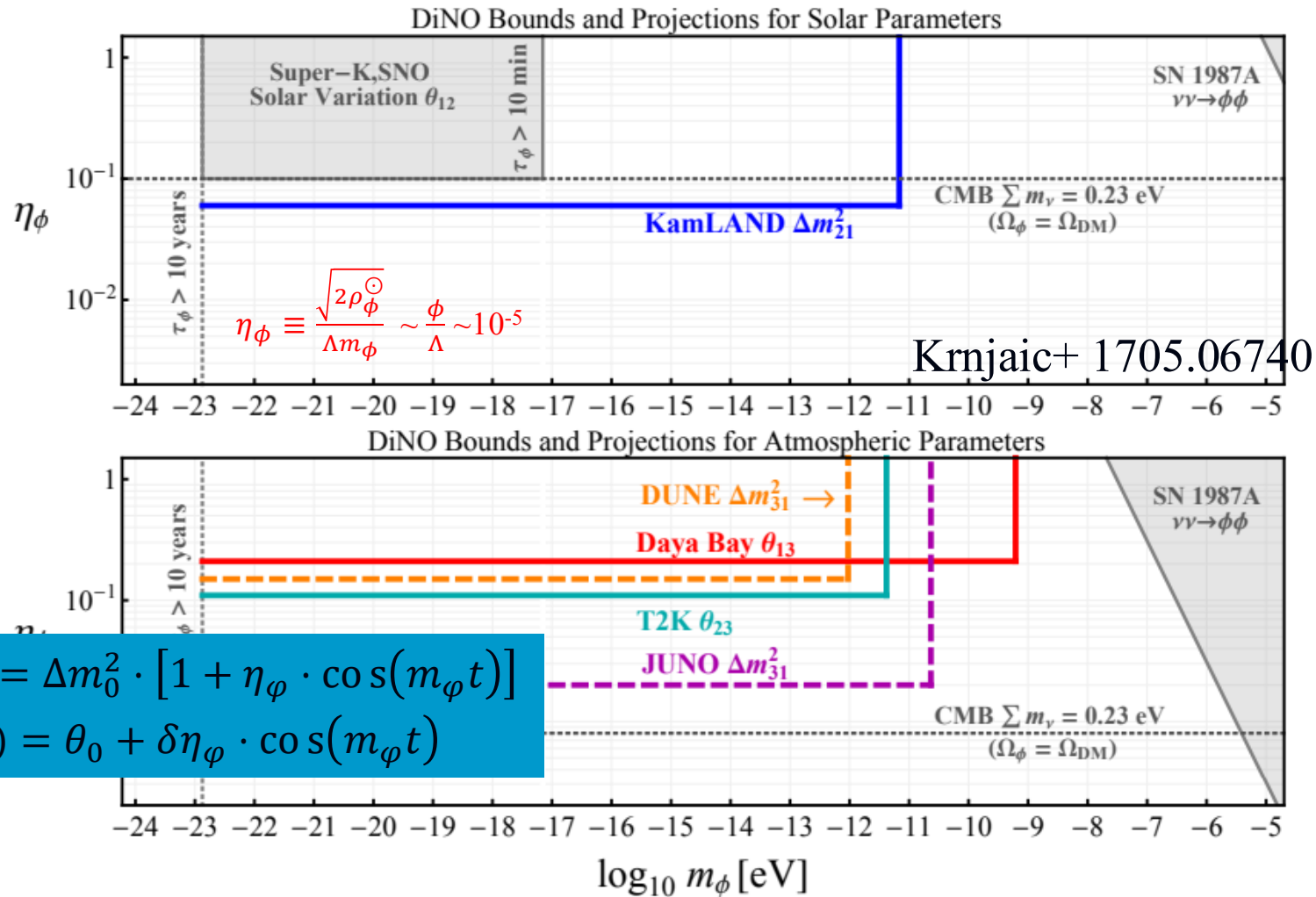
Spin-Based Sensors

- Measure shifts in Larmor precession of electron or nuclear spins.
- Effective Lagrangian:

$$\mathcal{L} \supset \kappa \phi d_{me} m_e \bar{e} e.$$



neutrino oscillation with ULDM

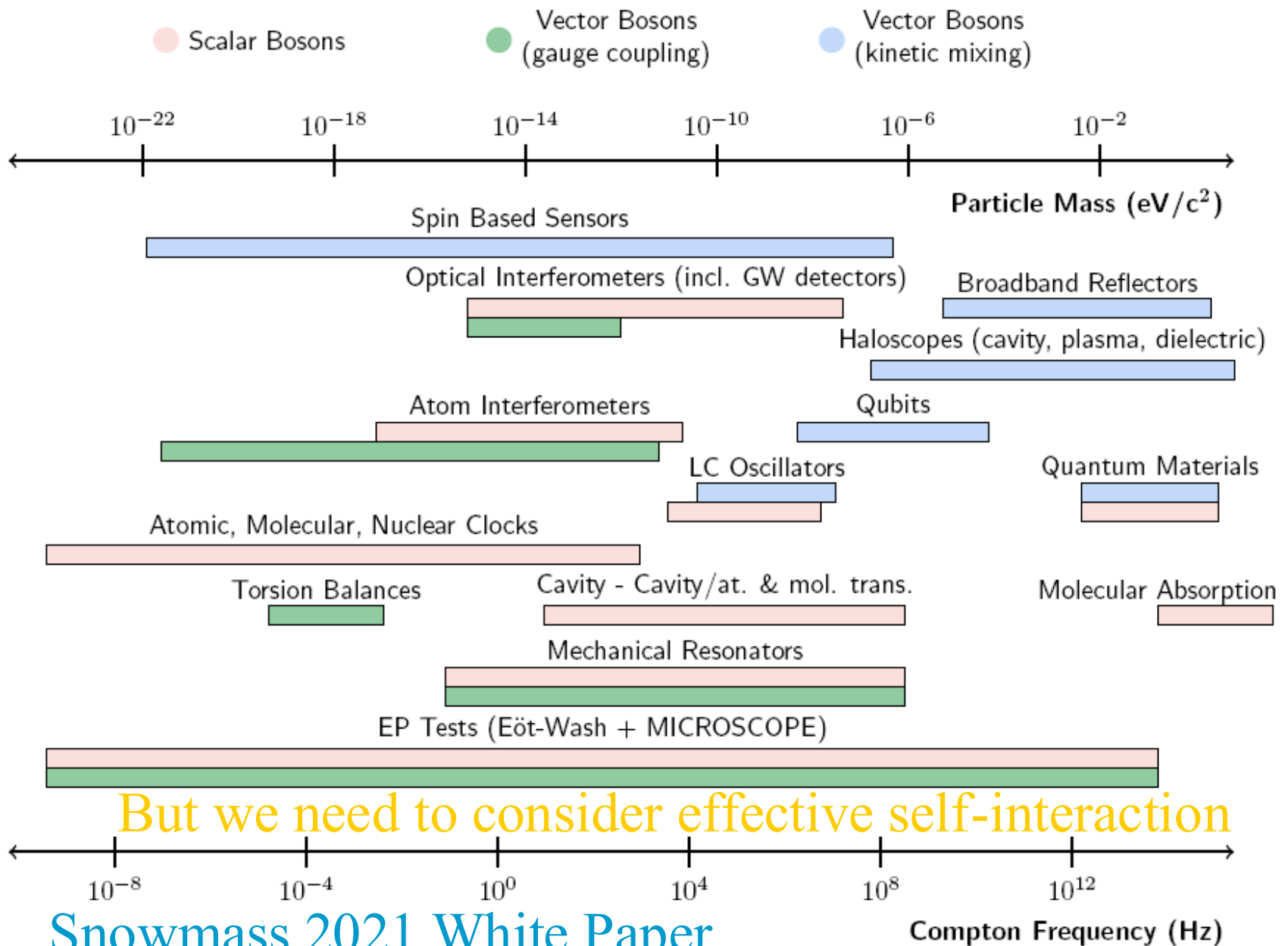


$$\Delta m^2(t) = \Delta m_0^2 \cdot [1 + \eta_\phi \cdot \cos(m_\phi t)]$$

$$\theta(t) = \theta_0 + \delta\eta_\phi \cdot \cos(m_\phi t)$$

$$\mathcal{L}_{\text{eff}} = -m_\nu \left(1 + y_1 \frac{\phi}{\Lambda} \right) \nu\nu, \quad \phi(x, t) \simeq \frac{\sqrt{2\rho_\phi^\odot}}{m_\phi} \cos[m_\phi(t - \vec{v} \cdot \vec{x})]$$

Dark Matter Candidates



Conclusions

ULDM with $m \sim 10^{-21}$ eV or

self-interacting ULDM with $\frac{m}{\lambda^{1/4}} \sim 1 \text{ eV}$

seems to be a viable alternative to CDM

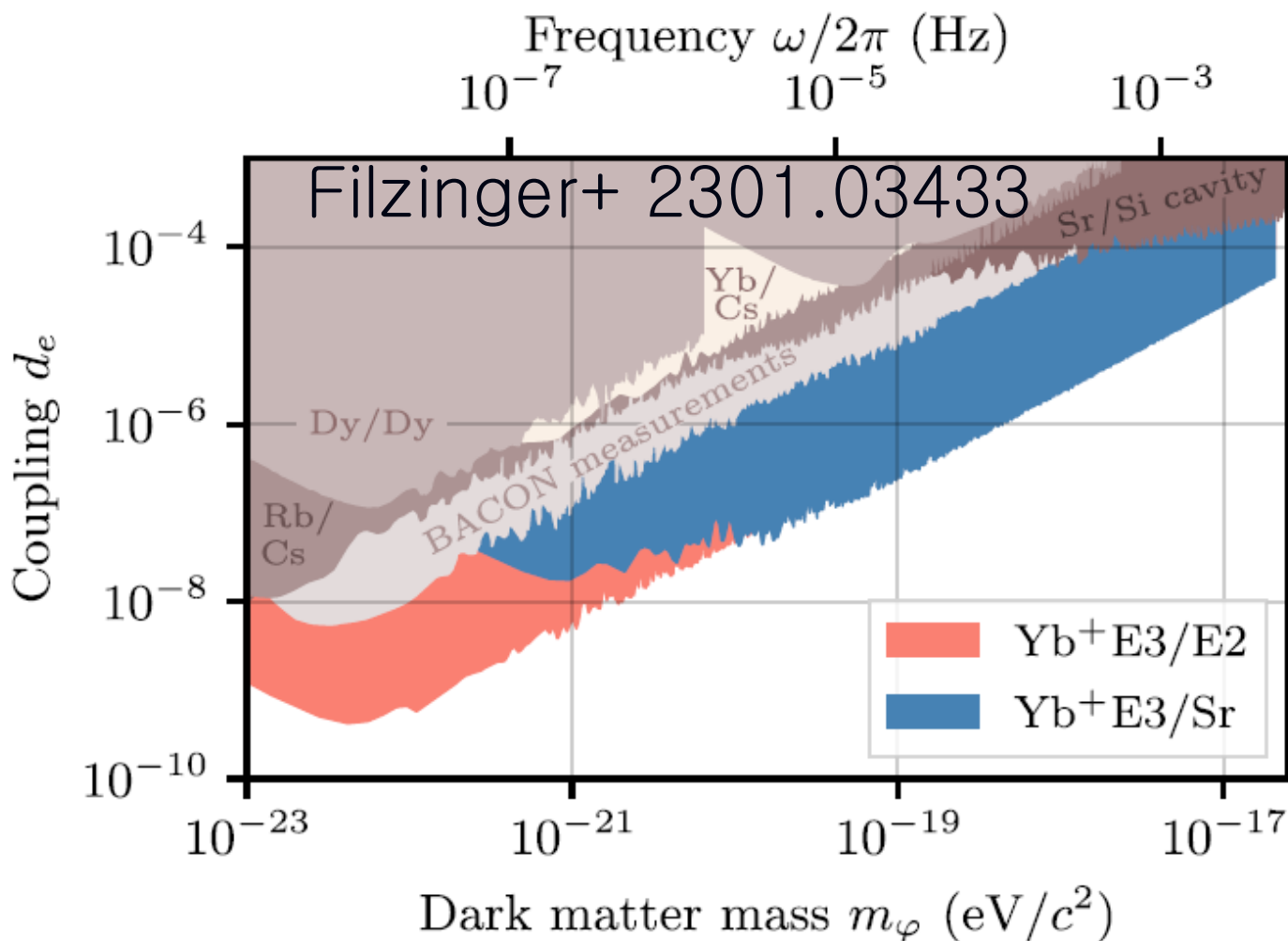
→ might solve many mysteries of cosmology and astrophysics.

→ can be detected soon!

detection

Oscillation of fine structure constant

$$\alpha(t) \approx \alpha [1 + d_e \varphi_0 \cos(\omega t + \delta)]$$



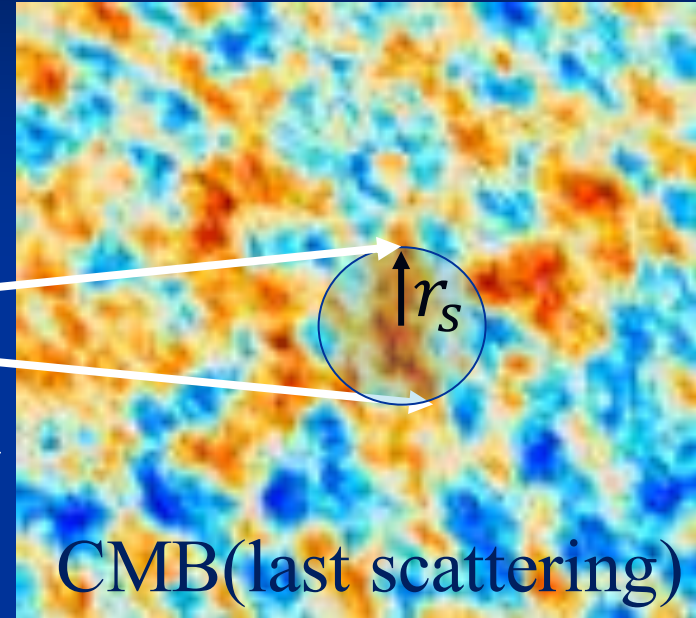
H_0 estimation from CMB

$\theta_s = \frac{r_s}{D_A}$ is a standard ruler with 0.04% precision. (fixed)



$$\theta_s \sim 1^\circ$$

$$D_A$$



Sound Horizon
$$r_s = \frac{c}{\sqrt{3}H_{ls}} \int_{z_{ls}}^{\infty} \frac{dz}{[\rho(z)/\rho(z_{ls})]^{1/2} (1 + R(z))^{1/2}}$$

Angular distance
$$D_A = \frac{c}{H_0} \int_0^{z_{ls}} \frac{dz}{[\rho(z)/\rho_0]^{1/2}} \sim \frac{1}{H_0}$$

r_s & D_A depend on expansion history (matter contents $\rho(z)$)

$$H_0 = \theta_s H_{ls} \frac{\int_{z_{ls}}^0 \frac{c dz}{[\rho(z)/\rho_0]^{1/2}}}{\int_{\infty}^{z_{ls}} \frac{c_s(z) dz}{[\rho(z)/\rho(z_{ls})]^{1/2}}}$$

late time (decrease $\rho(z)$)

early time (increase $\rho(z)$)

To increase H_0 we can ($\theta_s = \frac{r_s}{D_A}$ fixed)

1) in early time solutions (decrease r_s to decrease D_A)

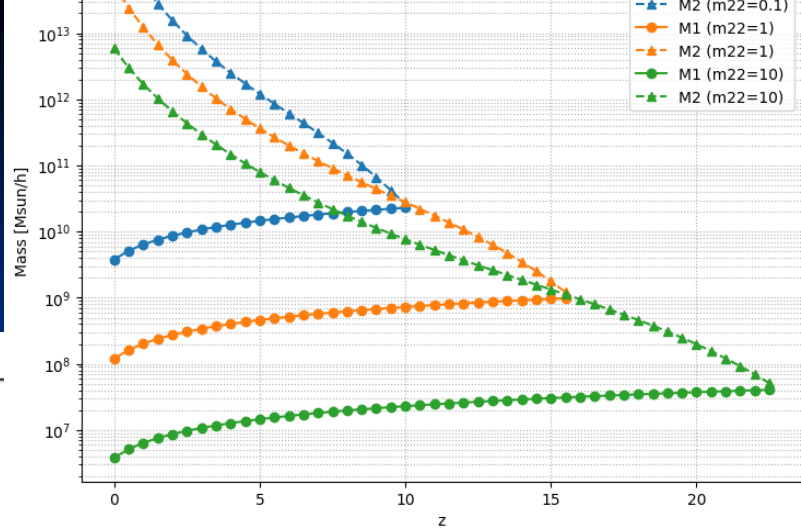
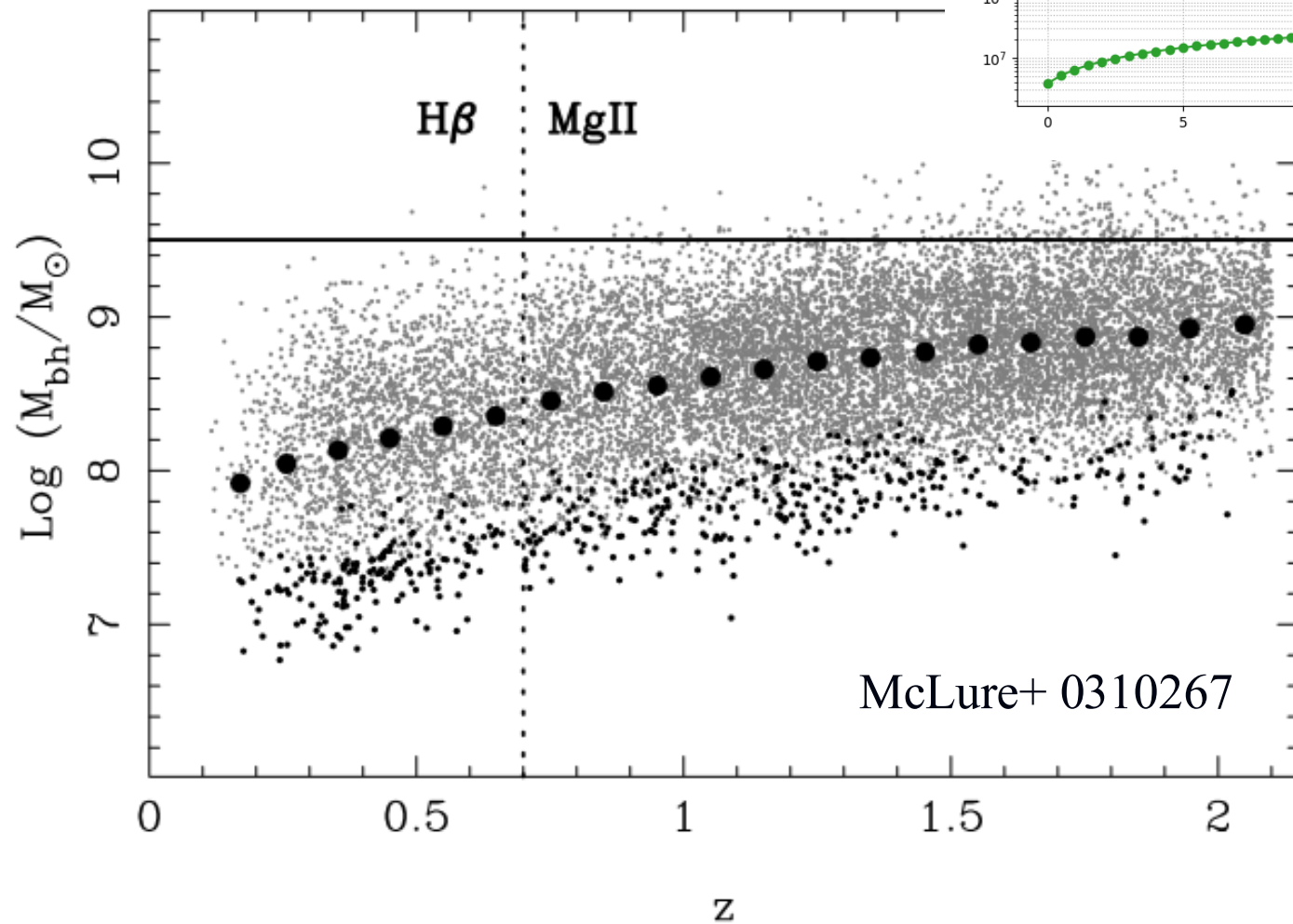
- increase $\rho(z)$ (extra radiation, EDE) just before ls
- but need more perturbation
- worsen S_8 tension, ad hoc matter, coincidence?

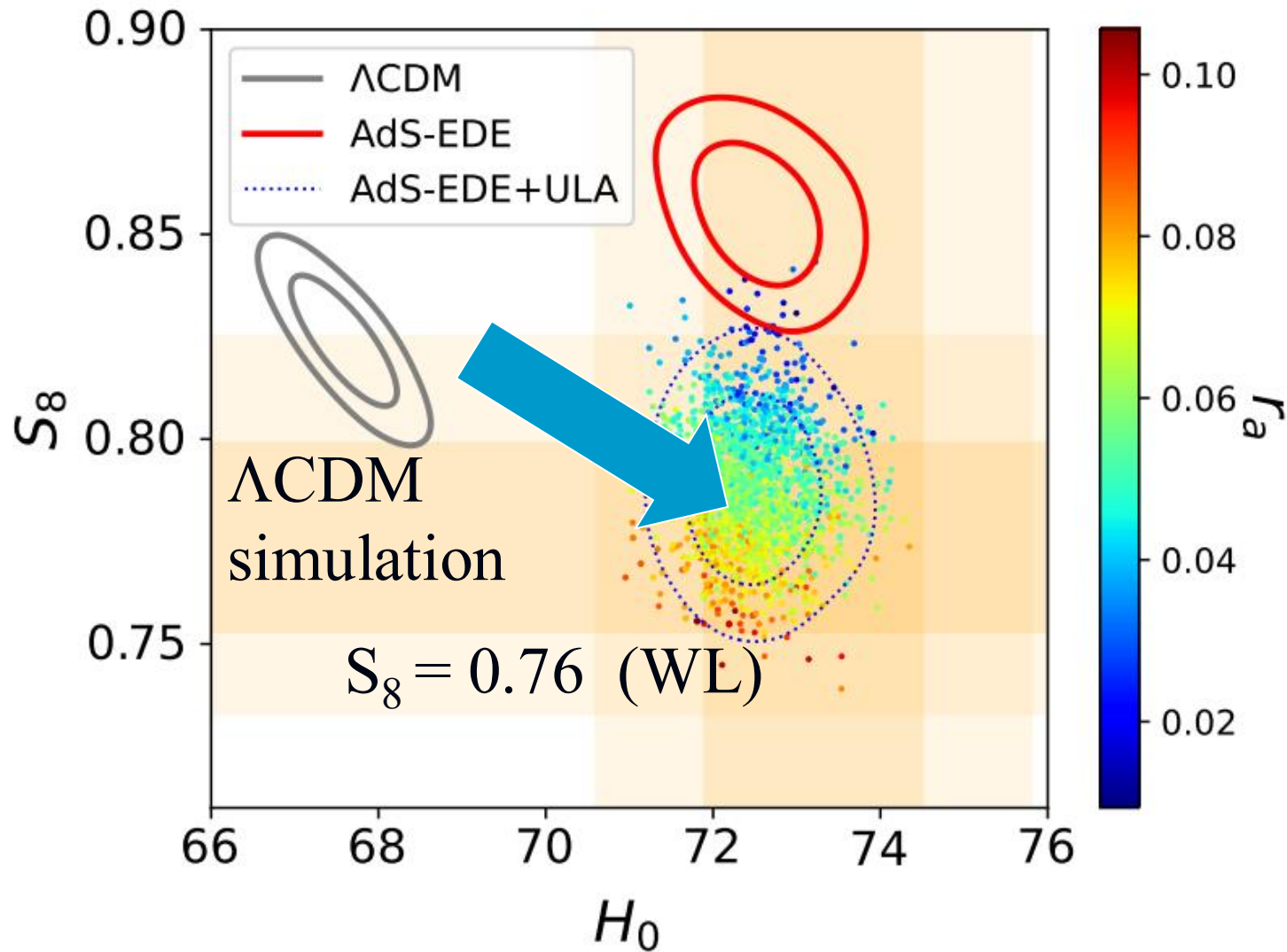
2) in late time solutions (r_s & D_A fixed, increase the integral)

- decrease $\rho(z)$ by decaying DM after ls and/or increase DE to increase the integral
- late evolution changes
- tensions with BAO, SN; null energy violation; fine tuning

$$D_A = \frac{c}{H_0} \int_0^{z_{ls}} \frac{dz}{[\rho(z)/\rho_0]^{1/2}}$$

Do we need both solutions?

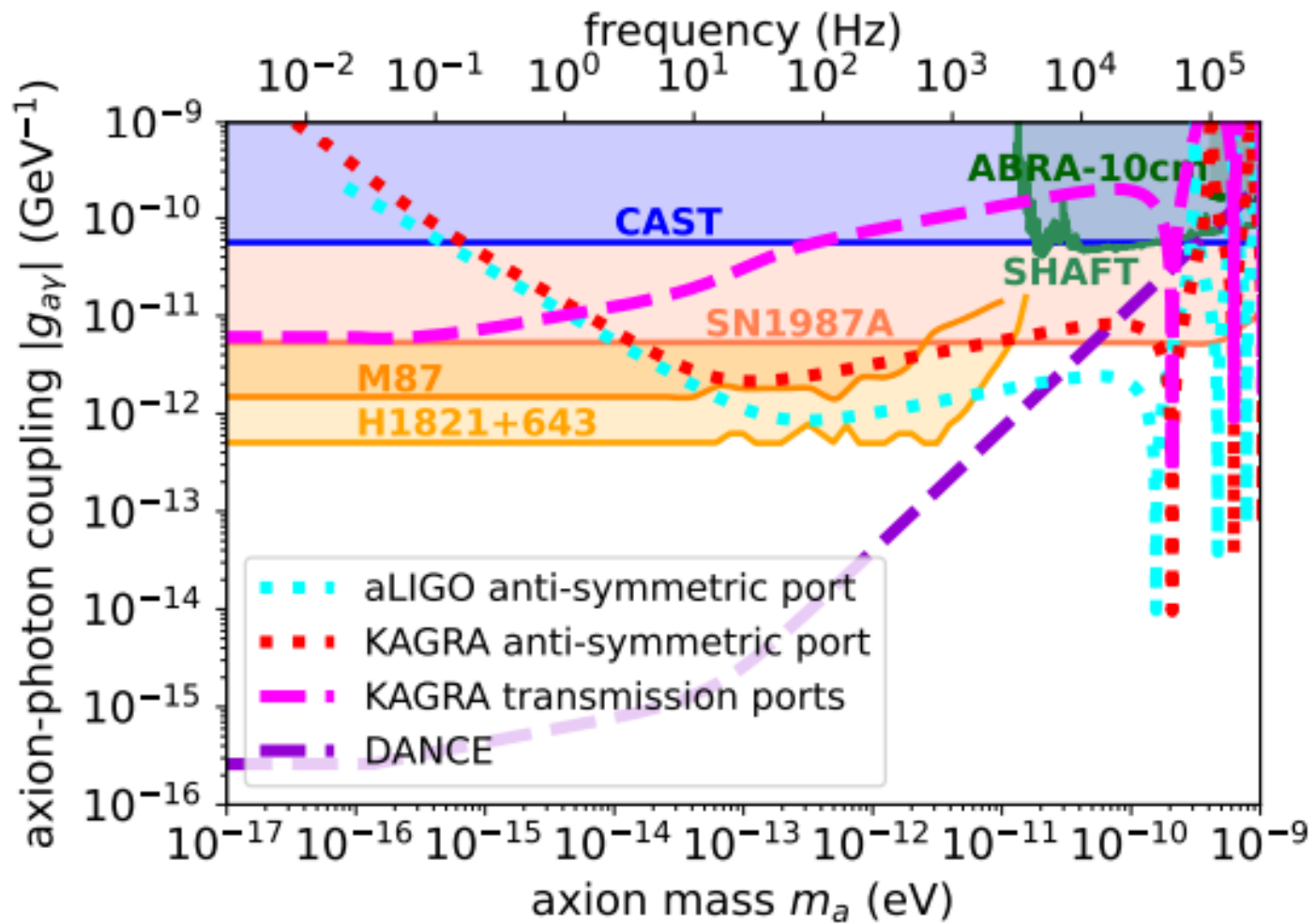


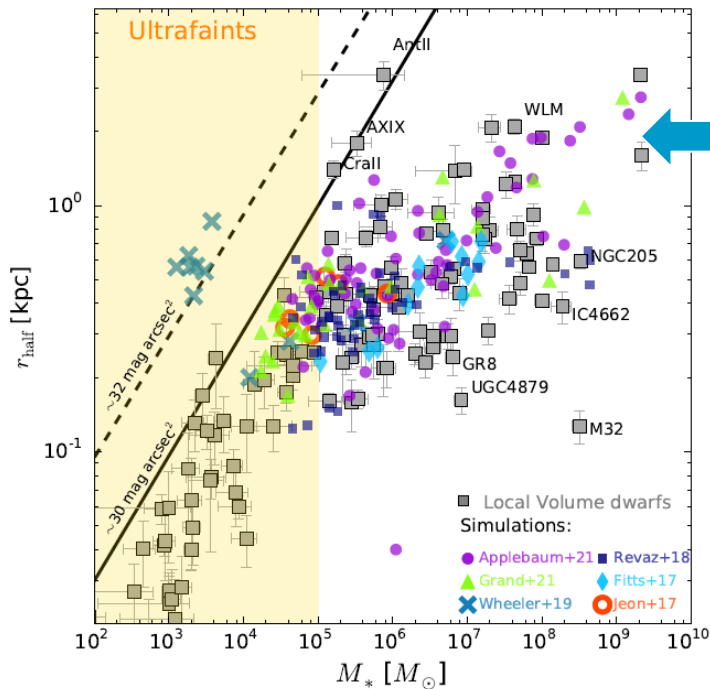
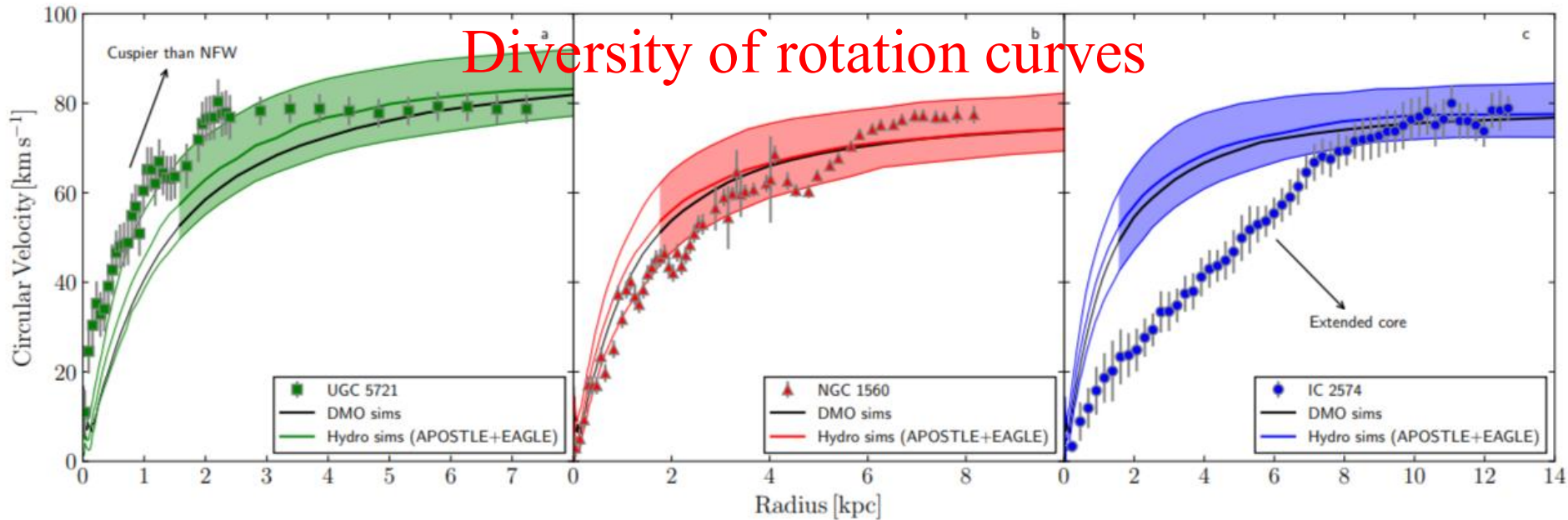


Can ULDM solve both H_0 and S_8 ? Ye+ 2107.13391
 → need extra DE & ULDM with $m < 10^{-24}$ eV

GW

2501.08930





a wide range of sizes at fixed stellar mass

No DM models is successful in reproducing RCs of all galaxies so far.

Sales+ 2206.05295

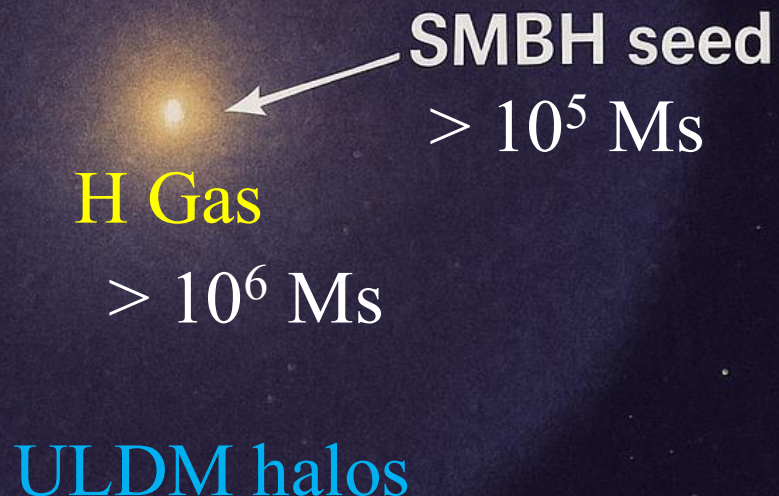
Other cosmological Constraints

- BEC phase transition before nucleosynthesis: $m < 10^2 \text{ eV}$
- field oscillation before equality $m > 10^{-28} \text{ eV}$
- Maximum mass of galaxies from BS theory
 - spiral $1.04 \times 10^{12} M_{\odot} < O(1) M_p^2/m \rightarrow m < O(1) 1.28 \times 10^{-22} \text{ eV}$
 - elliptical $1 \times 10^{13} M_{\odot} < O(1) M_p^2/m \rightarrow m < O(1) 1.28 \times 10^{-23} \text{ eV}$
- $\text{Ly}\alpha$ forest $m > 10^{-21} \text{ eV}$
- high-redshift galaxy luminosity $\rightarrow m > 1.2 \times 10^{-22} \text{ eV}$
- Stella subpopulations in Fornax $\rightarrow m < 1.1 \times 10^{-22} \text{ eV}$
- Ultra-faint dSphs $\rightarrow m \sim 3.7\text{-}5.6 \times 10^{-22} \text{ eV}$

fiducial value $m \sim 10^{-22} \text{ eV}$

Formation of SMBH seeds in ULDM halo

ULDM galactic cores $> 10^7 M_{\odot}$



in collaboration with Dongsu BAK (Seoul city Univ.)